PHYSICS 140A : STATISTICAL PHYSICS HW ASSIGNMENT #4 SOLUTIONS

(1) Consider a *d*-dimensional ultrarelativistic gas of classical indistinguishable particles with a dispersion $\varepsilon(\mathbf{p}) = c |\mathbf{p}|$.

- (a) Find an expression for the grand potential $\Omega(T, V, \mu)$.
- (b) Find the average number of particles $N(T, V, \mu)$.
- (c) Find the entropy $S(T, V, \mu)$.
- (d) Express the RMS fluctuations in the number of particle number, $(\Delta N)_{\text{RMS}}$, in terms of the volume *V*, temperature *T*, and the pressure *p*.

Solution :

(a) The OCE partition function Z(T, V, N) is computed in §4.2.4 of the Lecture Notes. One finds

$$Z(T, V, N) = \frac{V^N}{N!} \left(\frac{\Gamma(d) \,\Omega_d}{(\beta h c)^d}\right)^N \,.$$

From $\Xi=e^{-\beta \varOmega}=\sum_{N=0}^{\infty}e^{\beta \mu N}\,Z(T,V,N)$, we obtain

$$\Omega(T, V, \mu) = -\frac{\Gamma(d)\Omega_d}{(hc)^d} V (k_{\rm B}T)^{d+1} e^{\mu/k_{\rm B}T} .$$

(b) The particle number is

$$N(T, V, \mu) = -\left(\frac{\partial \Omega}{\partial \mu}\right)_{T, V} = -(d+1)\frac{\Gamma(d) \Omega_d}{(hc)^d} V (k_{\rm B}T)^d e^{\mu/k_{\rm B}T}$$

(c) The entropy is

$$S(T, V, \mu) = -\left(\frac{\partial \Omega}{\partial T}\right)_{V,\mu} = \left((d+1)k_{\rm B} - \frac{\mu}{T}\right)\frac{\Gamma(d)\Omega_d}{(hc)^d}V(k_{\rm B}T)^d e^{\mu/k_{\rm B}T}.$$

(d) The variance of the number \hat{N} is (see eqn. 4.138 of the Lecture Notes)

$$\mathrm{var}(\hat{N}) = k_{\mathrm{B}}T \, \left(\frac{\partial N}{\partial \mu}\right)_{\!\!T,V} = N = \frac{pV}{k_{\mathrm{B}}T} \, . \label{eq:var}$$

Thus,

$$(\Delta N)_{\rm RMS} = \sqrt{{\rm var}(\hat{N})} = \left(\frac{pV}{k_{\rm B}T}\right)^{\!\!\!\!1/2}. \label{eq:mass_mass_star}$$

(2) Consider again the *d*-dimensional classical ultrarelativistic gas with $\varepsilon(\mathbf{p}) = c\mathbf{p}$.

- (a) If d = 3, find the momentum distribution function g(p).
- (b) Again for d = 3, find a general formula for the moments $\langle |\mathbf{p}|^k \rangle$.
- (c) Repeat parts (a) and (b) for the case d = 2.
- (d) In d = 3, what is the distribution function f(v) for velocities?

Solution :

(a) We have

$$g(\mathbf{p}) = \left\langle \delta(\mathbf{p} - \mathbf{p}_1) \right\rangle = \frac{e^{-\beta cp}}{\int d^3p \, e^{-\beta cp}} = \frac{c^3}{8\pi (k_{\rm B}T)^3} \, e^{-\beta cp} \, .$$

(b) The moments are

$$\left< |\mathbf{p}|^k \right> = \frac{1}{2} (\beta c)^3 \int_0^\infty dp \, p^{2+k} \, e^{-\beta cp} = \frac{1}{2} (k+2)! \, (\beta c)^{-k}$$

(c) In d = 2,

$$g(\boldsymbol{p}) = \left\langle \delta(\boldsymbol{p} - \boldsymbol{p}_1) \right\rangle = rac{e^{-eta c p}}{\int d^2 p \, e^{-eta c p}} = rac{c^2}{2\pi (k_{
m B}T)^2} \, e^{-eta c p}$$

and

$$\langle |\mathbf{p}|^k \rangle = (\beta c)^2 \int_0^\infty dp \, p^{1+k} \, e^{-\beta cp} = (k+1)! \, (\beta c)^{-k}$$

(d) The velocity is $v = \frac{\partial \varepsilon}{\partial p} = c \hat{p}$. Thus, the magnitude is fixed at |v| = c and the direction is distributed isotropically, *i.e.*

$$f(\boldsymbol{v}) = \frac{\delta(v-c)}{4\pi c^2} \,.$$

(3) A classical gas of indistinguishable particles in three dimensions is described by the Hamiltonian

$$\hat{H} = \sum_{i=1}^{N} \left\{ A |\mathbf{p}_{i}|^{3} - \mu_{0} H S_{i} \right\},\,$$

where A is a constant, and where $S_i \in \{-1, 0, +1\}$ (*i.e.* there are three possible spin polarization states).

