

$$d\sigma = \frac{w(\vec{p}'|\vec{p})}{n_s |\vec{v}|} d^3 p' \quad ; \quad \vec{v}(\vec{p}) = \frac{\partial \mathcal{E}}{\partial \vec{p}} = \frac{\vec{p}}{m}$$

and \uparrow density of scatterers

$$\left(\frac{df}{dt}\right)_{\text{coll}} = \int d^3 p' \left\{ w(\vec{p}|\vec{p}') f(\vec{r}, \vec{p}', t) - w(\vec{p}'|\vec{p}) f(\vec{r}, \vec{p}, t) \right\}$$

This is called single particle scattering. We also have two particle scattering. Let

$$w(\vec{p}', \vec{p}'_1 | \vec{p}, \vec{p}_1) f(\vec{p}) f(\vec{p}_1) d^3 p d^3 p_1 d^3 p' d^3 p'_1 =$$

rate per (volume)² to scatter two particles
 $\{ |\vec{p} + d\vec{p}, \vec{p}_1 + d\vec{p}_1\rangle \} \rightarrow \{ |\vec{p}' + d\vec{p}', \vec{p}'_1 + d\vec{p}'_1\rangle \}$

Then

$$\left(\frac{df}{dt}\right)_{\text{coll}} = \int d^3 p_1 \int d^3 p' \int d^3 p'_1 \left\{ w(\vec{p}, \vec{p}_1 | \vec{p}', \vec{p}'_1) f_2(\vec{r}, \vec{p}'; \vec{r}, \vec{p}'_1; t) - w(\vec{p}', \vec{p}'_1 | \vec{p}, \vec{p}_1) f_2(\vec{r}, \vec{p}; \vec{r}, \vec{p}_1; t) \right\}$$

Approximation:

$$f_2(\vec{r}, \vec{p}; \vec{r}, \vec{p}_1; t) \approx f(\vec{r}, \vec{p}; t) f(\vec{r}, \vec{p}_1; t)$$

Then we have

$$\left(\frac{df}{dt}\right)_{\text{coll}} = \int d^3 p_1 \int d^3 p'_1 \int d^3 p''_1 \left\{ W(p, \vec{p}_1 | \vec{p}', \vec{p}'_1) f(\vec{r}, \vec{p}; t) f(\vec{r}, \vec{p}'_1; t) \right. \\ \left. - W(\vec{p}', \vec{p}'_1 | p, \vec{p}_1) f(\vec{r}, \vec{p}; t) f(\vec{r}, \vec{p}_1; t) \right\}$$

Detailed Balance

DB is a constraint on the equilibrium distⁿ $f^0(\vec{p})$.

For one-body scattering, the integrand of the collision term is

$$W(\vec{p} | \vec{p}') f(\vec{r}, \vec{p}', t) - W(\vec{p}' | \vec{p}) f(\vec{r}, \vec{p}, t)$$

Thus, a uniform and time-independent $f^0(\vec{p})$ satisfying

$$W(\vec{p} | \vec{p}') f^0(\vec{p}') = W(\vec{p}' | \vec{p}) f^0(\vec{p})$$

$$\Rightarrow \frac{f^0(\vec{p}')}{f^0(\vec{p})} = \frac{W(\vec{p}' | \vec{p})}{W(\vec{p} | \vec{p}')}$$

will make the integrand of $\left(\frac{df}{dt}\right)_{\text{coll}}$ vanish for all \vec{p} and \vec{p}' . Thus $f^0(\vec{p})$ is a solⁿ to

the BE with collisions. For two-body scattering, DB requires

$$W(p, \vec{p}_1 | \vec{p}', \vec{p}'_1) f^0(\vec{p}') f^0(\vec{p}'_1) \\ = W(\vec{p}', \vec{p}'_1 | p, \vec{p}_1) f^0(\vec{p}) f^0(\vec{p}_1)$$

Boltzmann's H - theorem

Although the microscopic laws governing a system may be time-reversal invariant, the Boltzmann equation itself establishes an arrow of time.

Define

$$h(\vec{r}, t) = \int d^3 p f(\vec{r}, \vec{p}, t) \log [f(\vec{r}, \vec{p}, t) a^3] \\ \vec{j}(\vec{r}, t) = \int d^3 p f(\vec{r}, \vec{p}, t) \log [f(\vec{r}, \vec{p}, t) a^3] \frac{d\vec{r}}{dt}$$

where a is any constant with dimensions of action. After some grinding (see § 8.3.7 in the lecture notes), we arrive at

$$\frac{\partial h(\vec{r}, t)}{\partial t} + \vec{\nabla} \cdot \vec{j}(\vec{r}, t) = \sigma(\vec{r}, t) < 0$$

with

$$\sigma(\vec{r}, t) = -\frac{1}{2} \int d^3 p \int d^3 p_1 \int d^3 p' \int d^3 p'_1 \left\{ \omega(\vec{p}', \vec{p}'_1 | \vec{p}, \vec{p}_1) \right. \\ \left. \times f(\vec{r}, \vec{p}', t) f(\vec{r}, \vec{p}'_1, t) (x \log x - x + 1) \right.$$

where

$$x(\vec{p}, \vec{p}_1, \vec{p}', \vec{p}'_1; \vec{r}, t) = \frac{f(\vec{r}, \vec{p}, t) f(\vec{r}, \vec{p}_1, t)}{f(\vec{r}, \vec{p}', t) f(\vec{r}, \vec{p}'_1, t)}$$

Integrating,

$$\frac{dH(t)}{dt} = \int d^3 r \sigma(\vec{r}, t) < 0$$

$$H(t) = \int d^3 r h(\vec{r}, t)$$

So $H(t)$ always decreases -- the arrow of time!

To make $\sigma = 0$ we set $x=1$ when $f(\vec{r}, \vec{p}, t) = f^0(\vec{p})$.

Then $f^0(\vec{p}) f^0(\vec{p}_1) = f^0(\vec{p}') f^0(\vec{p}'_1)$, i.e.

$$\log f^0(\vec{p}) + \log f^0(\vec{p}_1) = \log f^0(\vec{p}') + \log f^0(\vec{p}'_1)$$

But this means that $\log f^0(\vec{p})$ is a collisional invariant! For point

particles there are five such CIs:

$$1, p^x, p^y, p^z, \epsilon(\vec{p})$$

Thus we take

$$-\log f^0(\vec{p}) = \frac{\epsilon(\vec{p}) - \vec{V} \cdot \vec{p} - \mu}{k_B T}$$

This is a Boltzmann distⁿ in a moving frame.

