28. The energy of the 2p to 1s Lyman transition is

$$E = (-13.60570 \text{ eV}) \left(\frac{1}{2^2} - \frac{1}{1^2}\right) = 10.20428 \text{ eV}$$

and its wavelength (in the absence of fine structure) is

$$\lambda = \frac{hc}{E} = \frac{1239.842 \text{ eV} \cdot \text{nm}}{10.20428 \text{ eV}} = 121.5022 \text{ eV}$$

With the fine structure energy splitting of  $4.5 \times 10^{-5}$  eV, the wavelength splitting is

$$\Delta \lambda = \frac{\lambda^2}{hc} \Delta E = \frac{(121.5 \text{ nm})^2}{1240 \text{ eV} \cdot \text{nm}} (4.5 \times 10^{-5} \text{ eV}) = 0.00054 \text{ nm}$$

The fine structure splits one level up by  $0.5\Delta E$  and the other down by the same amount, so the wavelengths are

$$\lambda + \frac{1}{2}\Delta\lambda = 121.5024 \text{ nm}$$
 and  $\lambda - \frac{1}{2}\Delta\lambda = 121.5019 \text{ nm}$ 

1. (a) For a 2*p* electron, n = 2, l = 1,  $m_l = 0,\pm 1$  and  $m_s = \pm \frac{1}{2}$ , so the possible sets of quantum numbers  $(n,l,m_l,m_s)$  are:

 $(2,1,+1,+\frac{1}{2}), (2,1,+1,-\frac{1}{2}), (2,1,0,+\frac{1}{2}), (2,1,0,-\frac{1}{2}), (2,1,-1,+\frac{1}{2}), (2,1,-1,-\frac{1}{2})$ (b) There are 6 possible sets of quantum numbers for each electron, so the total number of possibilities for 2 electrons is  $6 \times 6 = 36$ .

(c) The Pauli principle prevents the two sets from being identical. There will be 6 combinations in which the two sets are identical; eliminating these combinations leaves 30 allowed combinations.

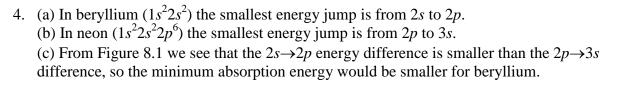
(d) Because the n values are different, the Pauli principle does not restrict the number of combinations, to there will be 36 possible combinations.

2. (a) The two electrons in the 1s level have  $m_s$  of  $\pm 1/2$  and  $\pm 1/2$ , so they do not contribute to the total  $m_s$ , and the same is true for the two electrons in the 2s level. In the 2p level, there are three different possible values of  $m_l$ , and for each of those values we can assign a set of quantum numbers with  $m_s = \pm 1/2$ , so the maximum possible value of the total  $m_s$  is  $\pm 3/2$ .

(b)  $(n, l, m_l, m_s) = (2, 1, +1, +1/2), (2, 1, 0, +1/2), (2, 1, -1, +1/2)$ 

(c) There is only one possible value of the total  $m_l$  in the states that maximize  $m_s$ , and from the states listed in (b) that value is +1 + 0 + (-1) = 0.

(d) We could maximize the total  $m_l$  by giving the first 2p electron  $m_l = +1$ , and the second electron can also have  $m_l = +1$  if we give these two electrons opposite values of  $m_s$ . The third electron cannot have  $m_l = +1$ , so we must assign it  $m_l = 0$  and the maximum total  $m_l$  is +2.



11. Singly ionized lithium has two electrons. When one of those is excited to a higher level, it is screened by the one electron remaining in the 1s level so  $Z_{eff} = 3 - 1 = 2$ . The expected energy when the outer electron is excited to the 2p level is

$$E_n = (-13.6 \text{ eV}) \frac{Z_{\text{eff}}^2}{n^2} = (-13.6 \text{ eV}) \frac{2^2}{2^2} = -13.6 \text{ eV}$$

which agrees very well with the measured value of -13.4 eV. When the outer electron is in the 3*d* level, its expected energy is

$$E_n = (-13.6 \text{ eV}) \frac{Z_{\text{eff}}^2}{n^2} = (-13.6 \text{ eV}) \frac{2^2}{3^2} = -6.0 \text{ eV}$$

in excellent agreement with the measured value.



16. Solving Equation 8.4 for Z with  $\Delta E = \frac{hc}{\lambda} = \frac{1240 \text{ eV} \cdot \text{nm}}{0.1940 \text{ nm}} = 6392 \text{ eV}$ , we obtain

$$Z = 1 + \sqrt{\frac{\Delta E}{10.2 \text{ eV}}} = 1 + \sqrt{\frac{6392 \text{ eV}}{10.2 \text{ eV}}} = 26$$

so the element is iron.

17. Ca (Z = 20): 
$$\Delta E = (10.2 \text{ eV})(Z - 1)^2 = (10.2 \text{ eV})(19)^2 = 3.68 \text{ keV}$$
  
Zr (Z = 40):  $\Delta E = (10.2 \text{ eV})(Z - 1)^2 = (10.2 \text{ eV})(39)^2 = 15.5 \text{ keV}$   
Hg (Z = 80):  $\Delta E = (10.2 \text{ eV})(Z - 1)^2 = (10.2 \text{ eV})(79)^2 = 63.7 \text{ keV}$ 

The values computed from Moseley's law are smaller than the measured values, and the discrepancy increases as *Z* increases.

27. (a) For the 3s outer electron of sodium, inserting  $E_3 = -5.14$  eV into Equation 8.1 gives

$$Z_{\rm eff} = n \sqrt{\frac{E_n}{-13.6 \,\mathrm{eV}}} = 3 \sqrt{\frac{-5.14 \,\mathrm{eV}}{-13.6 \,\mathrm{eV}}} = 1.84$$

The simple screening model predicts  $Z_{eff} = 1$ , so clearly the 3*s* electron is slightly penetrating the inner orbits and so is less screened by the inner electrons. (b) For the 4*f* state,

$$Z_{\rm eff} = n \sqrt{\frac{E_n}{-13.6 \,\mathrm{eV}}} = 4 \sqrt{\frac{-0.85 \,\mathrm{eV}}{-13.6 \,\mathrm{eV}}} = 1.00$$

so the screening is complete, with the 11 positive charges in the nucleus screened by the 10 electrons in the n = 1 and n = 2 shells.

30. The wavelength difference is  $\Delta \lambda = 0.59$  nm. By taking differentials of  $E = hc/\lambda$ , we can find the corresponding energy difference:

$$\Delta E = \frac{hc}{\lambda^2} \Delta \lambda = \frac{1240 \text{ eV} \cdot \text{nm}}{(590 \text{ nm})^2} (0.59 \text{ nm}) = 2.1 \times 10^{-3} \text{ eV}$$

This energy difference comes from the interaction of a magnetic field *B* with a magnetic moment that we assume is of the order of 1  $\mu_{\rm B}$ . The energy difference between the cases with the magnetic moment parallel to *B* and antiparallel to *B* is (see Figure 7.25)  $\Delta E = 2\mu_{\rm B}B$ , so

$$B = \frac{\Delta E}{2\mu_B} = \frac{2.1 \times 10^{-3} \text{ eV}}{2(5.8 \times 10^{-5} \text{ eV/T})} = 18 \text{ T}$$

This is quite a large magnetic field, of the order of the largest that can be produced in the laboratory with superconducting electromagnets.