

Physics 4A  
Lecture 8: Feb. 17, 2015  
Sunil Sinha  
UCSD Physics

What's In Common ?



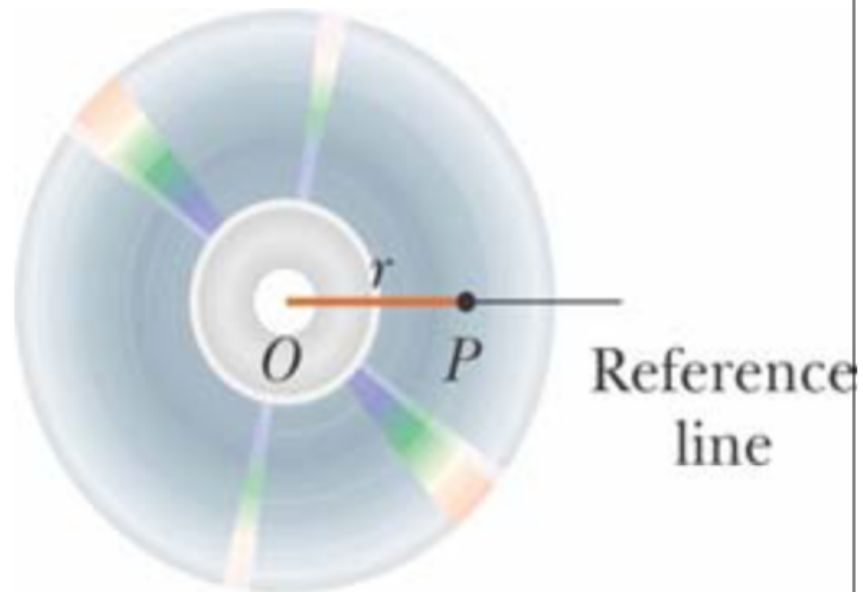
Each involves a body that rotates about a fixed axis

# Rotation of Rigid Bodies

- Rigid body is one that is non-deformable
  - relative location of all particles making up the object remain **constant**
- All real objects are deformable so “rigid body” is an *idealized model* but a useful one
- **This week:**
  - kinematic language to describe rotation
    - Like chapter 2, derive equation of motion
  - **Kinetic Energy in rotation**
  - **Dynamics of rotation**

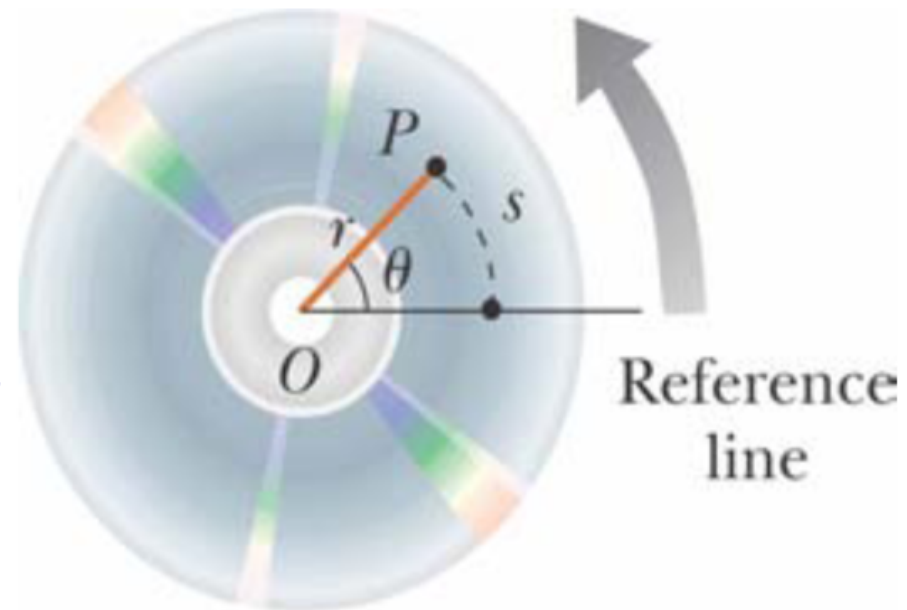
# Angular Position In Rotation

- Axis of rotation passes thru center of disc & is  $\perp$  to plane of picture
- Choose a fixed reference line
- Point  $P$  is at a **fixed** distance  $r$  from the origin
- Point  $P$  will rotate about the origin in a circle of radius  $r$



# Change In Angular Position

- As the particle moves, the only coordinate that changes is  $\theta$
- As the particle moves through  $\theta$ , it moves through an **arc length  $s$**
- The arc length and  $r$  are related:  
–  $s = r \theta$



# Angular Displacement

$$\text{Angular displacement } \theta = \frac{s}{r}$$



$\theta$  is a pure number, but

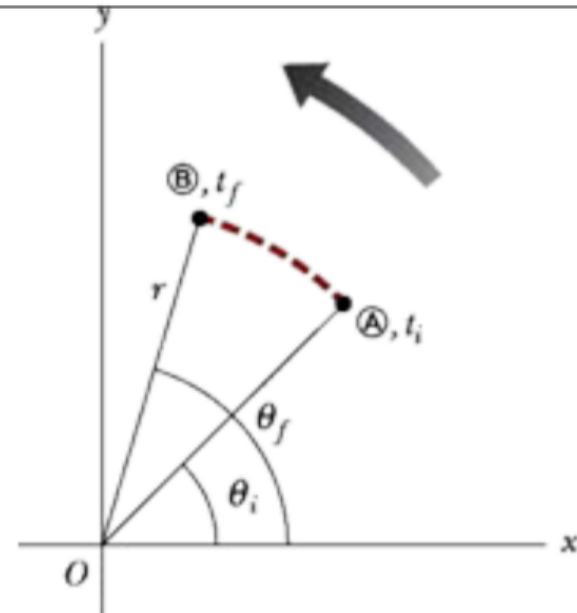
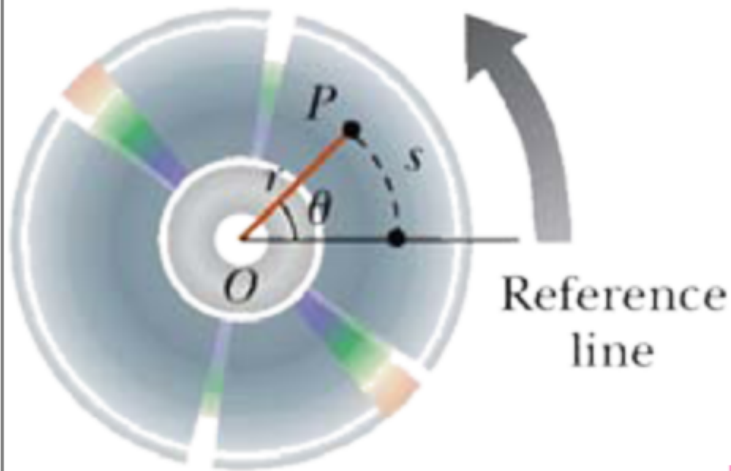
is given the artificial unit, **radian**

One radian is the angle subtended by an arc length equal to the radius of the arc

Comparing Degree and Radians:

$$1 \text{ rad} = \frac{360^\circ}{2\pi \text{ rad}} \approx 57.3^\circ$$

# Angular Velocity $\omega$



$$\omega_{av-z} = \frac{\theta_f - \theta_i}{t_f - t_i} = \frac{\Delta\theta}{\Delta t}$$

$$\omega_z = \lim_{\Delta t \rightarrow 0} \frac{\Delta\theta}{\Delta t} = \frac{d\theta}{dt}$$

