

PHYSICS 210B : NONEQUILIBRIUM STATISTICAL PHYSICS
HW ASSIGNMENT #4 SOLUTIONS

(1) Evaluate, for general α , the averages of the following stochastic integrals:

$$\int_0^t dW(s) W(s) s \quad , \quad \int_0^t dW(s) W^3(s) e^{-\lambda s} \quad , \quad \int_0^t dW(s) W^{2k+1}(s) \quad .$$

Solution:

We evaluate the general stochastic integral,

$$\begin{aligned} I_k[\phi] &= \int_0^t dW(s) W^{2k+1}(s) \phi(s) \\ &= \sum_{j=0}^{N-1} \left[(1-\alpha) W_j^{2k+1} \phi_j + \alpha W_{j+1}^{2k+1} \phi_{j+1} \right] (W_{j+1} - W_j) \quad . \end{aligned}$$

Therefore,

$$\begin{aligned} \langle I_k[\phi] \rangle &= \alpha \sum_{j=0}^{N-1} \phi_{j+1} \left(\langle W_{j+1}^{2k+2} \rangle - \langle W_j W_{j+1}^{2k+1} \rangle \right) \\ &= \frac{(2k+2)!}{2^{k+1}(k+1)!} \cdot \alpha \int_0^t ds s^k \phi(s) \quad . \end{aligned}$$

Thus,

$$\begin{aligned} \left\langle \int_0^t dW(s) W(s) s \right\rangle &= \frac{1}{2} \alpha t^2 \\ \left\langle \int_0^t dW(s) W^3(s) e^{-\lambda s} \right\rangle &= 3\alpha \int_0^t ds s e^{-\lambda s} = 3\alpha \left[\frac{1}{\lambda^2} (1 - e^{-\lambda t}) - \frac{t}{\lambda} e^{-\lambda t} \right] \\ \left\langle \int_0^t dW(s) W^{2k+1}(s) \right\rangle &= \frac{(2k+2)!}{2^{k+1}(k+1)!} \cdot \frac{\alpha t^{k+1}}{k+1} \quad . \end{aligned}$$

Note that in the limit $\lambda \rightarrow 0$ the second integral yields $\frac{3}{2}\alpha t^2$, which agrees with the results from the third expression.

(2) Derive Eqn. 3.107 of the lecture notes.

Solution:

Starting from

$$\begin{aligned} du &= -\beta u dt + \beta dW(t) \\ \frac{dz}{dt} &= i\nu z + i\lambda u(t) z, \end{aligned}$$

we obtained the solution

$$z(t) = z(0) \exp \left\{ i\nu t + \frac{i\lambda}{\beta} u(0) (1 - e^{-\beta t}) + i\lambda \int_0^t dW(s) (1 - e^{-\beta(t-s)}) \right\}.$$

We wish to compute the quantity $Y(s) = \lim_{t \rightarrow \infty} \langle z(t+s) z^*(t) \rangle$. We therefore have

$$\begin{aligned} \langle z(t+s) z^*(t) \rangle &= |z(0)|^2 e^{i\nu s} \exp \left\{ \frac{i\lambda}{\beta} u(0) e^{-\beta t} (e^{-\beta s} - 1) \right\} \times \\ &\quad \exp \left\{ -\frac{\lambda^2}{2} \left\langle \left(\int_0^t dW(\sigma) [1 - e^{-\beta(t-\sigma)}] - \int_0^{t+s} dW(\sigma) [1 - e^{-\beta(t+s-\sigma)}] \right)^2 \right\rangle \right\} \end{aligned}$$

We now invoke the result of Eqn. 3.33,

$$\left\langle \int_0^t dW(s) F(s) \int_0^{t'} dW(s') G(s') \right\rangle = \int_0^{\tilde{t}} ds F(s) G(s),$$

where $\tilde{t} = \min(t, t')$, to obtain

$$\begin{aligned} &\left\langle \left(\int_0^t dW(\sigma) [1 - e^{-\beta(t-\sigma)}] - \int_0^{t+s} dW(\sigma) [1 - e^{-\beta(t+s-\sigma)}] \right)^2 \right\rangle \\ &= t - \frac{2}{\beta} (1 - e^{-\beta t}) + \frac{1}{2\beta} (1 - e^{-2\beta t}) + (t+s) - \frac{2}{\beta} (1 - e^{-\beta(t+s)}) + \frac{1}{2\beta} (1 - e^{-2\beta(t+s)}) \\ &\quad - 2t + \frac{2}{\beta} (1 - e^{-\beta t}) + \frac{2}{\beta} (1 - e^{-\beta t}) e^{-\beta s} - \frac{2}{2\beta} (1 - e^{-2\beta t}) e^{-\beta s} \\ &\stackrel{t \rightarrow \infty}{=} s - \frac{1}{\beta} (1 - e^{-\beta s}), \end{aligned}$$

where we have assumed $s > 0$. For $s < 0$, it is clear that we must replace s with $|s|$. The final result is

$$Y(s) = \lim_{t \rightarrow \infty} \langle z(t+s) z^*(t) \rangle = |z(0)|^2 \exp \left\{ i\nu s - \frac{1}{2} \lambda^2 |s| + \frac{\lambda^2}{2\beta} (1 - e^{-\beta|s|}) \right\}.$$

(3) For the colored noise example in §3.5.3 of the notes, compute numerically $\hat{Y}(\omega)$ and plot your results as a function of $\omega - \nu$. Set $\lambda \equiv 1$ and plot your results for a representative set of different values of the parameter β .

