Fall Quarter 2009 UCSD Physics 214 & UCSB Physics 225 Homework 4

This Homework assignment covers a mix of problems from symmetries and neutrino oscillations. For the latter, you might find the following general hints useful:

General hints

A useful trick to simplify neutrino oscillation problems is to make use of unitarity, for example:

$$\begin{split} U_{ei}U_{ej}^{*} + U_{\mu i}U_{\mu j}^{*} + U_{\tau i}U_{\tau j}^{*} &= \delta_{ij} \\ \text{Then for } i \neq j \text{ multiply by } U_{ei}^{*}U_{ej} \text{ to get} \\ |U_{ei}U_{ej}^{*}|^{2} + U_{ei}^{*}U_{ej}(U_{\mu i}U_{\mu j}^{*} + U_{\tau i}U_{\tau j}^{*}) &= 0 \\ U_{ei}^{*}U_{ej}(U_{\mu i}U_{\mu j}^{*} + U_{\tau i}U_{\tau j}^{*}) &= -|U_{ei}U_{ej}^{*}|^{2} \\ \text{Another useful trick is} \\ \sum U_{\alpha i}^{*}U_{\beta i} &= \delta_{\alpha\beta} \\ \text{Take the square of this:} \\ \sum U_{\alpha i}^{*}U_{\beta i}U_{\alpha j}U_{\beta j}^{*} &= \delta_{\alpha\beta} \\ \sum_{i}U_{\alpha i}^{*}U_{\beta i}U_{\alpha j}U_{\beta j}^{*} &= \delta_{\alpha\beta} \\ \sum_{i}U_{\alpha i}^{*}U_{\beta i}U_{\alpha i}U_{\beta i}^{*} + 2\sum_{i < j}U_{\alpha i}^{*}U_{\beta i}U_{\alpha j}U_{\beta j}^{*} &= \delta_{\alpha\beta} \\ \text{Rearranging:} \\ \sum_{i}U_{\alpha i}^{*}U_{\beta i}U_{\alpha i}U_{\beta i}^{*} &= \delta_{\alpha\beta} - 2\sum_{i < j}U_{\alpha i}^{*}U_{\beta i}U_{\alpha j}U_{\beta j}^{*} \end{split}$$

You can do similar things interchanging the roles of rows and columns.

- 1. The amplitude for $v_{\alpha} \rightarrow v_{\beta}$ in the general case of N neutrino generations is given by equation (10) in hep-ph/0506165.
 - a. Show that the probability for $v_{\alpha} \rightarrow v_{\beta}$ oscillations is given by equation (11) in the same paper. If you can not obtain this result, just take it as a given and proceed.
 - b. If the neutrinos are Majorana particles, the neutrino mixing matrix includes two extra phases, as shown in equation (40) of the same paper. Verify that these Majorana phases have no effect on the oscillation probability of equation (11).
 - c. As discussed on page 4 of this paper, there is CP-violation in neutrino oscillations if the imaginary part $\Im(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \neq 0$ for α not equal

β. Use the unitarity of the neutrino mixing matrix to show that for a given pair of flavors, α not equal β, the quantity $\Im(U^*_{\alpha i}U_{\beta i}U_{\alpha j}U^*_{\beta j})$ is the same (up to a sign) for any *i*,*j*. This argument, which is based on unitarity, can be expanded by interchanging the roles of rows and

columns of U to conclude that for a given pair of *i*, *j* this quantity is independent of α and β (up to a sign). Thus CP violation is universal in the sense that is independent of the choice of $i_{i,j}$, α and β . To estimate whether the amount of CP-violation allowed in the Standard Model, after including guark and lepton sector, is sufficient to explain the lack of anti-matter in the universe, we thus need to measure this invariant in both the quark and lepton sector. Use the expression for U given in equation (40) of this paper (without the Majorana phases) to work out the value of $\Im(U^*_{\alpha i}U_{\beta i}U_{\alpha i}U^*_{\beta i})$. Verify that this quantity is proportional to s12 s23 s13 sin δ . This shows that CP-violation needs not only a complex phase δ but also non-trivial mixing between all three neutrino states. Mixing provides the CP-invariant phase we talked about in class for the simpler example of CP violation in decay of particles. It explains why measuring the smallest of the three angles determines how difficult the measurement of CP-violation in the neutrino sector will be because it sets the scale. It thus explains the importance of various attempts at measuring s13 worldwide. Note: To save a couple lines of trivial algebra, you might want to take the mixing matrix from 2006 PDG, page 138, rather than this paper. The PDG matrix is for quarks rather than neutrinos, but the functional form is identical.

- 2. This problem is for you to explore the argument behind the "off-axis" neutrino oscillation experiments like Nova.
 - a. Consider a pion of energy E in the lab frame. Find the lab energy of the neutrino in the decay: $\pi^+ \rightarrow \mu^+ v_{\mu}$ as a function of the lab angle θ that the emitted neutrino makes with the original flight direction of the charged pion.
 - b. Plot the energy of the neutrino for pion energies between 2GeV and 20GeV in the case $\theta = 0$ and $\theta = 15$ mrad. You should find that the spread in energy for the "off-axis" case is much reduced. This is used by some of the next generation of neutrino oscillation experiments to tune energy and baseline to sit right at the maximum of the oscillations.
- 3. Reactors are a source of electron anti-neutrinos. The next generation of neutrino reactor experiments will attempt to measure $\sin^2 2\theta_{13}$ of the neutrino mixing matrix. They will measure the probability for electron anti-neutrino disappearance a distance L from the reactor. Show that this probability is given by: $P = \sin^2 2\theta_{13} \sin^2 \Delta_{13} + \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \Delta_{21}$ where θ_{ij} are the

angles of the neutrino mixing matrix and $\Delta_{ij} = \frac{1}{4} (m_i^2 - m_j^2) L/E$. You can

make the approximation that $\Delta_{31} = \Delta_{32}$. Next, verify that for the typical E = 3 MeV and L = 1.5 Km, Δ_{31} is close to maximal (90°) and Δ_{21} is quite small. Take $|\Delta m_{21}^2| = |m_2^2 - m_1^2| = 8.0 \ 10^{-5} \ eV^2$ and $|\Delta m_{31}^2| = |m_3^2 - m_1^2| = 2.4 \ 10^{-3} \ eV^2$. These new experiments, Daya Bay in China and Double Chooz in France, propose to reach sensitivities of 0.01 and 0.03 respectively. This means that they will be looking for disappearance probabilities at the percent level or below.