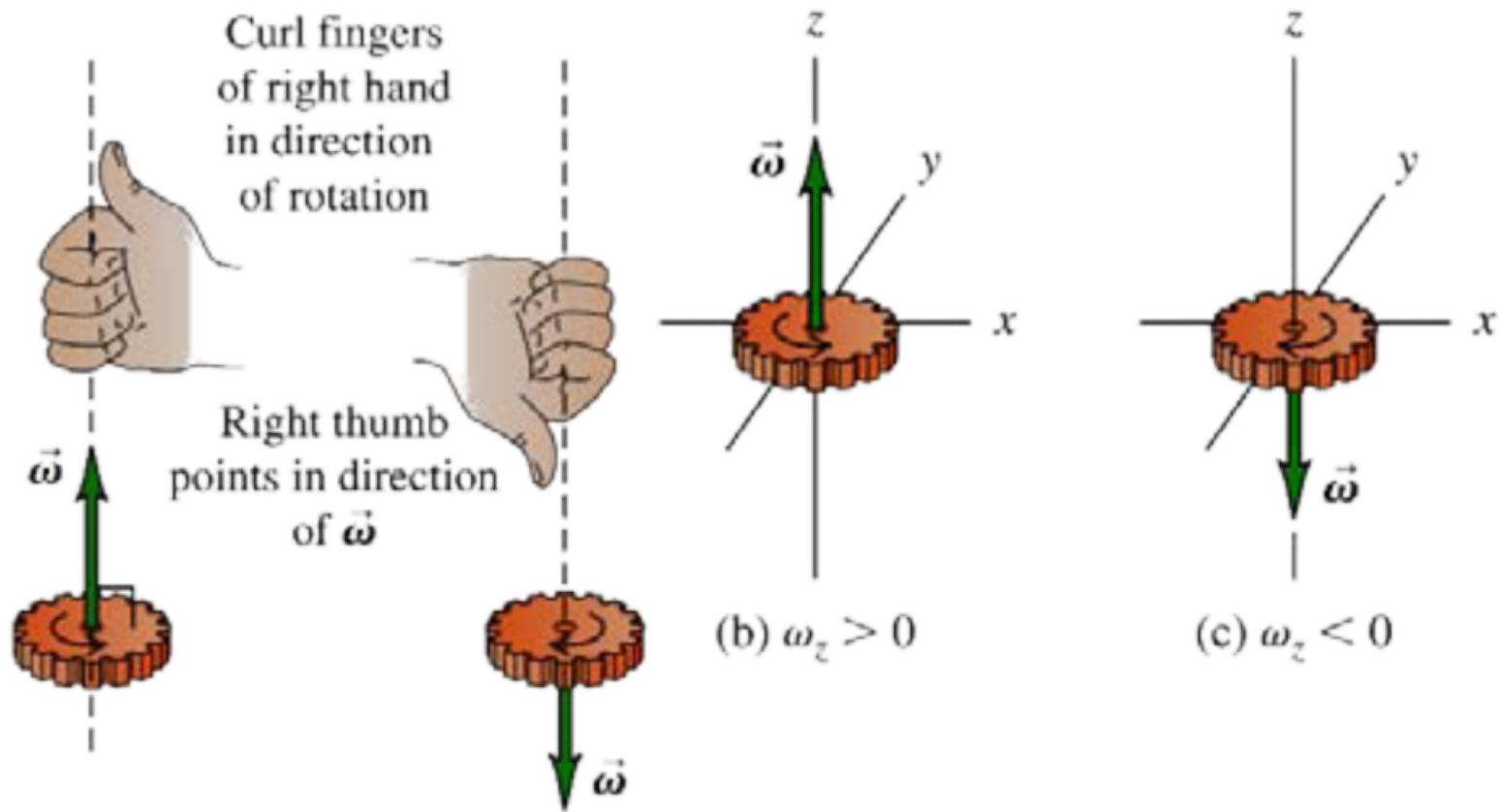
Unit of  $\omega$ : rad/s

rpm : revs/second

$$1 \text{ rpm} = \frac{2\pi}{60} \text{ rad/s}$$

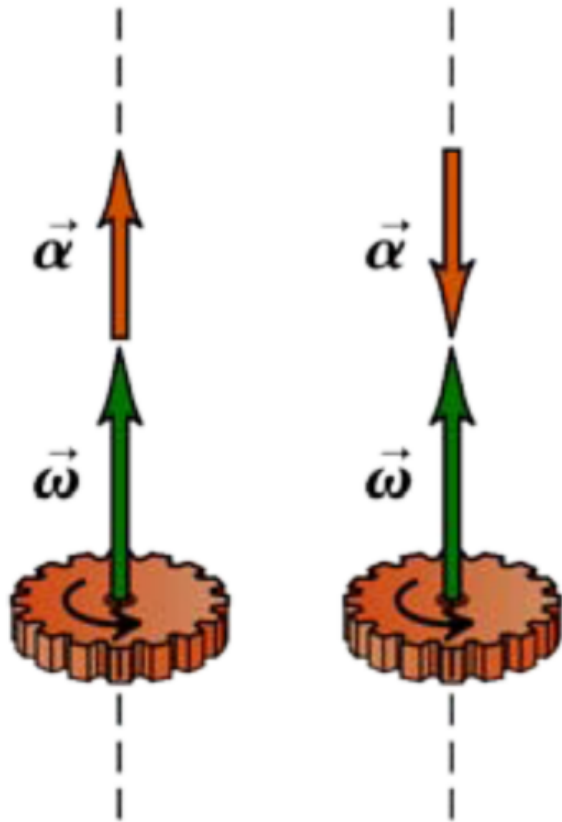
$$1 \text{ rad/s} \approx 10 \text{ rpm}$$

# Angular Velocity $\vec{\omega}$ & Right Hand Rule





# Angular Acceleration $\vec{\alpha}$



Speeding up    Slowing up

$$\alpha_{av-z} = \frac{\omega_2 - \omega_1}{t_2 - t_1} = \frac{\Delta\omega}{\Delta t}$$

$$\alpha_z = \lim_{\Delta t \rightarrow 0} \frac{\Delta\omega}{\Delta t} = \frac{d\omega_z}{dt} = \frac{d^2\theta}{dt^2}$$

$$\vec{\alpha}_z > 0 \Rightarrow \vec{\omega} \text{ increasing}$$

## Eqns. of Motion: Rotation with constant $\vec{\alpha}$

At  $t = 0$ , body at  $\theta_0$  moves with ang. vel.  $= \vec{\omega}_{0z}$  &  $\vec{\alpha}_z = \text{const}$

at time  $t=t$ , body at  $\theta$  moves with ang. vel.  $= \vec{\omega}_z$  &  $\vec{\alpha}_z = \text{const}$

$$\text{By def: } \alpha_z = \frac{\omega_z - \omega_{0z}}{t - 0} \Rightarrow \omega_z = \omega_{0z} + \alpha_z t$$

$$\text{By def: } \omega_{\text{av-z}} = \frac{\omega_{0z} + \omega_z}{2} = \frac{\theta - \theta_0}{t - 0} \Rightarrow \theta - \theta_0 = \frac{\omega_{0z} + \omega_z}{2} t$$

$$\begin{aligned} \text{substitute for } \omega_z &\Rightarrow \theta - \theta_0 = (1/2)[\omega_{0z} + \omega_{0z} + \alpha_z t]t \\ &\Rightarrow \theta - \theta_0 = \omega_{0z} t + (1/2)\alpha_z t^2 \quad \dots\dots \text{for constant } \alpha_z \end{aligned}$$

$$\text{Substitute } t = (\omega_z - \omega_{0z}) / \alpha_z \Rightarrow \omega_z^2 = \omega_{0z}^2 + 2\alpha_z(\theta - \theta_0)$$

## Translational & Rotation Motion Compared

Straight-Line Motion with Constant Linear Acceleration	Fixed-Axis Rotation with Constant Angular Acceleration
---	---

$$a_x = \text{constant}$$

$$v_x = v_{0x} + a_x t$$

$$x = x_0 + v_{0x} t + \frac{1}{2} a_x t^2$$

$$v_x^2 = v_{0x}^2 + 2a_x(x - x_0)$$

$$x - x_0 = \frac{1}{2}(v_x + v_{0x})t$$

$$\alpha_z = \text{constant}$$

$$\omega_z = \omega_{0z} + \alpha_z t$$

$$\theta = \theta_0 + \omega_{0z} t + \frac{1}{2} \alpha_z t^2$$

$$\omega_z^2 = \omega_{0z}^2 + 2\alpha_z(\theta - \theta_0)$$

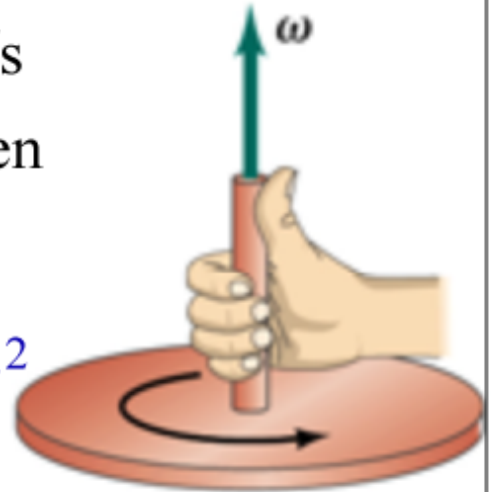
$$\theta - \theta_0 = \frac{1}{2}(\omega_z + \omega_{0z})t$$

A wheel rotates with constant angular acceleration of  $3.50 \text{ rad/s}^2$ . If the angular speed of wheel is  $2.00 \text{ rad/s}$  at  $t = 0$ , thru what angle does the wheel rotate between  $t = 0$  and  $t = 2\text{s}$  ? what's wheel's ang. speed after  $2 \text{ s}$ ?

$$\theta = \theta_0 + \omega_{0z}t + \frac{1}{2}\alpha t^2 \Rightarrow \theta - \theta_0 = \omega_{0z}t + \frac{1}{2}\alpha t^2$$

$$\Delta\theta = (2.00 \text{ rad/s})(2.00\text{s}) + \frac{1}{2}(3.50 \text{ rad/s}^2)(2.00\text{s})^2$$

$$= 11.0 \text{ rad} \equiv \frac{11 \text{ rad}}{2\pi \text{ rad/rev}} = 1.75 \text{ rev}$$

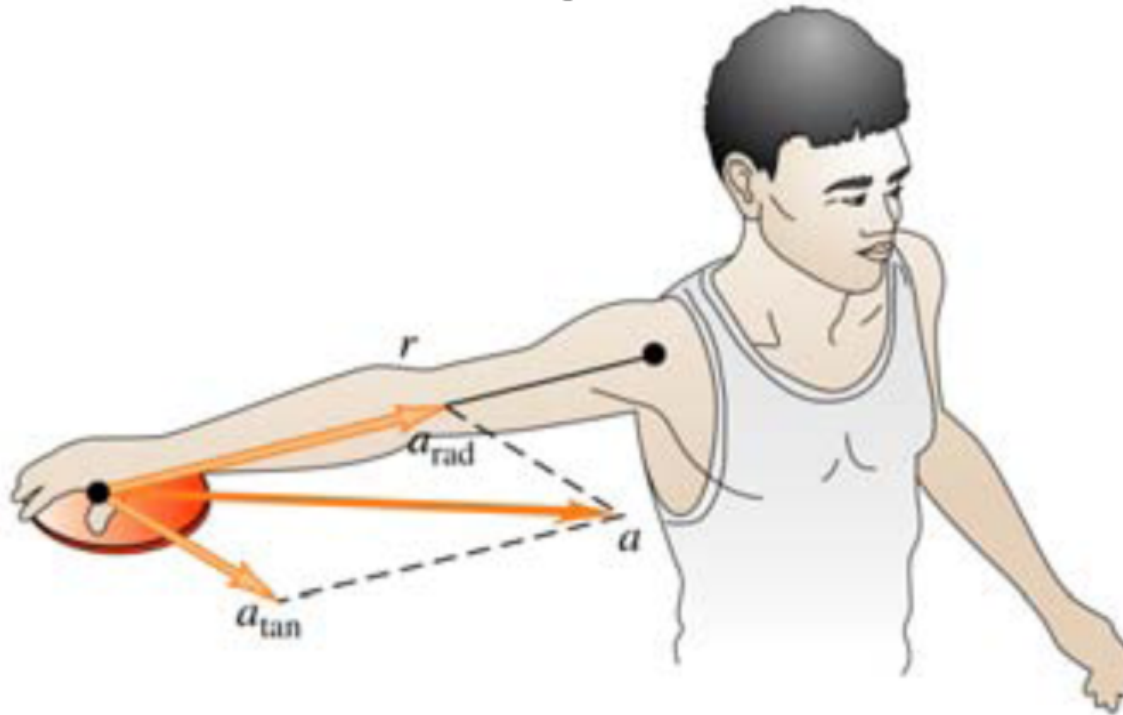


After 2 sec, ang. speed  $\omega_z = \omega_{0z} + \alpha t$

$$= 2.00 \text{ rad/s} + (3.50 \text{ rad/s}^2)(2.00\text{s}) = 9.00 \text{ rad/s}$$

