

Class Intro

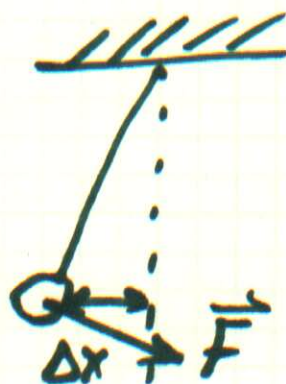
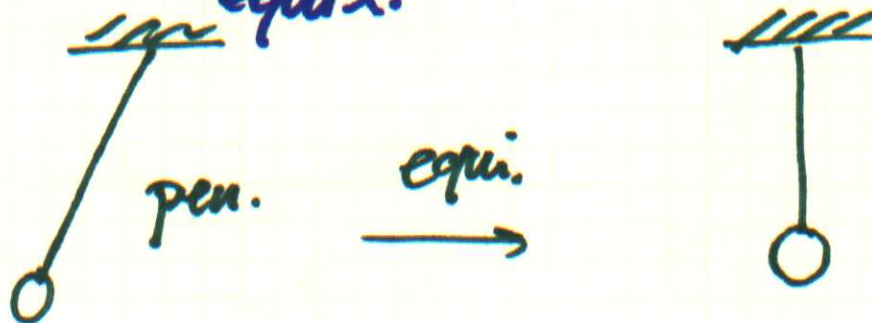
Oscillatory Motion

periodic motion: examples

↳ special kind in mech. sys.

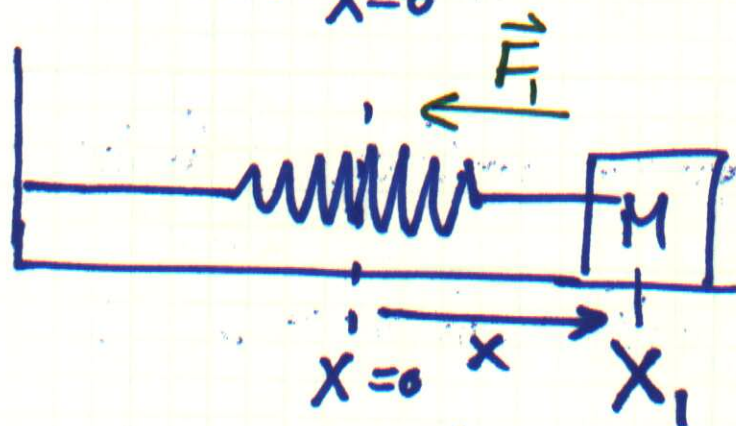
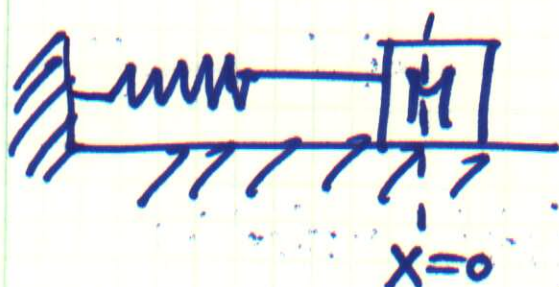
when Force is proportional
to the position of obj.

↳ Force always towards
equil.



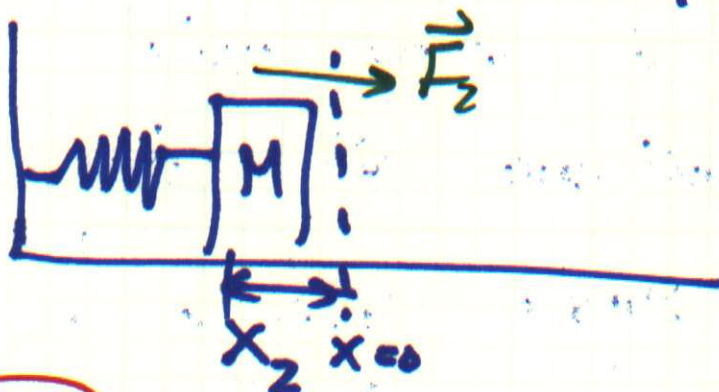
$$\vec{F} \propto \vec{\Delta x}$$

②



$$\vec{F} \propto \vec{x}$$

$$\vec{F}_1 = -k \vec{x}_1$$



$$\vec{F}_2 = -k \vec{x}_2$$

spring
const.

$$\vec{F} = -k \vec{x}$$

← Hooke's law

Newton's 1st law $\rightarrow \vec{F} = m\vec{a}$

$$\Rightarrow m \vec{a} = -k \vec{x}$$

$$m \frac{d^2 x}{dt^2} = -k x \xrightarrow{\text{ODE}} \frac{d^2 x}{dt^2} = -\frac{k}{m} x$$

③

$$a = \frac{d^2 x}{dt^2} = -\left(\frac{k}{m}\right) x \rightarrow \underline{a} \text{ propor. to } \underline{\text{negative of position}}$$

→ SHM whenever

Note:- $\vec{a}$ is not const. & changes during as $\vec{x}$ changes
(move?)