- (a) Compute the free energy $F_{gas}(T, H, V, N)$.
- (b) Compute the magnetization density $m_{gas} = M_{gas}/V$ as a function of temperature, pressure, and magnetic field.

The gas is placed in thermal contact with a surface containing N_s adsorption sites, each with adsorption energy $-\Delta$. The surface is metallic and shields the adsorbed particles from the magnetic field, so the field at the surface may be approximated by H = 0.

- (c) Find the Landau free energy for the surface, $\Omega_{surf}(T, N_s, \mu)$.
- (d) Find the fraction $f_0(T, \mu)$ of empty adsorption sites.
- (e) Find the gas pressure $p^*(T, H)$ at which $f_0 = \frac{1}{2}$.

Solution :

(a) The single particle partition function is

$$\zeta(T,V,H) = V \int \frac{d^3p}{h^3} e^{-Ap^3/k_{\rm B}T} \sum_{S=-1}^{1} e^{\mu_0 H S/k_{\rm B}T} = \frac{4\pi V k_{\rm B}T}{3Ah^3} \cdot \left(1 + 2\cosh(\mu_0 H/k_{\rm B}T)\right).$$

The *N*-particle partition function is $Z_{gas}(T, H, V, N) = \zeta^N / N!$, hence

$$F_{\rm gas} = -Nk_{\rm B}T \left[\ln\left(\frac{4\pi Vk_{\rm B}T}{3NAh^3}\right) + 1 \right] - Nk_{\rm B}T\ln\left(1 + 2\cosh(\mu_0 H/k_{\rm B}T)\right)$$

(b) The magnetization density is

$$m_{\rm gas}(T,p,H) = -\frac{1}{V} \frac{\partial F}{\partial H} = \frac{p\mu_0}{k_{\rm B}T} \cdot \frac{2\sinh(\mu_0 H/k_{\rm B}T)}{1+2\cosh(\mu_0 H/k_{\rm B}T)}$$

We have used the ideal gas law, $pV = Nk_{\rm B}T$ here.

(c) There are four possible states for an adsorption site: empty, or occupied by a particle with one of three possible spin polarizations. Thus, $\Xi_{surf}(T, N_s, \mu) = \xi^{N_s}$, with

$$\xi(T,\mu) = 1 + 3 e^{(\mu+\Delta)/k_{\rm B}T}$$
.

Thus,

$$\label{eq:surf} \varOmega_{\rm surf}(T,N_{\rm s},\mu) = -N_{\rm s}k_{\rm B}T\ln\Bigl(1+3\,e^{(\mu+\Delta)/k_{\rm B}T}\Bigr)$$

(d) The fraction of empty adsorption sites is $1/\xi$, *i.e.*

$$f_0(T,\mu) = \frac{1}{1 + 3 e^{(\mu + \Delta)/k_{\rm B}T}}$$

(e) Setting $f_0 = \frac{1}{2}$, we obtain the equation $3 e^{(\mu + \Delta)/k_{\rm B}T} = 1$, or

$$e^{\mu/k_{\rm B}T} = \frac{1}{3} e^{-\Delta/k_{\rm B}T}$$
.

We now need the fugacity $z = e^{\mu/k_BT}$ in terms of p, T, and H. To this end, we compute the Landau free energy of the gas,

$$\label{eq:gas} \varOmega_{\rm gas} = -pV = -k_{\rm B}T\,\zeta\,e^{\mu/k_{\rm B}T}\,.$$

Thus,

$$p^{*}(T,H) = \frac{k_{\rm B}T\,\zeta}{V}\,e^{\mu/k_{\rm B}T} = \frac{4\pi(k_{\rm B}T)^{2}}{9Ah^{3}}\cdot\Big(1+2\cosh(\mu_{0}H/k_{\rm B}T)\Big)e^{-\Delta/k_{\rm B}T}$$