# Relating Rotational & Linear Motion

Man moves discuss in circle of radius 80.0cm. At some instant when he is spinning at angular speed of 10.0rad/s, with the speed increasing at 50.0rad/s<sup>2</sup>. What is tangential & radial component of acceleration of discuss & magnitude of acceleration



## Relating Linear & Angular Kinetic Variables

consider a point P on rotating body

$$\text{since } \mathbf{s} = r\hat{e} \Rightarrow \left| \frac{ds}{dt} \right| = r \left| \frac{d\theta}{dt} \right| \Rightarrow v = r\omega$$

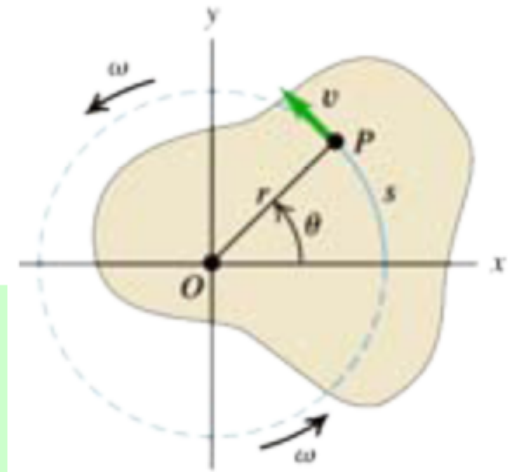
Tangential component of acceleration

$\vec{a}_{\text{tan}} \parallel \vec{v}$  changes magnitude of particle's

linear speed  $|\vec{v}|$  : 
$$\mathbf{a}_{\text{tan}} = \frac{dv}{dt} = r \frac{d\omega}{dt} = r\alpha$$

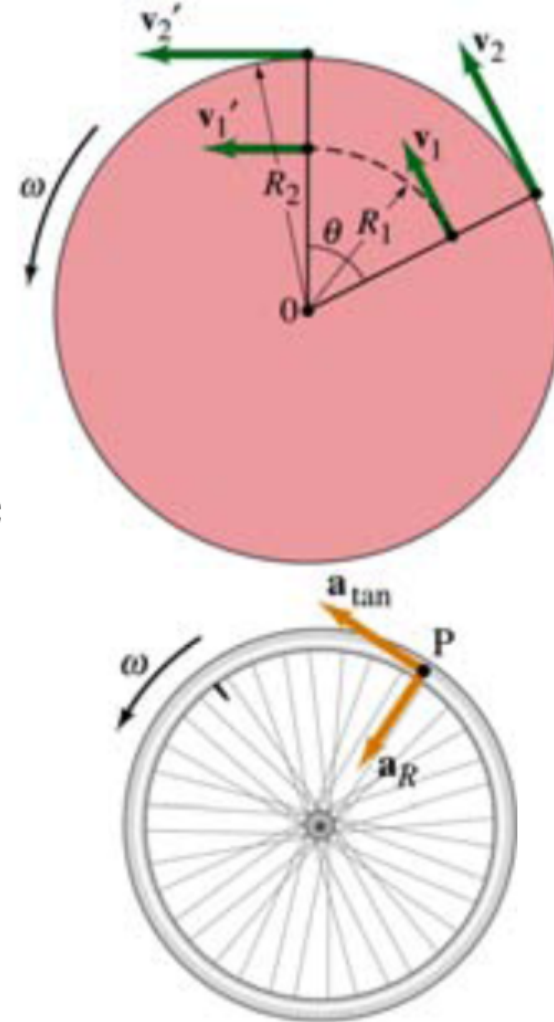
$\mathbf{a}_{\text{rad}}$  = centripetal component of body's acceleration  
related to the change in direction of linear velocity

$$\mathbf{a}_{\text{rad}} = \frac{v^2}{r} = \frac{\omega^2 r^2}{r} = \omega^2 r \quad (\text{directed radially in})$$



## Speed and Acceleration: Watch Out !

- All points on the rigid object will have the same *angular speed*, but **not** the same *tangential speed*
- All points on the rigid object will have the same *angular acceleration*, but **not** the same *tangential acceleration*
- The tangential quantities depend on  $r$ , and  $r$  is not the same for all points on the object



# The Joy of Rotation !





# Rotational Kinetic Energy

Rigid body = **collection of particles** of mass  $m_i$ , speed  $v_i$

Kin. energy of  $i$ th particle  $K_i = (1/2)m_i v_i^2 = (1/2)m_i r_i^2 \omega^2$

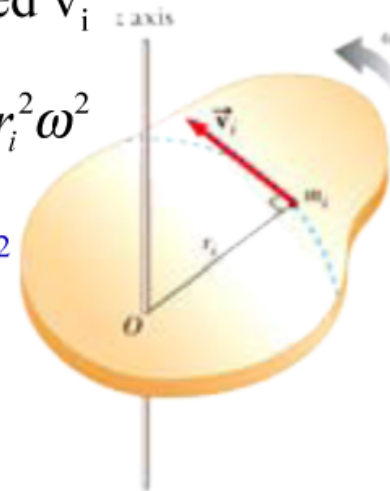
Total kin. energy  $K = \sum_i K_i = \frac{1}{2} (\sum_i m_i r_i^2) \omega^2 = \frac{1}{2} I \omega^2$

$I$  = body's moment of Inertia for *this* rotation axis

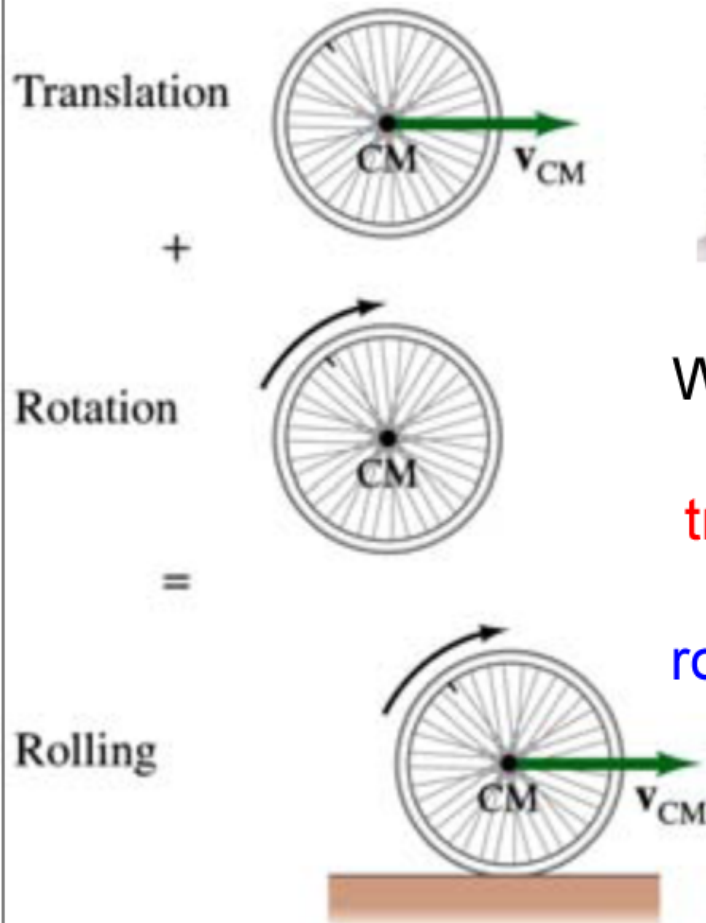
Moment of Inertia is measure of object's resistance to change in its angular speed.  $\Rightarrow$  **rotational inertia**

$$I = m_1 r_1^2 + m_2 r_2^2 + m_3 r_3^2 \dots = \sum_i m_i r_i^2$$

depends not just on the total mass but how it is **arranged** *around* the rotation axis



# Rolling Motion Of A Ball or Wheel



Wheel rolling without slipping  
can be considered as  
translation of the wheel as a  
whole with velocity  $v_{CM}$  +  
rotation about center of mass

