

- 1-5 This is a case of dilation. $T = \gamma T'$ in this problem with the proper time $T' = T_0$

$$T = \left[1 - \left(\frac{v}{c} \right)^2 \right]^{-1/2} T_0 \Rightarrow \frac{v}{c} = \left[1 - \left(\frac{T_0}{T} \right)^2 \right]^{1/2};$$

in this case $T = 2T_0$, $v = \left\{ 1 - \left[\frac{L_0/2}{L_0} \right]^2 \right\}^{1/2} = \left[1 - \left(\frac{1}{4} \right) \right]^{1/2}$ therefore $v = 0.866c$.

- 1-6 This is a case of length contraction. $L = \frac{L'}{\gamma}$ in this problem the proper length $L' = L_0$,

$$L = \left[1 - \frac{v^2}{c^2} \right]^{-1/2} L_0 \Rightarrow v = c \left[1 - \left(\frac{L}{L_0} \right)^2 \right]^{1/2}; \text{ in this case } L = \frac{L_0}{2}, v = \left\{ 1 - \left[\frac{L_0/2}{L_0} \right]^2 \right\}^{1/2} = \left[1 - \left(\frac{1}{4} \right) \right]^{1/2}$$

therefore $v = 0.866c$.

- 1-7 The problem is solved by using time dilation. This is also a case of $v \ll c$ so the binomial

expansion is used $\Delta t = \gamma \Delta t' \cong \left[1 + \frac{v^2}{2c^2} \right] \Delta t'$, $\Delta t - \Delta t' = \frac{v^2 \Delta t'}{2c^2}$; $v = \left[\frac{2c^2 (\Delta t - \Delta t')}{\Delta t'} \right]^{1/2}$;

$$\Delta t = (24 \text{ h/day})(3600 \text{ s/h}) = 86400 \text{ s}; \Delta t = \Delta t' - 1 = 86399 \text{ s};$$

$$v = \left[\frac{2(86400 \text{ s} - 86399 \text{ s})}{86399 \text{ s}} \right]^{1/2} = 0.0048c = 1.44 \times 10^6 \text{ m/s}.$$

1-8 $L = \frac{L'}{\gamma}$

$$\frac{1}{\gamma} = \frac{L}{L'} = \left[1 - \frac{v^2}{c^2} \right]^{1/2}$$

$$v = c \left[1 - \left(\frac{L}{L'} \right)^2 \right]^{1/2} = c \left[1 - \left(\frac{75}{100} \right)^2 \right]^{1/2} = 0.661c$$

1-9 $L_{\text{earth}} = \frac{L'}{\gamma}$

$$L_{\text{earth}} = L' \left[1 - \frac{v^2}{c^2} \right]^{1/2}, L', \text{ the proper length so } L_{\text{earth}} = L = L' \left[1 - (0.9)^2 \right]^{1/2} = 0.436L'.$$

- 1-10 (a) $\tau = \gamma \tau'$ where $\beta = \frac{v}{c}$ and

$$\gamma = (1 - \beta^2)^{-1/2} = \tau' \left(1 - \frac{v^2}{c^2} \right)^{-1/2} = (2.6 \times 10^{-8} \text{ s}) \left[1 - (0.95)^2 \right]^{-1/2} = 8.33 \times 10^{-8} \text{ s}$$

(b) $d = v\tau = (0.95)(3 \times 10^8)(8.33 \times 10^{-8} \text{ s}) = 24 \text{ m}$

1-11 $\Delta t = \gamma \Delta t'$

$$\Delta t = \Delta t' \left(1 - \frac{v^2}{c^2}\right)^{-1/2} \cong \left(1 + \frac{v^2}{2c^2}\right) \Delta t' \cong \left[1 + \frac{(4.0 \times 10^2 \text{ m/s})^2}{2(3.0 \times 10^8 \text{ m/s})^2}\right] (3600 \text{ s})$$

$$\cong (1 + 8.89 \times 10^{-13})(3600 \text{ s}) = (3600 + 3.2 \times 10^{-9}) \text{ s}$$

$$\Delta t - \Delta t' \cong 3.2 \text{ ns. (Moving clocks run slower.)}$$

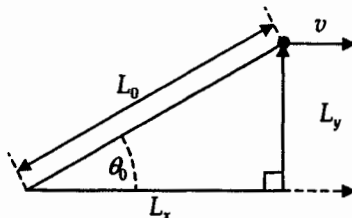
1-12 (a) 70 beats/min or $\Delta t' = \frac{1}{70} \text{ min}$