- no friction → SHM will continue forever

ODE

$$\frac{d^2 x}{dt^2} = -\frac{k}{m} x$$

** What should $\underline{x}$ be?

2nd deriv. of $\underline{x}$ is prop. to $\underline{x}$ w/ a (-) sign.

→ Sin & Cos

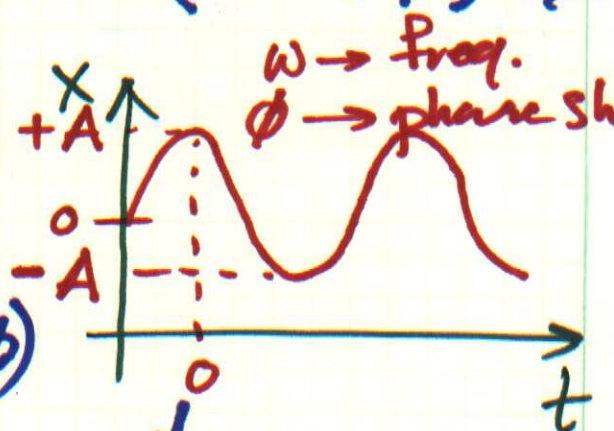
Try: $x(t) =$

(4)

Reminder:

$$\cos(\omega t + \phi) \rightarrow A \cos(\omega t + \phi)$$

try the gen. form: $x(t) = A \cos(\omega t + \phi)$



$$\frac{d^2 x(t)}{dt^2} = \frac{d}{dt} \left(\frac{dx(t)}{dt} \right) = \frac{d}{dt} (A \omega \sin(\omega t + \phi))$$

$$= \frac{d}{dt} (-A \omega^2 \cos(\omega t + \phi))$$

$x(t)$

$$\rightarrow \frac{d^2 x(t)}{dt^2} = -\omega^2 x(t)$$

$$a = -\left(\frac{k}{m}\right) x(t)$$

$$\omega = \sqrt{\frac{k}{m}}$$

SHM def.

(4)

$$X(t) = A \cos(\omega t + \phi)$$

Amplitude

(units: dist.)

Ang. frequency

(units: rad/s)

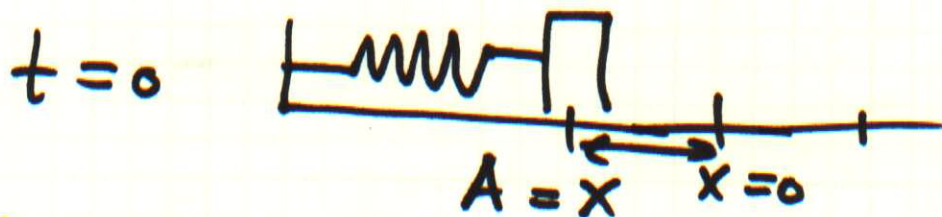
phase const.

(unit: phase (rad))

Amp.

A is max dist. from equil.

(if $x=A$ @ $t=0 \rightarrow \phi=0$)