# Rolling Without Slipping

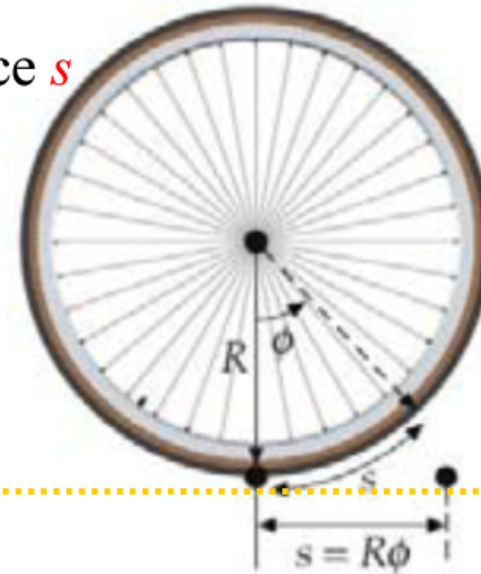
As a wheel rotates thru  $\phi$  without slipping, point of contact between wheel and surface moves distance  $s$

$$\text{Then } S = R\phi$$

If wheel rolling on flat surface then wheel's CM remains directly over point of contact, so it too moves a distance  $s = R\phi$

$$\Rightarrow v_{CM} = R\omega$$

$$\text{and } a_{CM} = R\alpha$$



# Pop Quiz

- A disc of radius 2 m is rolling on level ground without slipping. Its CM is moving horizontally at a velocity of 4 m/s. The instantaneous velocity of the top of the disc is
  - (A) 8 m/s
  - (B) 4 m/s
  - (C)  $-4$  m/s
  - (D) 0 m/s

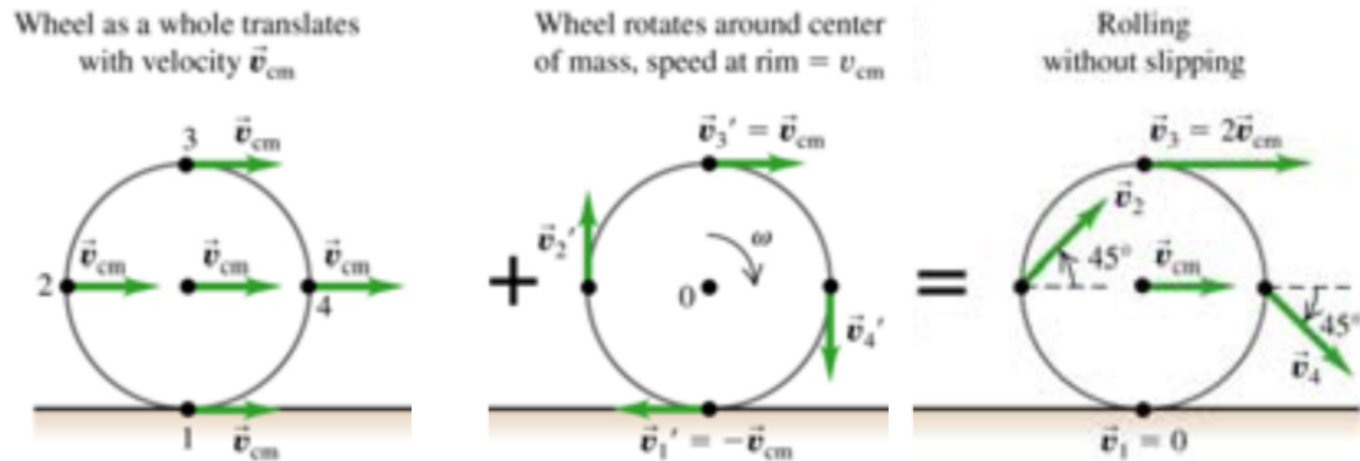
# Wheel (Radius $R$ , Ang. speed $\omega$ ) Rolls Without Slipping

Symmetric Wheel  $\Rightarrow$  CM = Geometric center

Imagine observer at rest w.r.t surface on which a wheel rolls

In Observer's frame: point on wheel touching surface must

**instantaneously be at rest** so that it does not **slip**



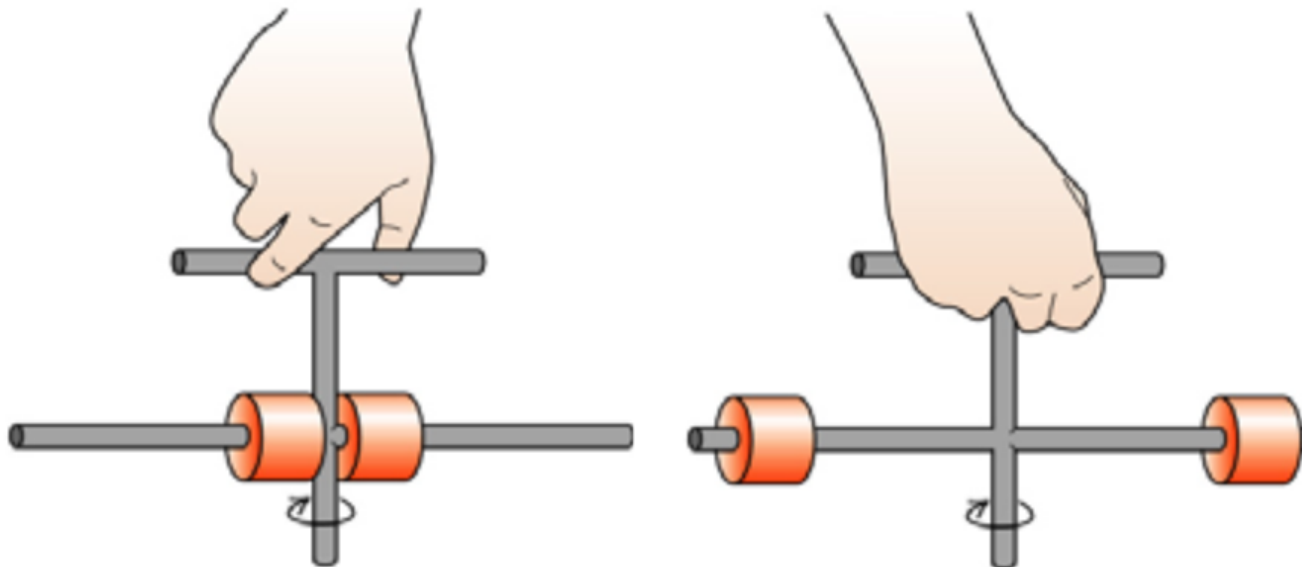
Velocity of point of contact  $\vec{v}_1' = R\vec{\omega}$  (rel. to CM) must have same magnitude but opposite direction as  $\vec{v}_{cm} \Rightarrow \boxed{\vec{v}_{cm} = -R\vec{\omega}}$

## Same Total Mass, Different Moment of Inertia

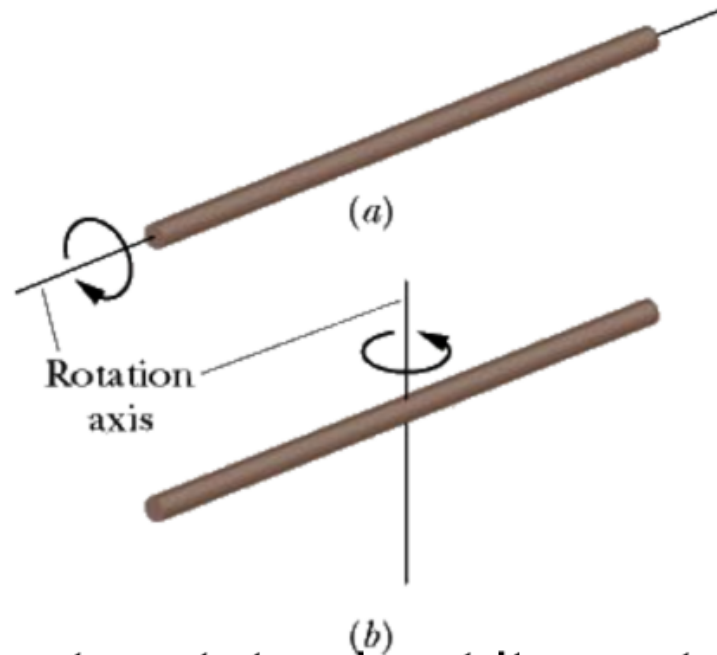
$$I = m_1 r_1^2 + m_2 r_2^2$$

- Mass close to axis
- Small moment of inertia
- Easy to start apparatus rotating

- Mass farther from axis
- Greater moment of inertia
- Harder to start apparatus rotating



## Diff. Rotation Axis $\Rightarrow$ Diff. MOI for Same Body



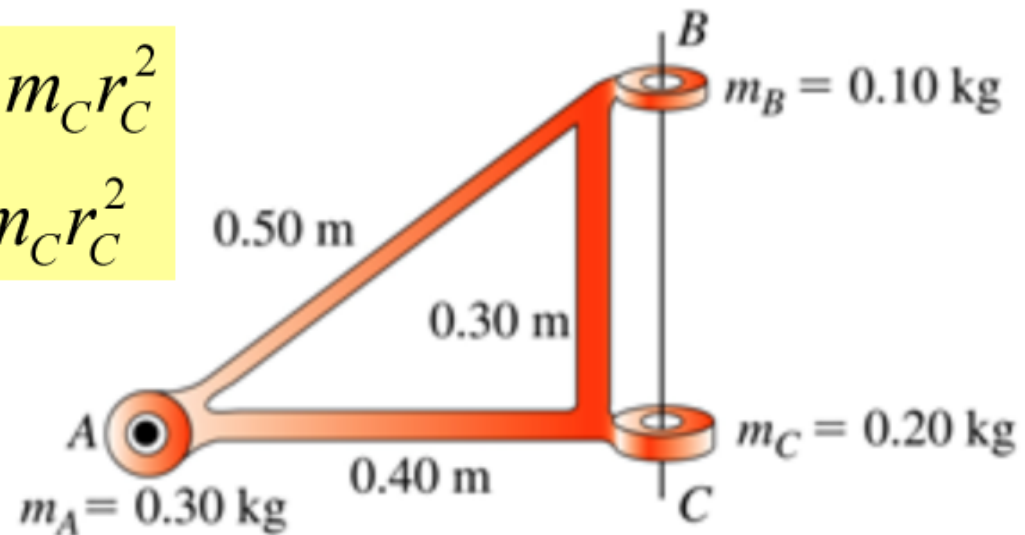
Long rod easier to rotate about its central axis  
(longitudinal) axis (a) than about an axis (b) through its  
center and  $\perp$  to its length