(b) $\Delta t = \gamma \Delta t' = [1 - (0.9)^2]^{-1/2} \left(\frac{1}{70}\right) \text{ min} = 0.0328 \text{ min/beat}$ or the number of beats per minute $\approx 30.5 \approx 31$.

1-13 (a) $\tau = \gamma \tau' = [1 - (0.95)^2]^{-1/2} (2.2 \mu\text{s}) = 7.05 \mu\text{s}$

(b) $\Delta t' = \frac{d}{0.95c} = \frac{3 \times 10^3 \text{ m}}{0.95c} = 1.05 \times 10^{-5} \text{ s}$, therefore,

$$N = N_0 \exp\left(-\frac{\Delta t}{\tau}\right) = (5 \times 10^4 \text{ muons}) \exp(-1.487) \approx 1.128 \times 10^4 \text{ muons.}$$

1-14 (a) Only the x -component of L_0 contracts.

$$L_{x'} = L_0 \cos \theta_0 \Rightarrow \frac{L_x [L_0 \cos \theta_0]}{\gamma}$$

$$L_{y'} = L_0 \sin \theta_0 \Rightarrow L_y = L_0 \sin \theta_0$$

$$L = [(L_x)^2 + (L_y)^2]^{1/2} = \left[\left(\frac{L_0 \cos \theta_0}{\gamma} \right)^2 + (L_0 \sin \theta_0)^2 \right]^{1/2}$$

$$= L_0 \left[\cos^2 \theta_0 \left(1 - \frac{v^2}{c^2}\right) + \sin^2 \theta_0 \right]^{1/2} = L_0 \left[1 - \frac{v^2}{c^2} \cos^2 \theta_0 \right]^{1/2}$$

(b) As seen by the stationary observer, $\tan \theta = \frac{L_y}{L_x} = \frac{L_0 \sin \theta_0}{L_0 \cos \theta_0 / \gamma} = \gamma \tan \theta_0$.

- 1-15 (a) For a receding source we replace v by $-v$ in Equation 1.15 and obtain:

$$\begin{aligned} f_{\text{ob}} &= \left\{ \frac{[c-v]^{1/2}}{[c+v]^{1/2}} \right\} f_{\text{source}} = \left\{ \frac{[1-v/c]^{1/2}}{[1+v/c]^{1/2}} \right\} f_{\text{source}} \cong \left(1 - \frac{v}{2c} \right) \left(1 - \frac{v}{2c} \right) f_{\text{source}} \\ &\cong \left(1 - \frac{v}{c} + \frac{v^2}{4c^2} \right) f_{\text{source}} \cong \left(1 - \frac{v}{c} \right) f_{\text{source}} \end{aligned}$$

where we have used the binomial expansion and have neglected terms of second and higher order in $\frac{v}{c}$. Thus, $\frac{\Delta f}{f_{\text{source}}} = \frac{f_{\text{ob}} - f_{\text{source}}}{f_{\text{source}}} = -\frac{v}{c}$.

- (b) From the relations $f = \frac{c}{\lambda}$, $\frac{df}{d\lambda} = -\frac{c}{\lambda^2}$ we find $\frac{df}{f} = -\frac{c/\lambda^2}{c/\lambda} d\lambda$, or $\frac{\Delta\lambda}{\lambda} = -\frac{\Delta f}{f} = \frac{v}{c}$.

