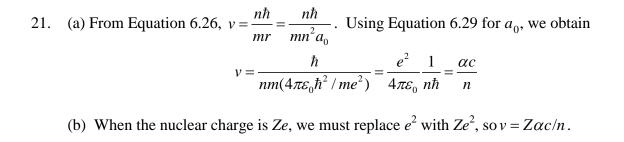
19.
$$r_3 = 9a_0 = 9(0.0529 \text{ nm}) = 0.476 \text{ nm}$$

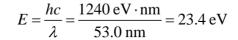
$$v = \frac{n\hbar}{mr} = c \frac{n\hbar c}{mc^2 r} = c \frac{3(1240 \,\text{eV} \cdot \text{nm})/2\pi}{(0.511 \times 10^6 \,\text{eV})(0.476 \,\text{nm})} = 2.43 \times 10^{-3} c = 7.30 \times 10^5 \,\text{m/s}$$

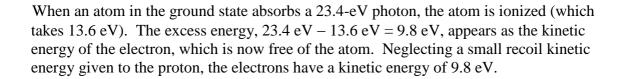
$$U = -\frac{e^2}{4\pi\varepsilon_0} \frac{1}{r} = -\frac{1.440 \text{ eV} \cdot \text{nm}}{0.476 \text{ nm}} = -3.02 \text{ eV}$$

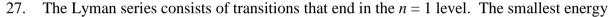
$$K = \frac{e^2}{8\pi\varepsilon_0} \frac{1}{r} = \frac{1.440 \,\mathrm{eV} \cdot \mathrm{nm}}{2(0.476 \,\mathrm{nm})} = 1.51 \,\mathrm{eV}$$

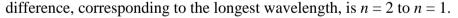


24. The photon energy of the incident light is









$$\Delta E = E_2 - E_1 = (-13.6 \text{ eV})2^2 \left(\frac{1}{2^2} - \frac{1}{1^2}\right) = 40.8 \text{ eV}$$
$$\lambda = \frac{hc}{\Delta E} = \frac{1240 \text{ eV} \cdot \text{nm}}{40.8 \text{ eV}} = 30.4 \text{ nm}$$

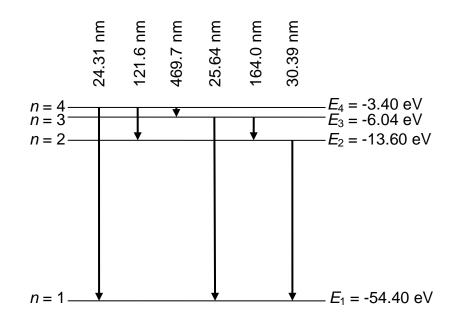
The largest energy difference would correspond to transitions from $n = \infty$ to n = 1:

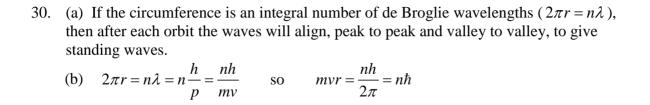
$$\Delta E = E_{\infty} - E_1 = (-13.6 \text{ eV})2^2 \left(0 - \frac{1}{1^2}\right) = 54.4 \text{ eV}$$

$$\lambda = \frac{hc}{\Delta E} = \frac{1240 \text{ eV} \cdot \text{nm}}{54.4 \text{ eV}} = 22.8 \text{ nm}$$

28. Using Equation 6.38, we have $E_n = (-13.6 \text{ eV})Z^2/n^2 = (-54.4 \text{ eV})/n^2$, so $E_1 = -54.40 \text{ eV}$, $E_2 = -13.60 \text{ eV}$, $E_3 = -6.04 \text{ eV}$, $E_4 = -3.40 \text{ eV}$. The possible transitions are:

$\lambda = hc / \Delta E = 24.31 \mathrm{nm}$
$\lambda = hc / \Delta E = 121.6 \text{ nm}$
$\lambda = hc / \Delta E = 469.7 \text{ nm}$
$\lambda = hc / \Delta E = 25.64 \text{ nm}$
$\lambda = hc / \Delta E = 164.0 \text{ nm}$
$\lambda = hc / \Delta E = 30.39 \text{ nm}$





35. (a) The frequency of revolution is given by Equation 6.41:

$$f_n = \frac{me^4}{32\pi^3\varepsilon_0^2\hbar^3} \frac{1}{n^3} = \frac{1}{\pi\hbar} \frac{me^4}{32\pi^2\varepsilon_0^2\hbar^2} \frac{1}{n^3} = \frac{13.6 \text{ eV}}{\pi\hbar} \frac{1}{n^3} = \frac{6.58 \times 10^{15} \text{ Hz}}{n^3}$$

A similar calculation gives the radiation frequency from Equation 6.42:

$$f = \frac{me^4}{64\pi^3\varepsilon_0^2\hbar^3} \frac{2n-1}{n^2(n-1)^2} = \frac{13.6 \text{ eV}}{2\pi\hbar} \frac{2n-1}{n^2(n-1)^2} = (6.58 \times 10^{15} \text{ Hz}) \frac{2n-1}{2n^2(n-1)^2}$$

For n = 10, we get $f_n = 6.58 \times 10^{12}$ Hz and $f = 7.72 \times 10^{12}$ Hz.

(b) For
$$n = 100$$
, $f_n = 6.58 \times 10^9$ Hz and $f = 6.68 \times 10^9$ Hz.

(c) For
$$n = 1000$$
, $f_n = 6.58 \times 10^6$ Hz and $f = 6.59 \times 10^6$ Hz.

(d) For n = 10,000, $f_n = 6.58 \times 10^3$ Hz and $f = 6.58 \times 10^3$ Hz. Note how f approaches f_n as n becomes large, in accordance with the correspondence principle.

36. The Rydberg constant in ordinary hydrogen is

$$R_{\rm H} = R_{\infty} \left(1 + \frac{m}{M_{\rm H}} \right) = R_{\infty} \left(1 + \frac{5.48580 \times 10^{-4} \text{ u}}{1.007825 \text{ u}} \right) = R_{\infty} (1.000544)$$

and in "heavy" hydrogen or deuterium:

$$R_{\rm D} = R_{\infty} \left(1 + \frac{m}{M_{\rm D}} \right) = R_{\infty} \left(1 + \frac{5.48580 \times 10^{-4} \text{ u}}{2.104102 \text{ u}} \right) = R_{\infty} (1.000272)$$

From Equation 6.33 the difference in wavelengths for the first line of the Balmer series (n = 3 to n = 2) is

$$\lambda_{\rm D} - \lambda_{\rm H} = \left(\frac{1}{R_{\rm D}} - \frac{1}{R_{\rm H}}\right) \left(\frac{3^2 2^2}{3^2 - 2^2}\right) = \frac{7.2}{1.09737 \times 10^7 \,{\rm m}^{-1}} \left(\frac{1}{1.000272} - \frac{1}{1.000544}\right) = 0.178 \,{\rm nm}$$