Because mass is distributed closer to  
longitudinal axis (a) than (b)

## Diff. Rotation Axis $\Rightarrow$ Diff. MOI for Same Body

3 heavy connectors linked by “massless” plastic strut

$$\begin{aligned} I_A &= m_A r_A^2 + m_B r_B^2 + m_C r_C^2 \\ &= 0 + m_B r_B^2 + m_C r_C^2 \end{aligned}$$



$$\begin{aligned} I_B &= m_A r_A'^2 + m_B r_B'^2 + m_C r_C'^2 \\ &= m_A r_A'^2 + 0 + 0 \end{aligned}$$

$I_A \neq I_B$ . Since  $K = \frac{1}{2} I \omega^2 \Rightarrow K_A \neq K_B$



## MOI For A Continuous Distribution Of Mass

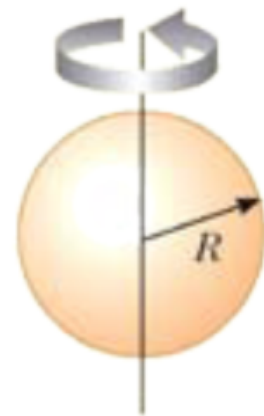
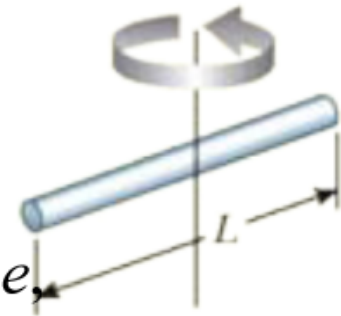
$$I = \lim_{\Delta m_i \rightarrow 0} \sum_i r_i^2 \Delta m_i = \int r^2 dm$$

For objects of uniform density  $\rho = \text{mass} / \text{volume}$ ,  
easiest to express  $dm = \rho dV$

$$I = \int r^2 dm = \int r^2 \rho dV = \rho \int r^2 dV$$

Volume element  $dV = dx dy dz$

Limits of integration defined by  
shape & dimension of the object

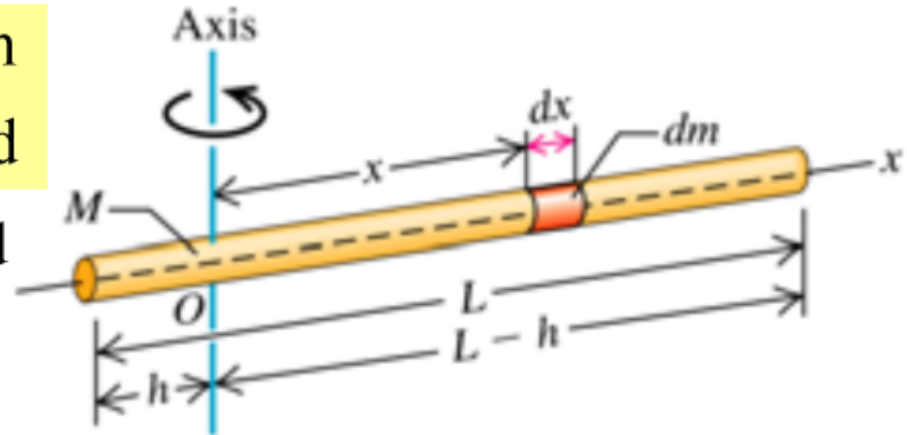


# Moment Of Inertia For Uniform Thin Rod

Calculate I about axis of rotation thru O, distance h from rod's end

Mass uniformly distributed on rod

$$\frac{dm}{M} = \frac{dx}{L} \Rightarrow dm = \frac{M}{L} dx$$



$$I_h = \int x^2 dm = \frac{M}{L} \int_{-h}^{L-h} x^2 dx = \left[ \frac{M}{L} \left( \frac{x^3}{3} \right)_{-h}^{L-h} \right] = \frac{1}{3} M(L^2 - 3Lh - 3h^2)$$

MOI if rod rotating thru left (h=0) or right (h=L) end:

$$I_{0,L} = \frac{1}{3} ML^2$$

## Moment Of Inertia For Thick Hollow Cylinder

Pick as volume element a thin cylindrical shell of radius  $r$   
thickness  $dr$  and length  $L \Rightarrow dV=(2\pi r)Ldr$

$\Rightarrow dm=\rho(2\pi r)Ldr$ . Cylinder thickness =  $R_2 - R_1$

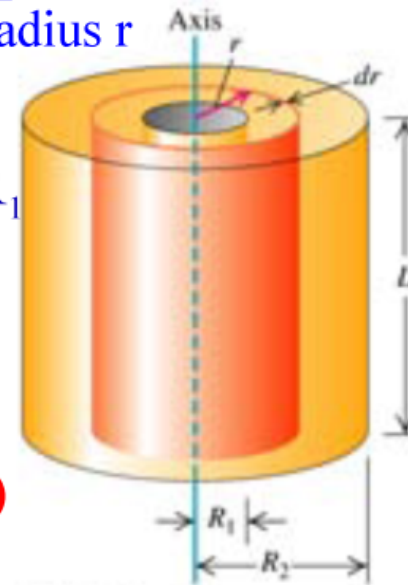
$$I = \int_{R_1}^{R_2} r^2 \rho(2\pi r)Ldr = 2\pi\rho L \int_{R_1}^{R_2} r^3 dr$$

$$= \frac{2\pi\rho L}{4}(R_2^4 - R_1^4) = \frac{2\pi\rho L}{4}(R_2^2 - R_1^2)(R_2^2 + R_1^2)$$

Cylinder volume  $V=\pi L(R_2^2 - R_1^2) \Rightarrow M = \rho\pi L(R_2^2 - R_1^2)$

$$\Rightarrow \boxed{I = \frac{1}{2} M(R_2^2 + R_1^2)}$$

For a solid cylinder  $R_1=0 \Rightarrow I = \frac{1}{2} MR^2$



# Pop Quiz

- A rope is being wound without slipping onto a cylinder of radius 2 m which is rotating at 1 revolution/minute. The speed with which a point on the rope is moving is
  - (A)  $2/60$  m/s
  - (B)  $4\pi/60$  m/s
  - (C)  $60/2\pi$  m/s
  - (D)  $4\pi$  m/s

# Cable Unwinding Off A Cylinder

Cable wrapped around cylinder (mass  $M$ ) is attached to object mass  $m$ . As cable unwinds,  $U_{\text{grav}}$  converted to kinetic energy.

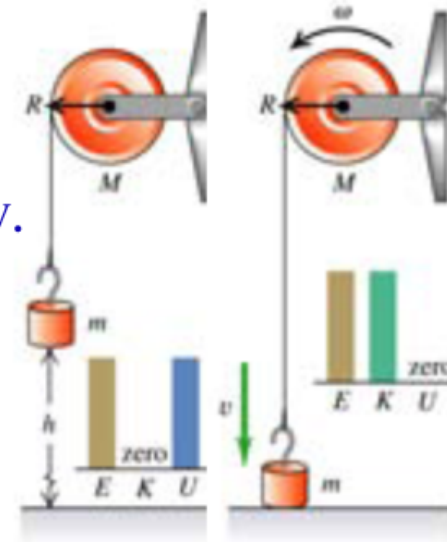
Find speed of object as it hits floor

$$K_1 + U_1 = K_2 + U_2$$

$$0 + mgh = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2 + 0$$

$$\Rightarrow mgh = \frac{1}{2}mv^2 + \frac{1}{2}\left(\frac{1}{2}MR^2\right)\left(\frac{v}{R}\right)^2$$

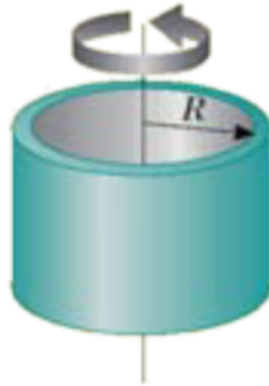
$$\Rightarrow v = \sqrt{\frac{2gh}{1 + M/2m}}$$



Cylinder is Solid

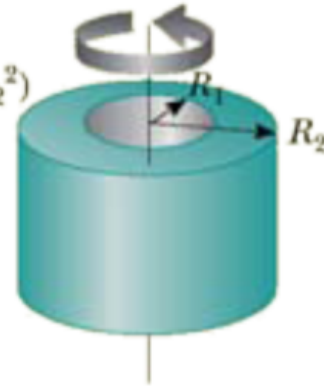
## Moments of Inertia of Homogeneous Rigid Objects With Different Geometries