Solution:

We may derive an expansion for $\hat{Y}(\omega)$ as follows. First, for convenience we set $|z(0)|^2 = 1$. Then we have

$$\begin{aligned} Y(s) &= \exp \left\{ i\nu s - \frac{1}{2}\lambda^2 |s| + \frac{\lambda^2}{2\beta} (1 - e^{-\beta|s|}) \right\} \\ &= e^{i\nu s} e^{-\lambda^2|s|/2} e^{\lambda^2/2\beta} \sum_{n=0}^{\infty} \frac{1}{n!} \left(-\frac{\lambda^2}{2\beta} \right)^n e^{-n\beta|s|} \end{aligned}$$

Taking the Fourier transform, we have

$$\hat{Y}(\omega) = e^{\lambda^2/2\beta} \sum_{n=0}^{\infty} \frac{1}{n!} \left(-\frac{\lambda^2}{2\beta} \right)^n \frac{2(n\beta + \frac{1}{2}\lambda^2)}{(\omega - \nu)^2 + (n\beta + \frac{1}{2}\lambda^2)^2}.$$

Define the parameter $\varepsilon \equiv \lambda^2/2\beta$, and define rescaled frequencies $\bar{\omega} \equiv \omega/\beta$ and $\bar{\nu} \equiv \nu/\beta$. Then $\hat{Y}(\omega) = \beta^{-1} \hat{\mathcal{Y}}_{\varepsilon}(\delta)$, where $\delta = \bar{\omega} - \bar{\nu}$ and

$$\begin{aligned} \hat{\mathcal{Y}}_{\varepsilon}(\delta) &= 2 \exp(\varepsilon) \sum_{n=0}^{\infty} \frac{(-\varepsilon)^n}{n!} \frac{n + \varepsilon}{\delta^2 + (n + \varepsilon)^2} \\ &= 2 \int_0^{\infty} d\tau \cos(\delta\tau) \exp \left\{ -\varepsilon (e^{-\tau} - 1 + \tau) \right\}. \end{aligned}$$

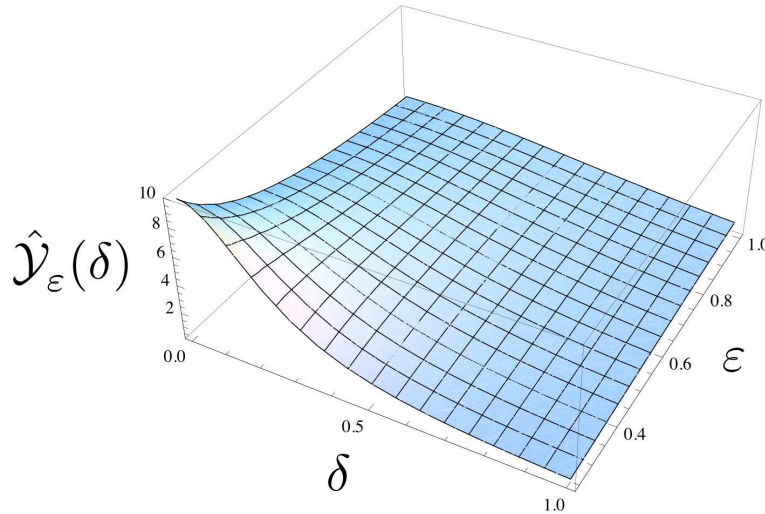


Figure 1: The integral $\hat{\mathcal{Y}}_{\varepsilon}(\delta)$ from problem (3).

Note that $f(\tau) = e^{-\tau} - 1 + \tau$ is nonnegative and monotonically increasing for $\tau \geq 0$, with $f(0) = 0$. For $\varepsilon \ll 1$, we can expand $f(\tau) = \frac{1}{2}\tau^2 + \mathcal{O}(\tau^3)$ and obtain

$$\hat{\mathcal{Y}}_\varepsilon(\delta) \simeq \sqrt{\frac{2\pi}{\varepsilon}} e^{-\delta^2/2\varepsilon} \quad (\varepsilon \rightarrow 0).$$

We evaluate numerically via Mathematica, *viz.*

```
Y[x_, a_] := NIntegrate[2 Cos[x * y] Exp[-a (Exp[-y] - 1 + y)], {y, 0, Infinity}]
Plot3D[Y[x, a], {x, 0, 1}, {a, 0.25, 1}, PlotRange -> Full]
```

The resulting plot is shown in Fig. 1.