ω

reminder $\rightarrow$ $\neq$ period T

time required to go complete 1 cycle

x, v, a are the same @

$t, t+T, t+2T, \dots$

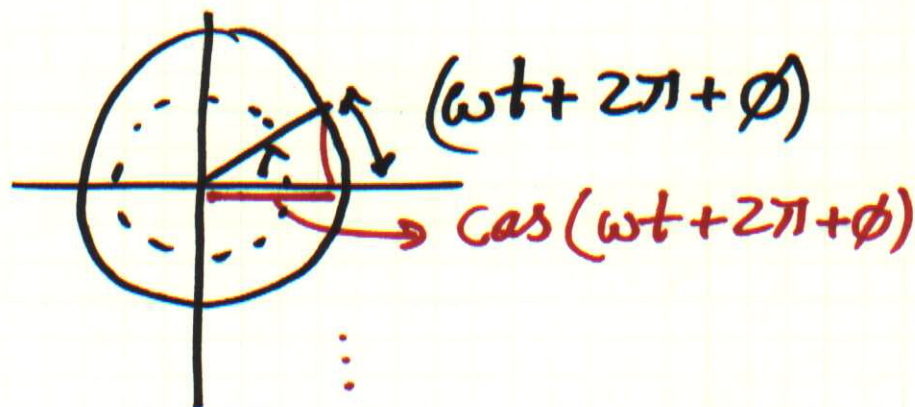
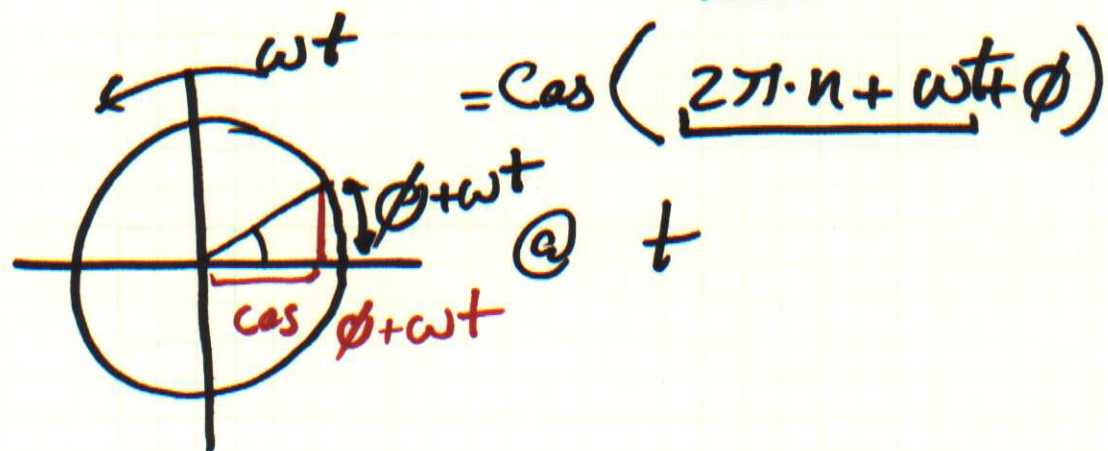
$$T = 2\pi / \omega \quad ; \quad f = 1/T = \omega / 2\pi$$

$$[T] = s \quad [f] = \frac{1}{s} = \text{Hz}$$

⑤

why $T = \frac{2\pi}{\omega}$?

$$\cos(\omega t + \phi) = \cos(2\pi + \omega t + \phi)$$



$$T = \frac{1}{f} = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{m}{k}}$$

$f = \dots$

only
dep. on
mass & k
nothing else

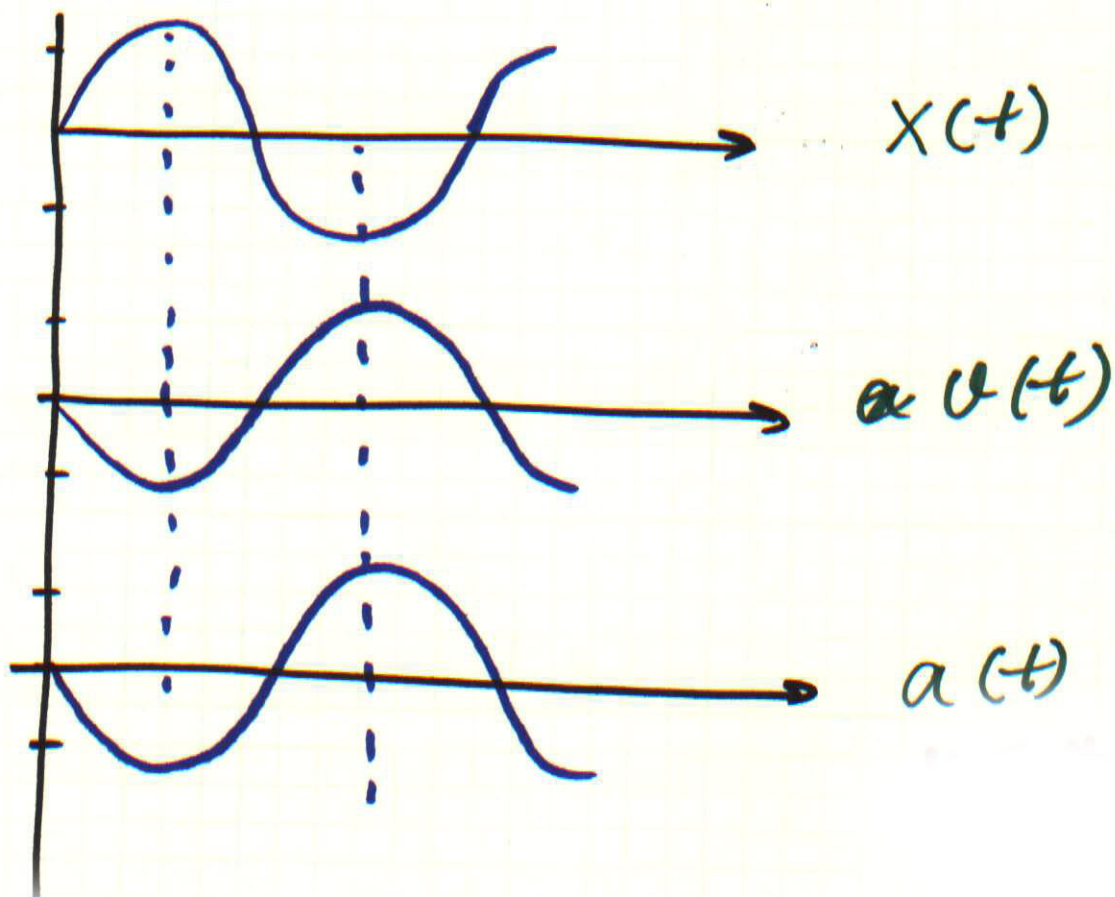
$$\vec{X}(t) \xrightarrow{dx/dt} \vec{V}(t) \xrightarrow{dv/dt} \vec{a} \quad (6)$$