Hoop or thin  
cylindrical shell  
 $I_{CM} = MR^2$

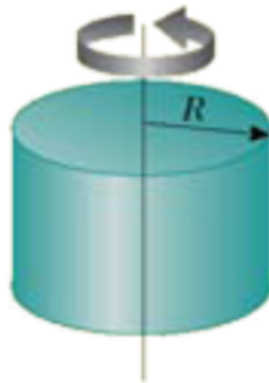


Hollow cylinder

$$I_{CM} = \frac{1}{2} M(R_1^2 + R_2^2)$$

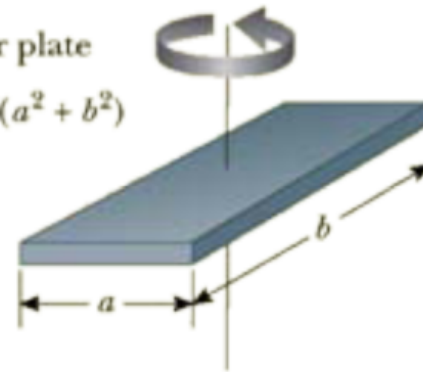


Solid cylinder  
or disk  
 $I_{CM} = \frac{1}{2} MR^2$



Rectangular plate

$$I_{CM} = \frac{1}{12} M(a^2 + b^2)$$



Moment of Inertia of rectangle about normal axis through center

$$I_{strip} = \frac{1}{12} b^2 dm$$

$$dm = \sigma b dy$$

$$I_{strip} = \frac{1}{12} b^2 \sigma b dy$$

$$dI_{cm} = I_{strip} + y^2 dm = \frac{1}{12} b^2 \sigma b dy + y^2 \sigma b dy$$

$$I_{cm} = \sigma b \int_{-a/2}^{a/2} \left[ \frac{1}{12} b^2 + y^2 \right] dy = \sigma b \left[ \frac{1}{12} b^2 a + \frac{1}{12} a^3 \right]$$

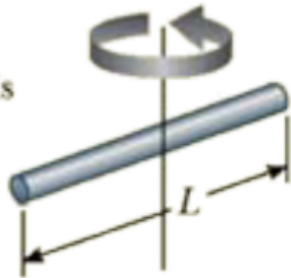
$$= \frac{1}{12} \sigma ab [a^2 + b^2] = \frac{1}{12} M [a^2 + b^2]$$

$$(M = \sigma ab)$$

## Moment of Inertia of Objects About Their Axis Of Rotation

Long thin rod  
with rotation axis  
through center

$$I_{\text{CM}} = \frac{1}{12} ML^2$$



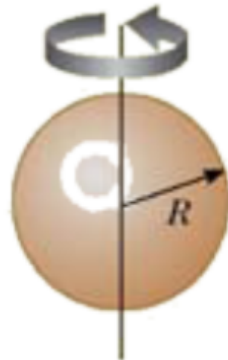
Long thin  
rod with  
rotation axis  
through end

$$I = \frac{1}{3} ML^2$$



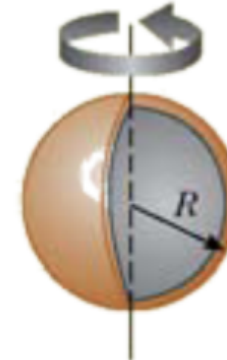
Solid sphere

$$I_{\text{CM}} = \frac{2}{5} MR^2$$



Thin spherical  
shell

$$I_{\text{CM}} = \frac{2}{3} MR^2$$



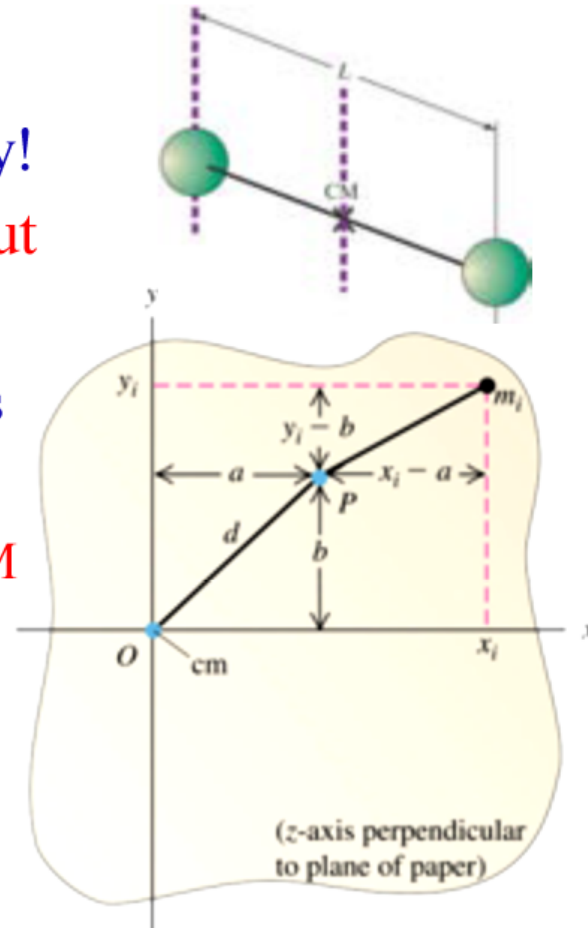


# Parallel Axes Theorem

A body does not have only one moment of Inertia, its has many!  
*As many as the many axes about which it can rotate*

Relation between MOI about an axis of rotation thru its center of mass & MOI about any other axis  $\parallel$  to C.O.M axis but displaced by distance  $d$

$$I_P = I_{cm} + Md^2$$



Moment of inertia of rectangle about in-plane axis going through center of side of length  $a = (1/12) Ma^2$

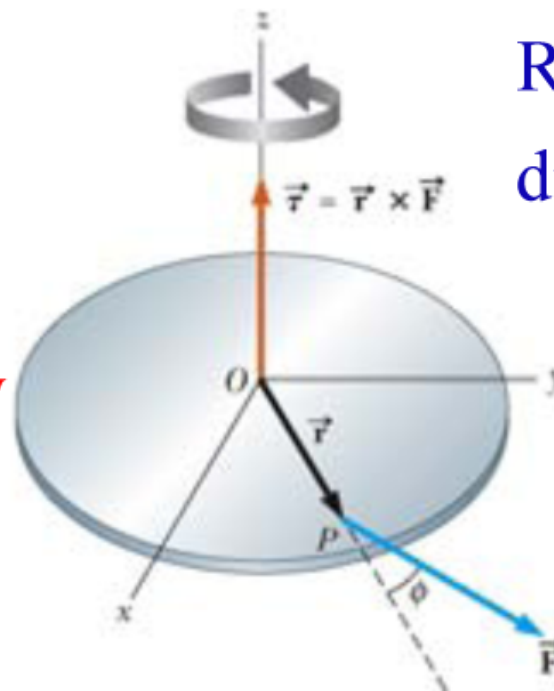
Moment of inertia of rectangle about in-plane axis going through center of side of length  $b = (1/12) Mb^2$

## Torque $\vec{\tau}$ : Rotational Analog Of Force

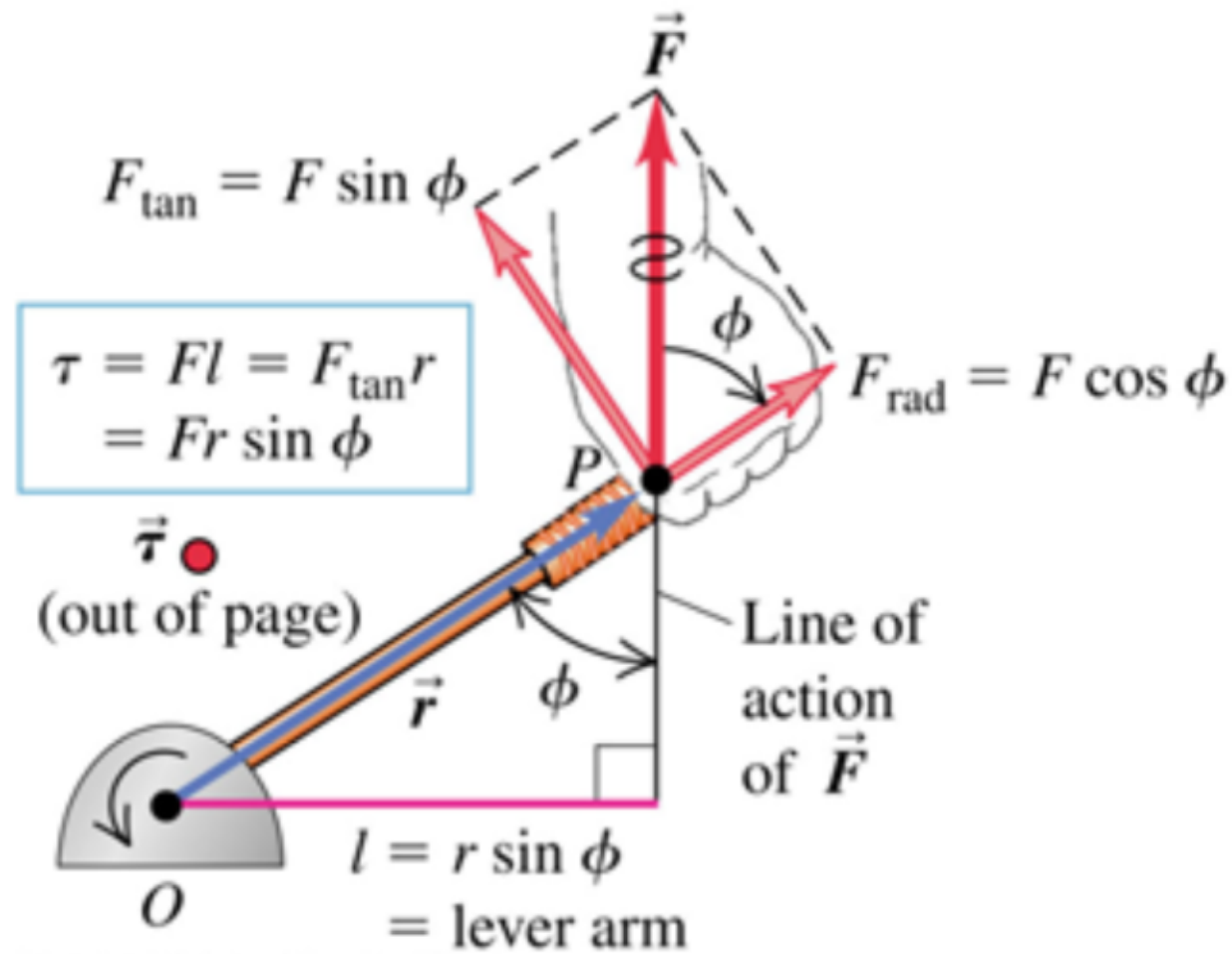


Torque at a point, due to an applied force causes angular acceleration

$\vec{\tau}$  lies in direction  
 $\perp$  to plane formed by  
position vector  $\vec{r}$  &  
applied force  $\vec{F}$



