Physics 141 Lecture 1 and Lab 1 David Moore, 1/8/2018

## Contents

1	One Particle Central Force		
	1.1	Acceleration Vector	1
	1.2	Numerics	2

## 1 One Particle Central Force

## 1.1 Acceleration Vector

Suppose a single particle of mass m is at position  $\vec{x}$ , and is attracted gravitationally to a single fixed particle of mass M at position  $\vec{x}_0$ . There is a force vector on  $\vec{x}$ , denoted  $\vec{a}$ .



The magnitude of acceleration is  $\|\vec{a}\| = \frac{GM}{\|\vec{x}_0 - \vec{x}\|^2}$  The direction of acceleration is from  $\vec{x}$ , to  $\vec{x}_0$ , so it points in the direction of  $\vec{x}_0 - \vec{x}$ . One way to remember the sign is to think of putting  $\vec{x}$  at the origin. We have the magnitude and direction of the acceleration, so we can find the vector itself:

$$\vec{a} = GM \frac{\vec{x}_0 - \vec{x}}{\|\vec{x}_0 - \vec{x}\|^3}$$

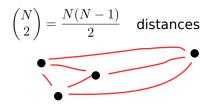
The quantity GM has units of length cubed per time squared, and it's the quantity that is easily experimentally measurable. G is known to four or five digits, but  $GM_{\odot}$  (G times the mass of the sun) is known to ten or eleven digits! GM is called the standard gravitational parameter for a given body.

As an aside, how do you do the same process if instead of one free body you have N free bodies interacting with each other? Sum the forces!

$$\vec{a_i} = \sum_{j=1, j \neq i}^{N} GM_j \frac{\vec{x_j} - \vec{x_i}}{\|\vec{x_j} - \vec{x_i}\|^3}$$

The distance from particle *i* to particle *j* is the same as the distance from *j* to *i*, and we ignore the j = i case, so there are actually only (*N* choose 2) = N(N-1)/2 calculations. Try plugging in  $N = 10^{10}$ , the particle number in some modern simulations. Not even modern supercomputers can run the brute force method<sup>1</sup>!

 $<sup>^{1}</sup>$ A back of the envelope estimate tells me, with a few tens of teraflops of computing power, it would take months to half a year to do finish all of these distance calculations = one timestep. The thing is you need thousands of timesteps to get a useful result.



## 1.2 Numerics

Let's run a simulation of our one body orbiting around a sun using **Euler's method**. I'm going to omit vector arrows, but keep in mind that x, v, and a(x) are all vectors here. We want to solve the second order ODE  $\ddot{x} = a(x)$ . This is equivalent to solving the two first order ODEs  $\dot{x} = v$ ,  $\dot{v} = a(x)$ . We have initial conditions  $v_0$  and  $x_0$  at t = 0. Let's use a fixed timestep  $\delta t$ , and let  $x_n$  denote the position of the particle at time  $n\delta t$ .

One way to approximate the solution is using the Euler scheme:

$$\dot{x} \approx \frac{x_{n+1} - x_n}{\delta t} = v_n$$
  $\dot{v} \approx \frac{v_{n+1} - v_n}{\delta t} = a(x_n)$ 

 $\mathrm{Then.}\,.$ 

$$x_{n+1} = x_n + \delta t \cdot v_n \qquad \qquad v_{n+1} = v_n + \delta t \cdot a(x_n)$$

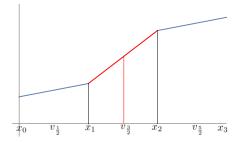
As we saw in the javascript applet in class, this is a pretty bad method. We should get an elliptical orbit, but the result isn't elliptical at all!!



For a function f(t), the quantity  $\frac{1}{\delta t}(f(t+\delta t) - f(t))$  is a pretty bad approximation to the derivative at t, and the error is proportional to  $\delta t$ . You can prove this by using Taylor expansions:

$$\frac{1}{\delta t}(f(t+\delta t) - f(t)) = f'(t) + O(\delta t)$$

There has to be a better way, and the better way is the **leapfrog method**. We can get something accurate to second order in  $\delta t$  if we use a bit of intuition: when we find the slope of a line segment, we're really finding a good approximation to the slope at the midpoint of the line segment.



You can prove the following formula if you Taylor expand both sides about  $t\colon$ 

$$\frac{f(t+\delta t) - f(t)}{\delta t} = f'\left(t + \frac{\delta t}{2}\right) + O(\delta t^2)$$

Then we have the following approximations:

$$\dot{x} \approx \frac{x_{n+1} - x_n}{\delta t} = v_{n+1/2}$$
  $\dot{v} \approx \frac{v_{n+1/2} - v_{n-1/2}}{\delta t} = a(x_n)$ 

which we turn into the Leapfrog numerical scheme:

$$v_{n+1/2} = v_{n-1/2} + \delta t \cdot a(x_n) \qquad \qquad x_{n+1} = x_n + \delta t \cdot v_{n+1/2}$$

Leapfrog is a great algorithm for four reasons:

- 1. It's second-order (whereas Euler is only first-order),
- 2. It only uses one force evaluation per timestep (force evaluations are expensive!),
- 3. It is time-reversible (unlike Euler),
- 4. It's *symplectic*, meaning it conserves phase space volume.

The power of the method is demonstrated by the Javascript implementation we went over class.