$$\vec{X}(t) = A \cos(\omega t + \phi)$$

$$\vec{V}(t) = -\omega A \sin(\omega t + \phi)$$

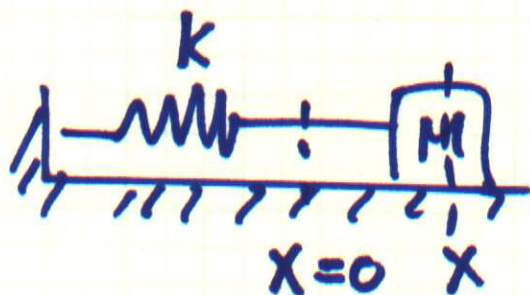
$$\vec{a}(t) = -\omega^2 A \cos(\omega t + \phi)$$

Note: not uniform acc. ^{time dep.}



Energy in SHM

(7)



(no friction)

$\downarrow$
E in const.

E?

$$E = K + U \quad \leftarrow \begin{array}{l} \text{pot. } E \\ \text{kinetic } E \end{array}$$

$$K = \frac{1}{2} m v^2 = \frac{1}{2} m \left(\frac{dx}{dt} \right)^2$$

$$= \frac{1}{2} m \left[-\omega A \sin(\omega t + \phi) \right]^2$$

$$= \frac{1}{2} m A^2 \omega^2 \sin^2(\omega t + \phi)$$

U=?

$$F = -\frac{dU}{dx} \rightarrow U = \int F dx$$

$$F = -kx \xrightarrow[\frac{dU}{dt}]{\text{int.}} U = \frac{1}{2} kx^2$$

$$U = \frac{1}{2} k x^2 = \frac{1}{2} k A^2 \cos^2(\omega t + \phi)$$

$$K = \frac{1}{2} m A^2 \omega^2 \sin^2(\alpha) \quad \textcircled{8}$$

$$\rightarrow m \omega^2 = m \left(\sqrt{\frac{k}{m}} \right)^2 = m \cdot \frac{k}{m} = k$$

$$K = \frac{1}{2} \cancel{A^2} k A^2 \sin^2(\alpha)$$

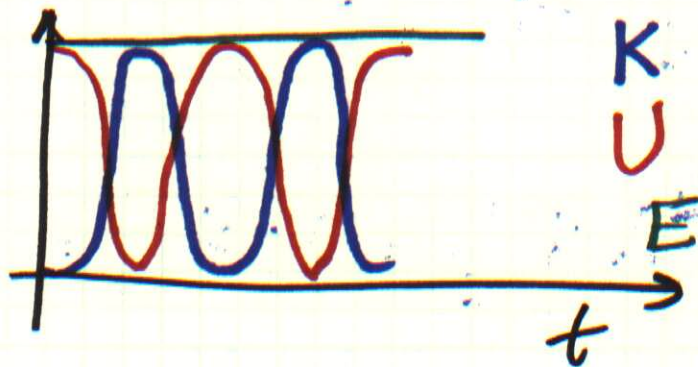
$$U = \frac{1}{2} k A^2 \cos^2(\alpha)$$

$$E = K + U = \frac{k}{2} A^2 \underbrace{(\sin^2 \alpha + \cos^2 \alpha)}_1$$

$$E = \frac{1}{2} k A^2$$

$\Rightarrow$ total mech. E is const.