Rotation in  
direction of  $\vec{\tau}$

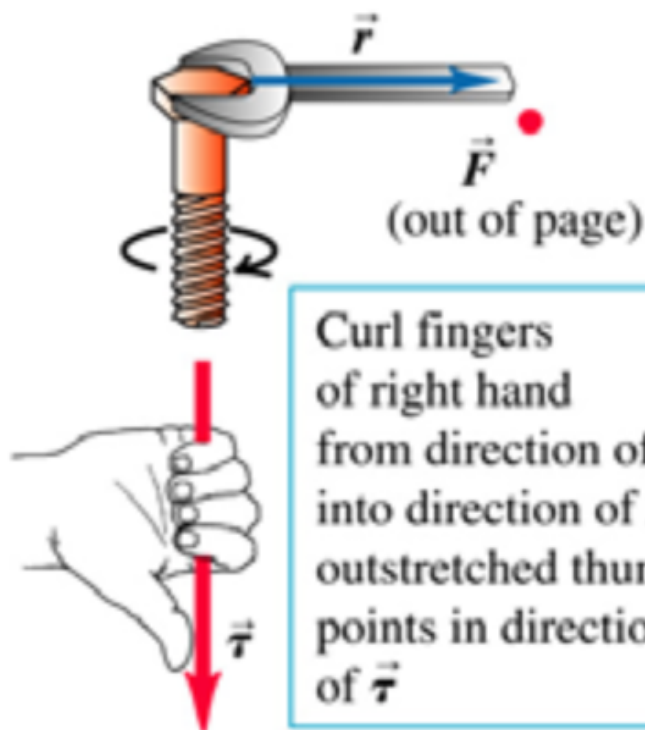


Lever arm =  $\perp$  distance between axis of rotation & line of action of force

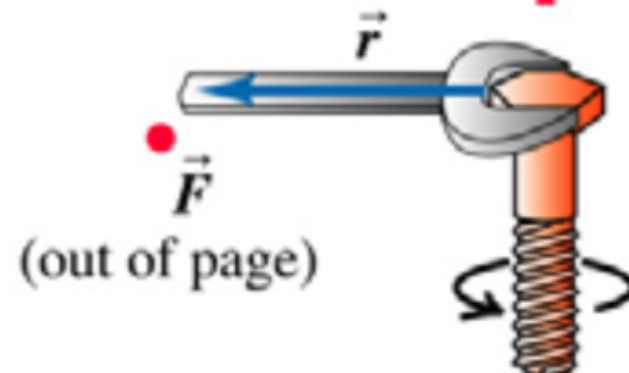
# Direction of The Torque Vector

$$\vec{\tau} = \vec{r} \times \vec{F}$$

$$|\vec{\tau}| = rF\sin\theta$$



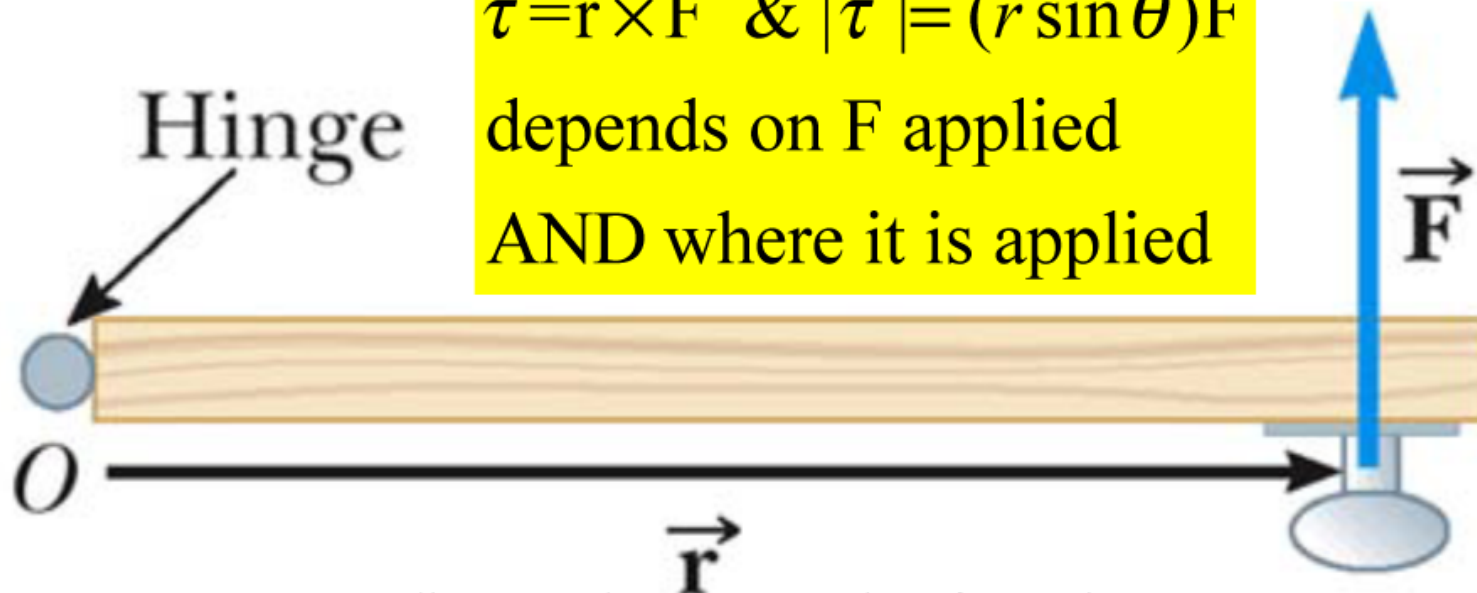
Curl fingers of right hand from direction of  $\vec{r}$  into direction of  $\vec{F}$ ; outstretched thumb points in direction of  $\vec{\tau}$



# Door Knob & The Hinge

Why are ALL doorknobs always located the furthest from the hinge around which they rotate ?

$\vec{\tau} = \vec{r} \times \vec{F}$  &  $|\vec{\tau}| = (r \sin \theta)F$   
depends on F applied  
AND where it is applied



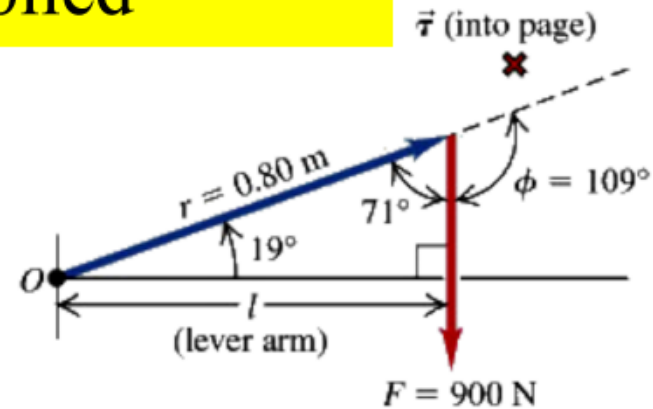
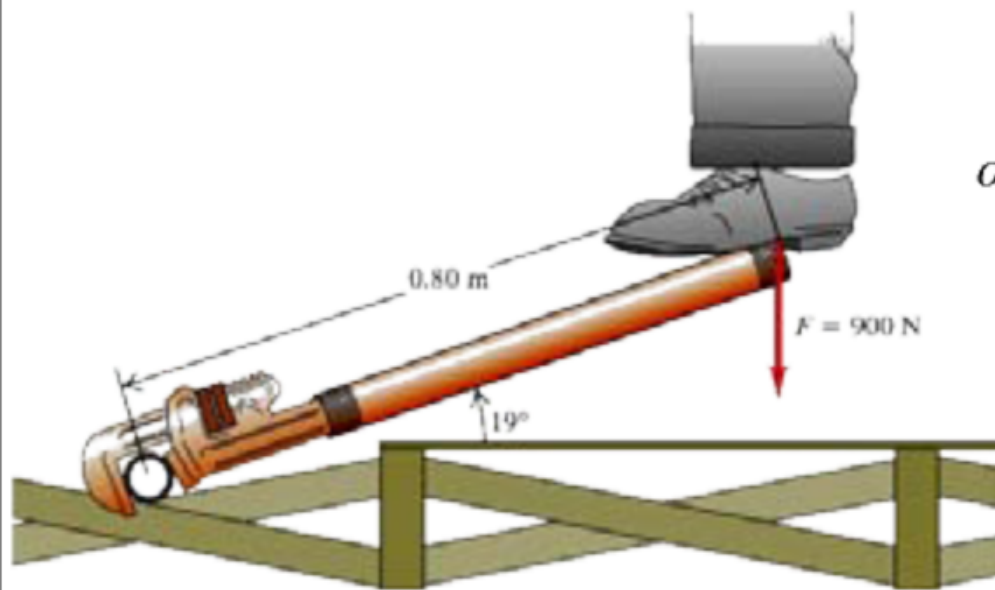
Lever arm =  $\perp$  distance between axis of rotation & line of action of force