- (c) Assuming $v \ll c$, $\frac{v}{c} \cong \frac{\Delta\lambda}{\lambda}$, or $v \cong \left(\frac{\Delta\lambda}{\lambda} \right) c = \left(\frac{20 \text{ nm}}{397 \text{ nm}} \right) c = 0.050c = 1.5 \times 10^7 \text{ m/s}$.

- 1-16 For an observer approaching a light source, $\lambda_{\text{ob}} = \left[\frac{(1-v/c)^{1/2}}{(1+v/c)^{1/2}} \right] \lambda_{\text{source}}$. Setting $\beta = \frac{v}{c}$ and after some algebra we find,

$$\beta = \frac{\lambda_{\text{source}}^2 - \lambda_{\text{obs}}^2}{\lambda_{\text{source}}^2 + \lambda_{\text{obs}}^2} = \frac{(650 \text{ nm})^2 - (550 \text{ nm})^2}{(650 \text{ nm})^2 + (550 \text{ nm})^2} = 0.166$$

$$v = 0.166c = (4.98 \times 10^7 \text{ m/s})(2.237 \text{ mi/h})(\text{m/s})^{-1} = 1.11 \times 10^8 \text{ mi/h}$$

- 1-17 (a) Galaxy A is approaching and as a consequence it exhibits blue shifted radiation. From Example 1.6, $\beta = \frac{v}{c} = \frac{\lambda_{\text{source}}^2 - \lambda_{\text{obs}}^2}{\lambda_{\text{source}}^2 + \lambda_{\text{obs}}^2}$ so that $\beta = \frac{(550 \text{ nm})^2 - (450 \text{ nm})^2}{(550 \text{ nm})^2 + (450 \text{ nm})^2} = 0.198$. Galaxy A is approaching at $v = 0.198c$.

- (b) For a red shift, B is receding. $\beta = \frac{v}{c} = \frac{\lambda_{\text{source}}^2 - \lambda_{\text{obs}}^2}{\lambda_{\text{source}}^2 + \lambda_{\text{obs}}^2}$ so that

$$\beta = \frac{(700 \text{ nm})^2 - (550 \text{ nm})^2}{(700 \text{ nm})^2 + (550 \text{ nm})^2} = 0.237. \text{ Galaxy B is receding at } v = 0.237c.$$

- 1-18 (a) Let f_c be the frequency as seen by the car. Thus, $f_c = f_{\text{source}} \sqrt{\frac{c+v}{c-v}}$ and, if f is the frequency of the reflected wave, $f = f_c \sqrt{\frac{c+v}{c-v}}$. Combining these equations gives

$$f = f_{\text{source}} \frac{(c+v)}{(c-v)}.$$

- (b) Using the above result, $f(c-v) = f_{\text{source}}(c+v)$, which gives

$$(f - f_{\text{source}})c = (f + f_{\text{source}})v = 2f_{\text{source}}v.$$

$$\text{The beat frequency is then } f_{\text{beat}} = f - f_{\text{source}} = \frac{2f_{\text{source}}v}{c} = \frac{2v}{\lambda}.$$

$$(c) \quad f_{\text{beat}} = \frac{2(30.0 \text{ m/s})(10.0 \times 10^9 \text{ Hz})}{3.00 \times 10^8 \text{ m/s}} = \frac{2(30.0 \text{ m/s})}{0.0300 \text{ m}} = 2000 \text{ Hz} = 2.00 \text{ kHz}$$

$$\lambda = \frac{c}{f_{\text{source}}} = \frac{3.00 \times 10^8 \text{ m/s}}{10.0 \times 10^9 \text{ Hz}} = 3.00 \text{ cm}$$

$$(d) \quad v = \frac{f_{\text{beat}} \lambda}{2} \text{ so,}$$

$$\Delta v = \frac{\Delta f_{\text{beat}} \lambda}{2} = \frac{(5 \text{ Hz})(0.0300 \text{ m})}{2} = 0.0750 \text{ m/s} \approx 0.2 \text{ mi/h}$$

$$1-19 \quad u_{XA} = -u_{XB}; u'_{XA} = 0.7c = \frac{u_{XA} - u_{XB}}{1 - u_{XA}u_{XB}/c^2}; 0.70c = \frac{2u_{XA}}{1 + (u_{XA}/c)^2} \text{ or } 0.70u_{XA}^2 - 2cu_{XA} + 0.7c^2 = 0.$$