$\hookrightarrow$ cont. transferred between K & U .



$$E = \frac{1}{2} k A^2 \quad \text{A can be solve for A} \rightarrow A = \pm \sqrt{\frac{2E}{k}} \quad (9)$$

or

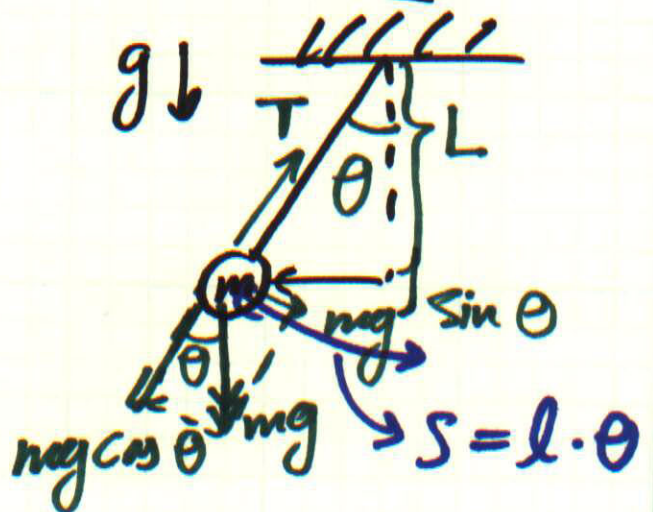
$$E = K + U = \frac{1}{2} m v^2 + \frac{1}{2} k x^2 = \frac{1}{2} k A^2$$

solve for v :

$$v^2 = k \frac{A^2 - x^2}{m} \rightarrow v = \pm \sqrt{\frac{k}{m} (A^2 - x^2)}$$

Simple Pendulum

- consist
- consist of mass & (m)
 - light string (l)
 - 2D motion



tangential component of (10)
the grav. force is the restoring
force.

$$F_t = ma_t \rightarrow -mg \sin \theta = m \frac{d^2 s}{dt^2}$$

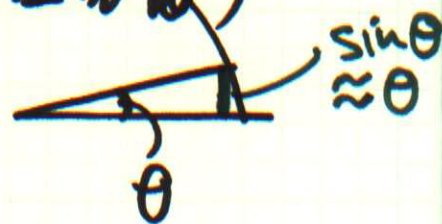
$$-g \sin \theta = \frac{d^2}{dt^2} (l \cdot \theta) \quad l \rightarrow \text{const.}$$
$$= l \frac{d^2 \theta}{dt^2}$$

$$\frac{d^2 \theta}{dt^2} = -\left(\frac{g}{l}\right) \sin \theta$$

$\theta \ll 1$, small θ : ($\theta < 2\pi$)

$$\sin \theta \approx \theta$$

[note: θ in Rad.]



$$\frac{d^2 \theta}{dt^2} \approx -\left(\frac{g}{l}\right) \theta \quad \xrightarrow[\text{as SHM}]{\text{same form}} \frac{d^2 \theta}{dt^2} = -\omega^2 \theta$$

$$\frac{d^2\theta}{dt^2} = -\omega^2 \theta \quad \text{or} \quad \omega = \sqrt{g/l} \quad \textcircled{11}$$

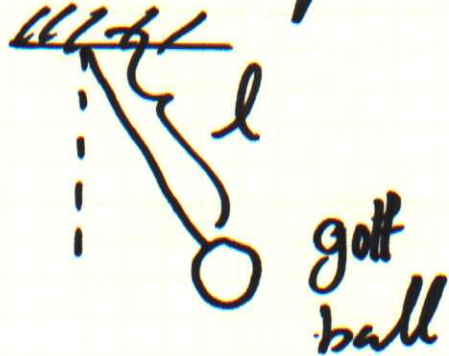
$$\theta(t) = \theta_{\max} \cos(\omega t + \phi)$$

$$\rightarrow T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{l}{g}}$$

$\Rightarrow$ T only dep.