Larger the  $r \Rightarrow$  less effort (F) needed

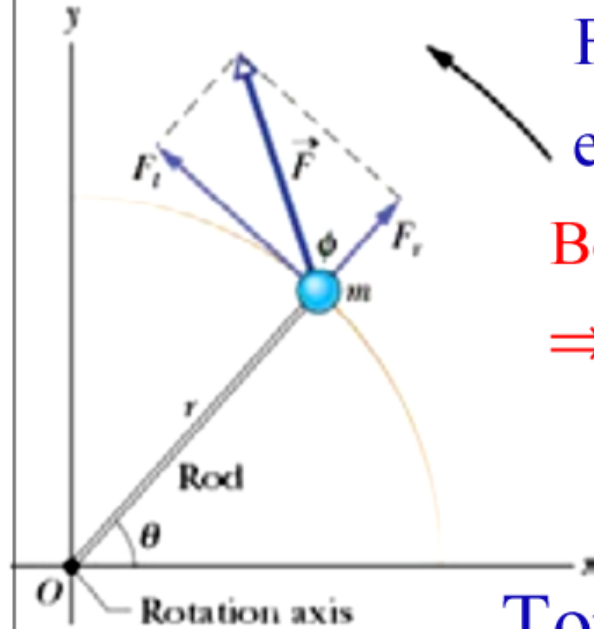
# How to Get *More* Out Of Same

$$\vec{\tau} = \vec{r} \times \vec{F} \quad \& \quad |\vec{\tau}| = (r \sin \theta) F$$

Torque depends on force applied  
AND where it is applied



# Torque & Angular Acceleration



Force  $\vec{F}$  acts on body of mass  $m$  on one end of a massless rod of length  $r$

Body rotates around an axis  $\perp$  to  $x$ - $y$  plane  
 $\Rightarrow$  circular motion in  $x$ - $y$  plane

$$F_{\text{tangential}} = ma_{\text{tangential}}$$

$$\text{Torque } \tau = F_{\text{tangential}} r = ma_{\text{tangential}} r$$

$$\Rightarrow \tau = m(r\alpha)r = (mr^2)\alpha \quad (\alpha \text{ in radians!})$$

$$\text{since } I = mr^2 \Rightarrow \boxed{\vec{\tau} = I\vec{\alpha}}$$

Newton's 2<sup>nd</sup> Law For Rotational Motion



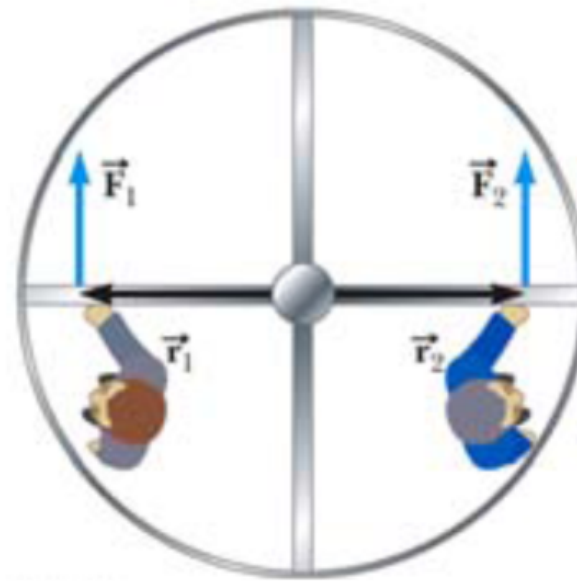
# Battle Of The Revolving Door !

Two fired Trump apprentices are trying to use a revolving door. If  $|r_1| > |r_2|$  &  $|F_1| = |F_2|$ , **which way** will the door turn?

If  $\tau_1 = -r_1 F_1$  then  $\tau_2 = +r_2 F_2$

$$\sum \vec{\tau} = \vec{\tau}_1 + \vec{\tau}_2 = -r_1 F + r_2 F < 0$$

$\Rightarrow$  Net torque is negative,  
will produce a **clockwise**  
rotation with **downward**  
angular acceleration  $\vec{\alpha}$

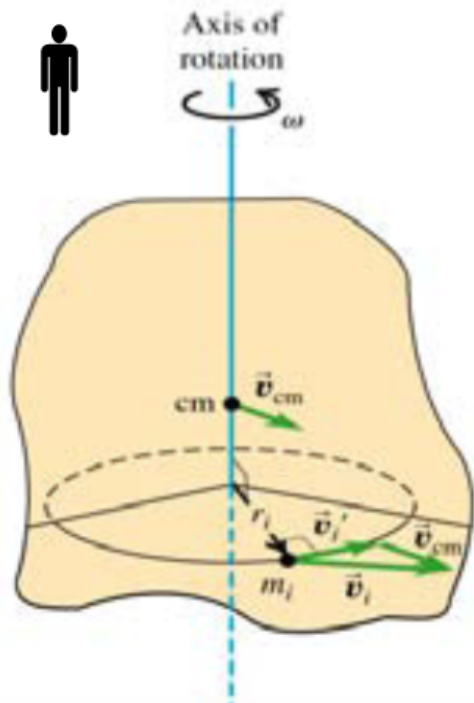


Which will no doubt upset the **blue suit**

## Rigid Body Rotation About a Moving Axis

A rigid body's motion = **sum** of translation motion  $\vec{v}_{cm}$  of CM  
& rotation about an axis through the CM

Component particle  $m_i$  at  $\vec{r}_i$  has  $\vec{v}_i = \vec{v}_{cm} + \vec{v}'_i$   $\Leftarrow$  **vel. rel. to CM**



$$K_i = \frac{1}{2} m_i (\vec{v}_{cm} + \vec{v}'_i) \cdot (\vec{v}_{cm} + \vec{v}'_i)$$

$$\Rightarrow K_i = \frac{1}{2} m_i (v_{cm}^2 + 2\vec{v}_{cm} \cdot \vec{v}'_i + v_i'^2)$$

$$\Rightarrow K = \frac{1}{2} \left( \sum_i m_i \right) v_{cm}^2 + \vec{v}_{cm} \cdot \left( \sum_i m_i \vec{v}'_i \right) + \sum_i \left( \frac{1}{2} m_i v_i'^2 \right)$$

Since  $\sum_i m_i \vec{v}'_i = M \times \text{Vel. of CM relative to CM} = 0$

$$\Rightarrow \boxed{K = \frac{1}{2} M v_{cm}^2 + 0 + \frac{1}{2} I_{cm} \omega^2}$$

## Work Done By Torque In Rotational Motion

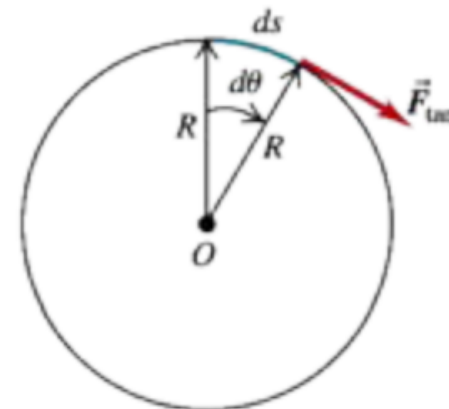
Tangential force  $\vec{F}$  over time  $dt$  applied at rim of a disk causes torque  $\vec{\tau}$ , leads to ang. displacement  $d\theta$

Work done  $dW = F_{\text{tan}} ds = F_{\text{tan}} R d\theta$

$$\Rightarrow dW = \tau_z d\theta \Rightarrow W = \int_{\theta_1}^{\theta_2} \tau_z d\theta$$

If applied torque is constant

$$\Rightarrow W = \int_{\theta_1}^{\theta_2} \tau_z d\theta = \tau_z (\theta_2 - \theta_1)$$



Overhead view of merry-go-round

# Work & Power In Rotational Motion

As result of work done by  $\vec{\tau}$ , kinetic energy changes

Since  $\vec{\tau} = I\vec{\alpha}_z \Rightarrow \tau_z d\theta = I\alpha_z d\theta$

$$\tau_z d\theta = I \frac{d\omega_z}{dt} d\theta = I \frac{d\theta}{dt} d\omega = I\omega_z d\omega_z$$

$$\Rightarrow W = \int_{\theta_1}^{\theta_2} \tau_z d\theta = \int_{\omega_1}^{\omega_2} I\omega_z d\omega_z = \frac{1}{2} I(\omega_2^2 - \omega_1^2) = \Delta K$$

work-energy theorem  
for rotating rigid bodies

Power associated with applied external torque:

$$P = \frac{dW}{dt} = \tau_z \frac{d\theta}{dt} = \tau_z \omega_z$$

# Cable Unwinding Off A Cylinder

Cable wrapped around cylinder (mass  $M$ ) is attached to object mass  $m$ . As cable unwinds,  $U_{\text{grav}}$  converted to kinetic energy.

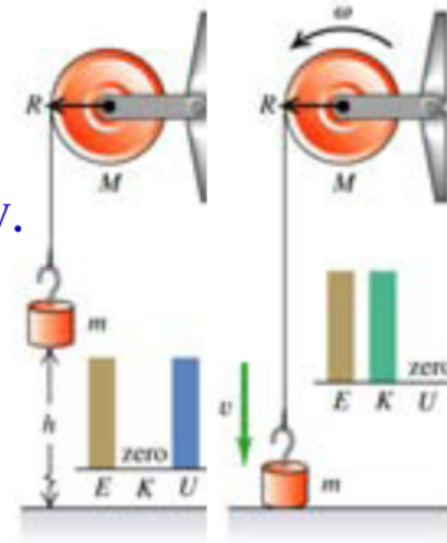
Find speed of object as it hits floor

$$K_1 + U_1 = K_2 + U_2$$

$$0 + mgh = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2 + 0$$

$$\Rightarrow mgh = \frac{1}{2}mv^2 + \frac{1}{2}\left(\frac{1}{2}MR^2\right)\left(\frac{v}{R}\right)^2$$

$$\Rightarrow v = \sqrt{\frac{2gh}{1 + M/2m}}$$



Cylinder is Solid