Solving this quadratic equation one finds $u_{XA} = 0.41c$ therefore $u_{XB} = -u_{XA} = -0.41c$.

$$1-20 \quad u = \frac{v + u'}{1 + vu'/c^2} = \frac{0.90c + 0.70c}{1 + (0.90c)(0.70c)/c^2} = 0.98c$$

$$1-21 \quad u'_X = \frac{u_X - v}{1 - u_X v/c^2} = \frac{0.50c - 0.80c}{1 - (0.50c)(0.80c)/c^2} = -0.50c$$

1-22 (a) The speed as observed in the laboratory is found by using Equation 1.30:

$$u_X = \frac{u'_X - v}{1 - u'_X v/c^2}. \text{ But } u'_X = \frac{c}{n} \text{ (speed measured by an observer moving with the fluid),}$$

$$\text{therefore } u_X = \frac{(c/n) - v}{1 + v/(nc)} = \frac{c}{n} \frac{1 + nv/c}{1 + v/(nc)}.$$

(b) $\frac{v}{c} \ll 1$. Use the binomial expansion,

$$u'_X \cong \frac{c}{n} \left[1 + n \left(\frac{v}{c} \right) \right] \left[1 - \frac{v}{nc} \right] \cong \frac{c}{n} \left[1 + \frac{nv}{c} - \frac{v^2}{c^2} \right] \cong \frac{c}{n} + v - \frac{v^2}{n^2}.$$

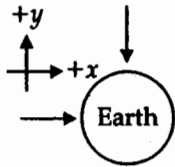
1-23 (a) Let event 1 have coordinates $x_1 = y_1 = z_1 = t_1 = 0$ and event 2 have coordinates $x_2 = 100 \text{ mm}$, $y_2 = z_2 = t_2 = 0$. In S' , $x'_1 = \gamma(x_1 - vt_1) = 0$, $y'_1 = y_1 = 0$, $z'_1 = z_1 = 0$, and $t'_1 = \gamma \left[t_1 - \left(\frac{v}{c^2} \right) x_1 \right] = 0$, with $\gamma = \left[1 - \frac{v^2}{c^2} \right]^{-1/2}$ and so $\gamma = \left[1 - (0.70)^2 \right]^{-1/2} = 1.40$. In system S' , $x'_2 = \gamma(x_2 - vt_2) = 140 \text{ m}$, $y'_2 = z'_2 = 0$, and

$$t'_2 = \gamma \left[t_2 - \left(\frac{v}{c^2} \right) x_2 \right] = \frac{(1.4)(-0.70)(100 \text{ m})}{3.00 \times 10^8 \text{ m/s}} = -0.33 \mu\text{s}.$$

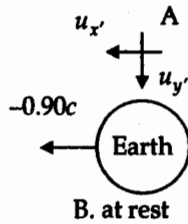
(b) $\Delta x' = x'_2 - x'_1 = 140 \text{ m}$

(c) Events are not simultaneous in S' , event 2 occurs $0.33 \mu\text{s}$ earlier than event 1.

1-24 A. $u_y = -0.90c$



B. $u_x = -0.90c$



$$u'_x = \frac{u_x - v}{1 - u_x v / c^2} = \frac{0 - 0.90c}{1 - (0)(0.90c) / c^2} = -0.90c$$

$$u'_y = \frac{u_y}{\gamma(1 - u_x v / c^2)} = \frac{0 - 0.90c}{[1 - 0.81]^{-1/2}} \equiv -0.392c$$

The speed of A as measured by B is

$$u_{AB} = [(u'_x)^2 + (u'_y)^2]^{1/2} = [(-0.90c)^2 + (-0.392c)^2]^{1/2} = 0.982c.$$

Classically, $u_{AB} = 1.3c$.