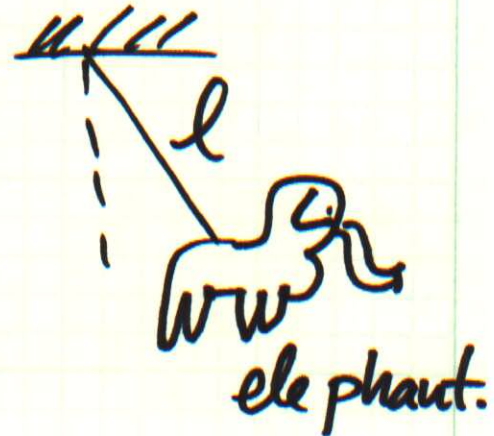
only ^{on} l & g

* ind. of mass (m)



&

SAME



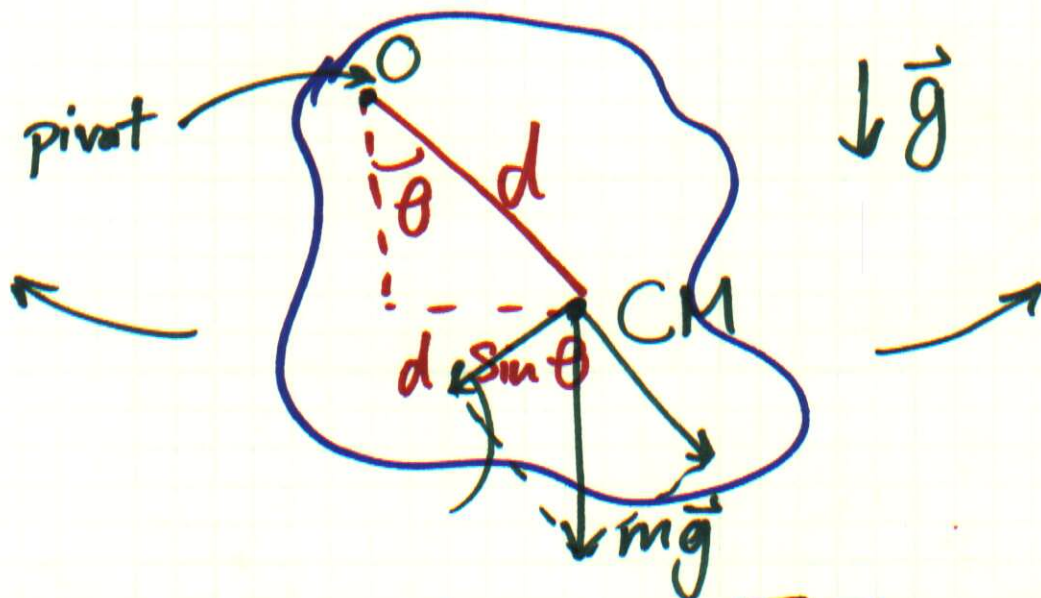
(quest. & examples)

AF 12.16

Physical Pendulum

(12)

- Obj. not approx. as a particle
- Obj. ^{not} hanging oscillating around a fixed axis that does not pass through its center of mass.
- - not a simple pendulum!!
- rigid obj model



Newton's 2nd law: $\sum_i \vec{F}_i = m\vec{a}$

$$\sum_i \tau_i = I \cdot \alpha$$

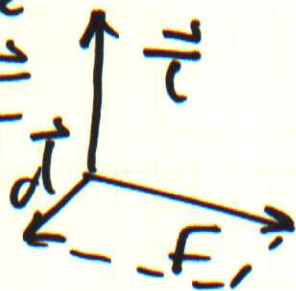
$$\vec{\tau} = I \cdot \vec{\alpha}$$

$$\alpha = d^2\theta / dt^2 \quad (13)$$

$I \equiv$ moment of inertia

$\tau \equiv$ torque

$$\tau = \vec{d} \times \vec{F}$$



$$\vec{\tau} = -mg \vec{d} \times \vec{d} = -mgd \sin \theta$$

$$\rightarrow -mgd \sin \theta = I \frac{d^2\theta}{dt^2}$$

$$\frac{d^2\theta}{dt^2} = -\left(\frac{mgd}{I}\right) \sin \theta$$

same form SHM

$$\frac{d^2\theta}{dt^2} = -\omega^2 \sin \theta$$

ω

$$\omega = \sqrt{\frac{mgd}{I}}$$

$$\Rightarrow T = 2\pi \sqrt{\frac{I}{mgd}}$$

* Phys. Pend. $\rightarrow$ measure I
(knowing d)

(14)

* If $I = md^2$ then phys. pend.
is the same as ~~sy~~ simple pend.
(mass all at the center)

