

Distⁿ of blackbody radiation:

EM radiation: $\nu = \frac{c}{\lambda} = \frac{ck}{2\pi}$, $\varepsilon_k = \hbar ck$

photons per unit frequency at temperature T :

$$N(\nu, T) d\nu = \frac{2V}{(2\pi)^3} \frac{d^3k}{e^{\hbar ck/k_B T} - 1}$$

Energy :

$$\mathcal{E}(\nu, T) = \frac{8\pi\hbar V}{c^3} \frac{\nu^3}{e^{\hbar\nu/k_B T} - 1}$$

$$= N(\nu, T) \times \hbar\nu \quad ; \quad k = \frac{2\pi\nu}{c} \text{ etc.}$$

Total energy :

$$E(T) = \int_0^\infty d\nu \mathcal{E}(\nu, T) = V \cdot \frac{\pi^2}{15} \frac{(k_B T)^4}{(\hbar c)^3}$$

Spectral density $\rho_{\mathcal{E}}(\nu, T)$:

$$\rho_{\mathcal{E}}(\nu, T) = \frac{\mathcal{E}(\nu, T)}{E(T)} = \frac{15}{\pi^4} \frac{\hbar}{k_B T} \cdot \frac{(\hbar\nu/k_B T)^3}{e^{\hbar\nu/k_B T} - 1}$$

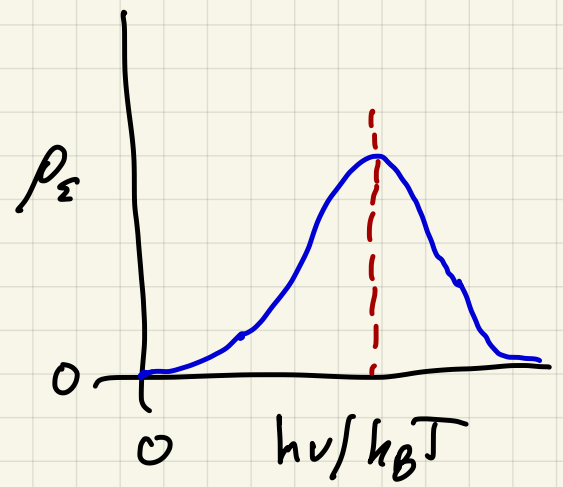
Scaled frequency $s \equiv \hbar\nu/k_B T$ (dim^{less}) $\Rightarrow$