1-25 We find Carpenter's speed: $\frac{mGM}{r^2} = \frac{mv^2}{r}$

$$v = \left[\frac{GM}{R+h} \right]^{1/2} = \left[\frac{(6.67 \times 10^{-11})(5.98 \times 10^{24})}{6.37 \times 10^6 + 0.16 \times 10^6} \right]^{1/2} = 7.82 \text{ km/s.}$$

Then the period of one orbit is $T = \frac{2\pi(R+h)}{v} = \frac{2\pi(6.53 \times 10^6)}{7.82 \times 10^3} = 5.25 \times 10^3 \text{ s.}$

(a) The time difference for 22 orbits is $\Delta t - \Delta t' = (\gamma - 1)\Delta t' = \left[\left(1 - \frac{v^2}{c^2} \right)^{-1/2} - 1 \right] (22)(T).$

Using the binomial expansion one obtains

$$\left(1 + \frac{1}{2} \frac{v^2}{c^2} - 1 \right) (22)(T) = \frac{1}{2} \left[\frac{7.82 \times 10^3 \text{ m/s}}{3 \times 10^8 \text{ m/s}} \right] (22)(5.5 \times 10^3 \text{ s}) = 39.2 \mu\text{s.}$$

(b) For one orbit, $\Delta t - \Delta t' = \frac{39.2 \mu\text{s}}{22} = 1.78 \mu\text{s} \approx 2 \mu\text{s}$. The press report is accurate to one significant figure.

- 1-26 The observed length of an object moving with speed v is $L = L' \left[1 - \left(\frac{v}{c} \right)^2 \right]^{1/2}$ with L' being the proper length. For the two ships, we know that $L_2 = L_1$, $L'_2 = 3L'_1$ and $v_1 = 0.35c$. Thus $L_2^2 = L_1^2$ and $(9L_1'^2) \left[1 - \left(\frac{v_2}{c} \right)^2 \right] = L_1'^2 \left[1 - (0.35)^2 \right]$, giving $9 - 9 \left(\frac{v_2}{c} \right)^2 = 0.8775$, or $v_2 = 0.95c$.

- 1-27 For the pion to travel 10 m in time Δt in our frame,

$$10 \text{ m} = v\Delta t = v(\gamma\Delta t') = v(26 \times 10^{-9} \text{ s}) \left[1 - \left(\frac{v}{c} \right)^2 \right]^{-1/2}$$

$$(3.85 \times 10^8 \text{ m/s})^2 \left(1 - \frac{v^2}{c^2} \right) = v^2$$

$$1.46 \times 10^{17} \text{ m}^2/\text{s}^2 = v^2(1 + 1.64)$$

$$v = 2.37 \times 10^8 \text{ m/s} = 0.789c$$

- 1-28 For Astronauts approaching Alpha Centauri $\gamma^{-1} = \left[1 - \left(\frac{v}{c} \right)^2 \right]^{1/2} = [1 - 0.902]^{1/2} = 0.312$.

(a) astronauts' time $t' = \gamma^{-1}t = (0.312)(4.4 \text{ y}) = 1.37 \text{ years}$,

(b) astronauts' distance $d' = (0.312)(4.2 \text{ light years}) = 1.31 \text{ light years}$.

- 1-29 (a) A spaceship, reference frame S' , moves at speed v relative to the Earth, whose reference frame is S . The space ship then launches a shuttle craft with velocity v in the forward direction. The pilot of the shuttle craft then fires a probe with velocity v in the forward direction. Use the relativistic compounding of velocities as well as its inverse transformation: $u'_x = \frac{u_x - v}{1 - (u_x v/c^2)}$, and its inverse $u_x = \frac{u'_x + v}{1 + (u'_x v/c^2)}$. The above

variables are defined as: v is the spaceship's velocity relative to S , u'_x is the velocity of the shuttle craft relative to S' , and u_x is the velocity of the shuttle craft relative to S . Setting u'_x equal to v , we find the velocity of the shuttle craft relative to the Earth to