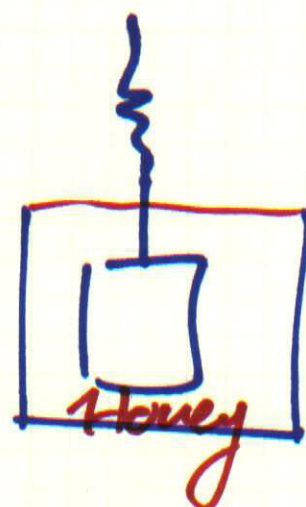
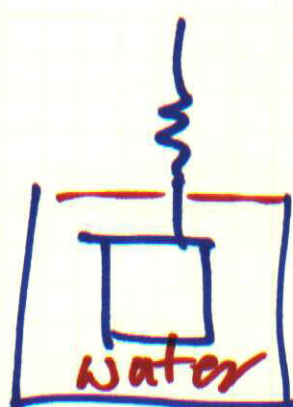
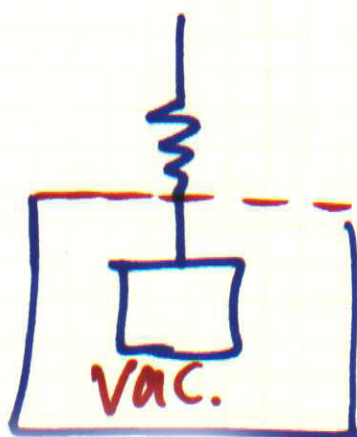
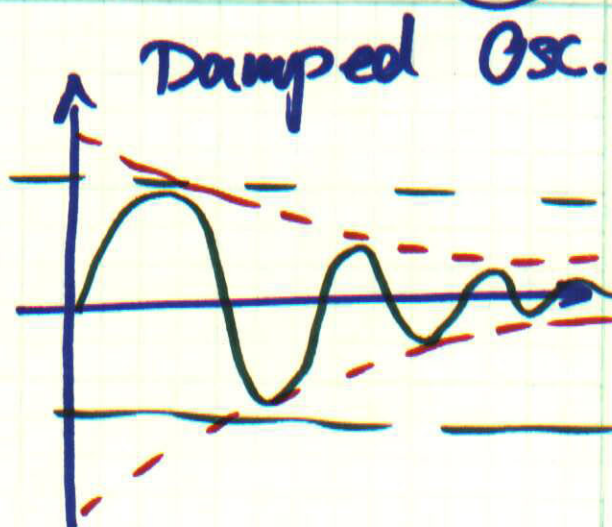
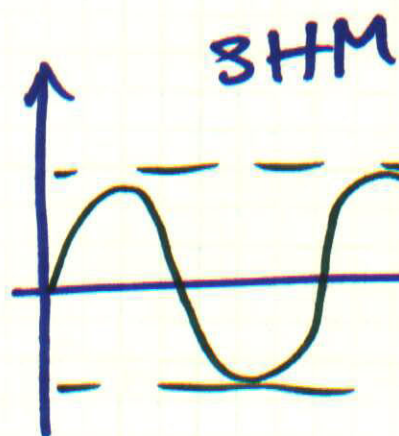
Damped Oscillation:

* In real sys. non cons. forces
are present.

$\rightarrow$ not an ideal sys.

$\rightarrow$ Friction present

* E is not a const anymore
 $\Rightarrow$ damped osc.



damped

presence of resistive force:

$$\vec{R} = -b \vec{v}$$

const.

$$\sum F_x = -kx - b\dot{x} = m a_x$$

restoring F

resistive force

$$-kx - b\dot{x} = m \frac{d^2x}{dt^2}$$

$$\Rightarrow m \frac{d^2x(t)}{dt^2} + b \frac{dx(t)}{dt} + kx(t) = 0$$

* the sol. of the above equation is beyond the scope of this class.

* we can solve this for the special case $b < \sqrt{4m\omega}$

$$x = \underbrace{\left[A e^{-(b/2m)t} \right]}_{A(t)} \cos(\omega t + \phi)$$

$$\omega = \sqrt{\frac{k}{m} - \left(\frac{b}{2m}\right)^2}$$

* $A(t)$ decays exp. w/ time

* motion ultimately stops

$$* \omega = \sqrt{\omega_0 - \left(\frac{b}{2m}\right)^2}$$

↑
natural freq.