$$\rho_\varepsilon = \text{const.} \cdot \frac{s^3}{e^s - 1}$$

$$\frac{d}{ds} \left(\frac{s^3}{e^s - 1} \right) \approx 0$$

$$\Rightarrow s^* = 2.82144$$

$$\nu_{\max} = s^* \cdot \frac{k_B T}{h} \propto T$$

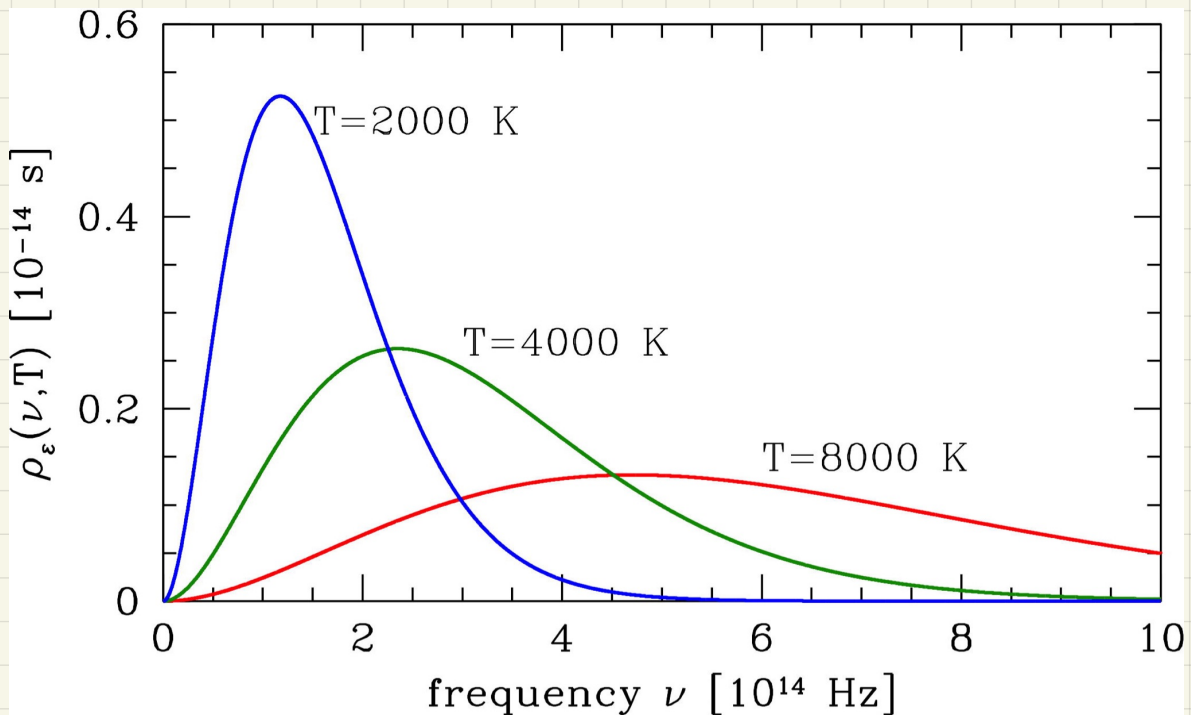


Photons γ ($d=3$)

- dispersion $\epsilon_{\vec{k}} = \hbar c |\vec{k}|$, $\vec{k} \in \mathbb{R}^3$
- $\hat{H} = \sum_{\vec{k}, \sigma} \epsilon_{\vec{k}} \hat{n}_{\vec{k}\sigma}$; $n_{\vec{k}\sigma} \in \{0, 1, \dots\}$
 $\nwarrow$ polarization
- bosons with $\mu = 0$, $g_d = 2$ (polarization)
- $n_{\vec{k}} = \langle \hat{n}_{\vec{k}} \rangle = \frac{1}{e^{\beta \epsilon_{\vec{k}}} - 1}$ occupancy of $|\vec{k}, \sigma\rangle$
- DOS $g(\epsilon) = \frac{\epsilon^2}{\pi^2 (\hbar c)^3}$ ($\epsilon \geq 0$)
- $n(T) = \int_0^\infty d\epsilon \frac{g(\epsilon)}{e^{\beta \epsilon} - 1} = \frac{23(3)}{\pi^2} \frac{(k_B T)^3}{(\hbar c)^3}$
- $p(T) = -k_B T \int_0^\infty d\epsilon g(\epsilon) \log(1 - e^{-\beta \epsilon})$
 $= \frac{23(4)}{\pi^2} \frac{(k_B T)^4}{(\hbar c)^3}$
- blackbody spectral density $\rho_\epsilon(\nu, T)$

$$\rho_{\varepsilon}(\nu, T) = \frac{15}{\pi^4} \frac{h}{k_B T} \frac{(h\nu/k_B T)^3}{e^{h\nu/k_B T} - 1}$$

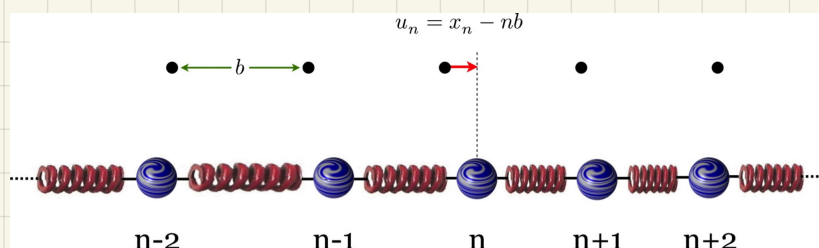
peaked at $\nu = 2.82 k_B T / h$



$\int_{\nu_1}^{\nu_2} \rho_{\varepsilon}(\nu, T) d\nu$ = fraction of $\varepsilon(T)$ in modes with $\nu \in [\nu_1, \nu_2]$

Phonons

- lattice vibrations in solids



Ideal Bose Gas

Classical NRIG: $Z = \frac{1}{N!} \left(\frac{V}{\lambda_T^3} \right)^N \Rightarrow$

$$F(T, V, N) = -k_B T \log Z(T, V, N)$$

$$= -Nk_B T \log \left(\frac{V}{N \lambda_T^3} \right) - Nk_B T$$

Use: $\log N! \approx N \log N - N + \mathcal{O}(\log N)$

Entropy:

$$S = - \left. \frac{\partial F}{\partial T} \right|_{V, N} = Nk_B \log \left(\frac{V}{N \lambda_T^3} \right) + \frac{5}{2} Nk_B$$

$$\lambda_T = \left(\frac{2\pi \hbar^2}{m k_B T} \right)^{1/2} \Rightarrow S \sim Nk_B \times \frac{3}{2} \log T$$

$$S(T \rightarrow 0^+) = -\infty$$

$$S(T=0) = k_B \log g_d$$

Noninteracting bosons : $\hat{h} |\alpha\rangle = \epsilon_\alpha |\alpha\rangle$
 1-body Hamiltonian,
 1-body spectrum

Occupancies:

$$n_\alpha = \langle \hat{n}_\alpha \rangle = \frac{1}{e^{\beta(\epsilon_\alpha - \mu)} - 1}$$

Number density:

$$n_0 = \frac{1}{z^{-1} - 1}$$

$$n(T, \mu) = -\frac{1}{V} \left. \frac{\partial \Omega}{\partial \mu} \right|_{T, V}$$

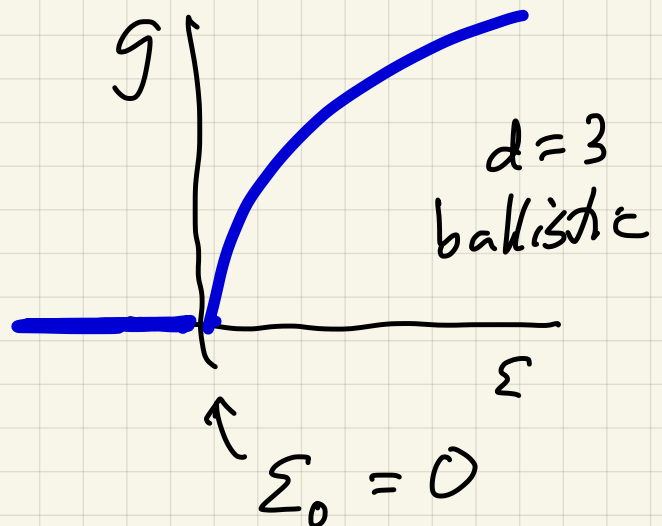
$$= \frac{1}{V} \sum_{\alpha} n_{\alpha} = \int d\epsilon \frac{g(\epsilon)}{e^{\beta(\epsilon - \mu)} - 1}$$

$$g(\epsilon) = \frac{1}{V} \sum_{\alpha} \delta(\epsilon - \epsilon_{\alpha}) \propto \epsilon^{\frac{d}{2} - 1}$$

E.g. particles with $S = \frac{1}{2}$

$$\epsilon(\mathbf{k}, \sigma) = \frac{\hbar^2 k^2}{2m} - \mu_0 B \sigma$$

$$\Rightarrow \epsilon_0 = -\mu_0 B$$



$$n(T, \mu) = \int_{\epsilon_0}^{\infty} d\epsilon \frac{g(\epsilon)}{e^{(\epsilon - \mu)/k_B T} - 1}$$

$$\epsilon_0 = 0$$

$$= \int_0^{\infty} d\epsilon \frac{g(\epsilon)}{z^{-1} e^{\epsilon/k_B T} - 1} = n(T, z)$$

Since $\mu < \epsilon_0 = 0 \Rightarrow z = e^{\beta\mu} \in [0, 1)$
 $\mu \in (-\infty, 0)$ $\mu = -\infty$ $\mu \rightarrow 0^-$

Clearly $n(T, z)$ increasing in z at fixed T

I.e. $\frac{\partial n}{\partial z} > 0$. What is maximum

allowed n ?

$$n_{\max} = n(T, z=1) = \int_0^{\infty} d\epsilon \frac{g(\epsilon)}{e^{\epsilon/k_B T} - 1}$$

Or, given n , there is a critical temperature $T_c(n)$ such that

$$n = \int_0^{\infty} d\varepsilon \frac{g(\varepsilon)}{e^{\varepsilon/k_B T_c} - 1}$$

Implicit eqn for $T_c(n)$.

Need to separate out $\hbar=0$ mode:

$$n(T, \mu) = \underbrace{\frac{1}{V} \frac{1}{z^{-1} - 1}}_{n_0(T)}$$

$$+ \underbrace{\int_0^{\infty} d\varepsilon \frac{g(\varepsilon)}{z^{-1} e^{\varepsilon/k_B T} - 1}}_{n'(T)}$$

$n_0(T)$: condensate
number density

$n'(T)$ = overcondensate

If n_0 finite, then

$$V n_0 = \frac{1}{z^{-1} - 1} \Rightarrow z^{-1} = 1 + \frac{1}{V n_0} \\ \approx e^{-\mu/k_B T}$$