(17)

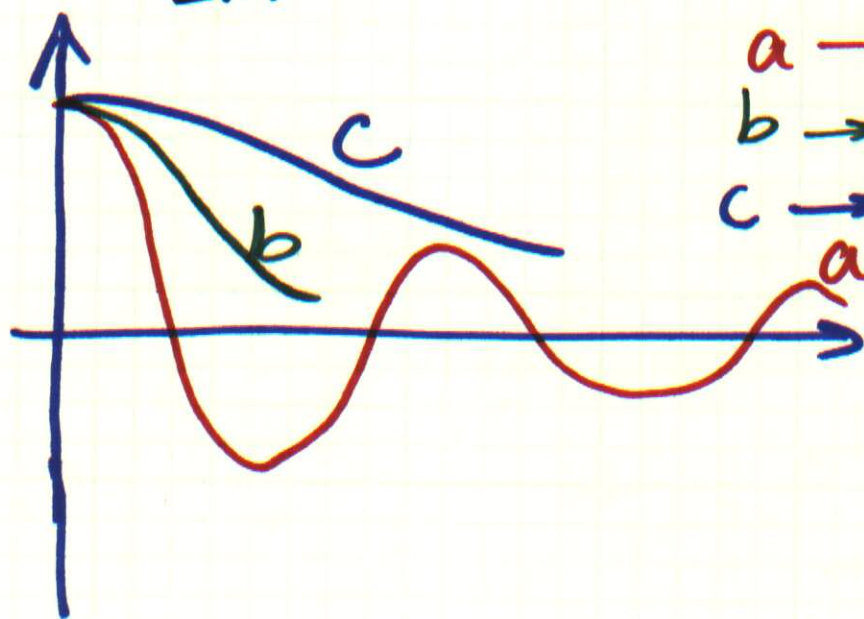
* If $R_{\max} = b v_{\max} < k A_{\max}$

→ underdamped

* If $b = b_{\text{crit.}} \xrightarrow{\text{s.t.}} \frac{b_c}{2m} = \omega_0$

→ no osc. (critic. damped)

* If $b/2m > \omega_0 \rightarrow \text{Over damped}$



a → und. damp.
b → crit. damp.
c → over damp.

(Q & E)

Forced Osc.

(18)

- * Compensate loss of energy by applying ext. force
- * A will remain const. if $\text{energy input/cycle} = \text{energy loss/cycle}$

stationary obj $\xrightarrow{\text{driving force}}$ osc. starts.

$\xrightarrow{\text{long enough time}}$ $E_{\text{driving}} = E_{\text{lost. to internal}}$

$\longrightarrow$ steady state

$\longrightarrow$ const A (ampl.)

$$m \frac{d^2 x(t)}{dt^2} + b \frac{dx(t)}{dt} + k x(t) = F(t) = F_0 \sin \omega t$$

$$A = \frac{F_0 / m}{\sqrt{(\omega^2 - \omega_0^2)^2 + \left(\frac{b\omega}{m}\right)^2}}$$

$$\left(\begin{array}{l} x(t) = \dots \\ \omega_0 = \dots \end{array} \right)$$

(19)

~~As $\omega \rightarrow \omega_0$~~

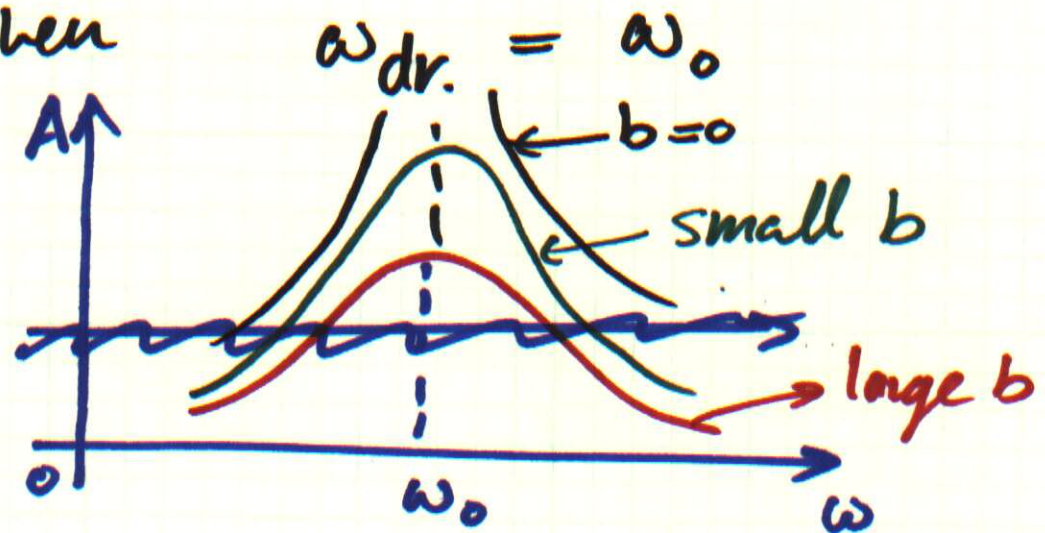
As $\omega \rightarrow \omega_0$ $A \rightarrow A_{\max}$

this dramatic increase in A
is known as Resonance

Resonance:

- when

-



(Q & E)