$$\text{be: } u_x = \frac{2v}{1 + (v/c)^2}.$$

- (b) If we now take S to be the shuttle craft's frame of reference and S' to be that of the probe whose speed is v relative to the shuttle craft, then the speed of the probe relative to the spacecraft will be, $u'_x = \frac{2v}{1 + (v/c)^2}$. Adding the speed relative to S yields:

$$u_x = \left[\frac{3 + (v/c)^2}{1 + 2(v/c)^2} \right] = \frac{3v + v^3/c^3}{1 + 2v^2/c^2}. \text{ Using the Galilean transformation of velocities, we see}$$

that the spaceship's velocity relative to the Earth is v , the velocity of the shuttle craft relative to the space ship is v and therefore the velocity of the shuttle craft relative to the Earth must be $2v$ and finally the speed of the probe must be $3v$. In the limit of low $\left(\frac{v}{c} \right)^2$, u_x reduces to $3v$. On the other hand, using relativistic addition of velocities, we find that $u_x = c$ when $v \rightarrow c$.

- 1-35 In the Earth frame, Speedo's trip lasts for a time $\Delta t = \frac{\Delta x}{v} = \frac{20.0 \text{ ly}}{0.950 \text{ ly/yr}} = 21.05$ Speedo's age advances only by the proper time interval: $\Delta t_p = \frac{\Delta t}{\gamma} = 21.05 \text{ yr} \sqrt{1 - 0.95^2} = 6.574 \text{ yr}$ during his trip. Similarly for Goslo, $\Delta t_p = \frac{\Delta x}{v} \sqrt{1 - \frac{v^2}{c^2}} = \frac{20.0 \text{ ly}}{0.750 \text{ ly/yr}} \sqrt{1 - 0.75^2} = 17.64 \text{ yr}$. While Speedo has landed on Planet X and is waiting for his brother, he ages by

$$\frac{20.0 \text{ ly}}{0.750 \text{ ly/yr}} - \frac{0.20 \text{ ly}}{0.950 \text{ ly/yr}} \sqrt{1 - 0.75^2} = 17.64 \text{ yr}.$$

Then Goslo ends up older by $17.64 \text{ yr} - (6.574 \text{ yr} + 5.614 \text{ yr}) = 5.45 \text{ yr}$.

- 1-36 Let Suzanne be fixed in reference frame S and see the two light-emission events with coordinates $x_1 = 0$, $t_1 = 0$, $x_2 = 0$, $t_2 = 3 \mu\text{s}$. Let Mark be fixed in reference frame S' and give the events coordinate $x'_1 = 0$, $t'_1 = 0$, $t'_2 = 9 \mu\text{s}$.

(a) Then we have

$$t'_2 = \gamma \left(t_2 - \frac{v}{c^2} x_2 \right) = 9 \mu\text{s} = \frac{1}{\sqrt{1 - v^2/c^2}} (3 \mu\text{s} - 0) = \sqrt{1 - \frac{v^2}{c^2}} = \frac{1}{3} = \frac{v^2}{c^2} = \frac{8}{9} = v = 0.943c.$$

(b) $x'_2 = \gamma(x_2 - vt_2) = 3(0 - 0.943c \times 3 \times 10^{-6} \text{ s}) \left(\frac{3 \times 10^8 \text{ m/s}}{c} \right) = 2.55 \times 10^3 \text{ m}$

- 1-37 Einstein's reasoning about lightning striking the ends of a train shows that the moving observer sees the event toward which she is moving, event B, as occurring first. We may take the S -frame coordinates of the events as $(x = 0, y = 0, z = 0, t = 0)$ and $(x = 100 \text{ m}, y = 0, z = 0, t = 0)$. Then the coordinates in S' are given by Equations 1.23 to 1.27. Event A is at $(x' = 0, y' = 0, z' = 0, t' = 0)$. The time of event B is:

$$t' = \gamma \left(t - \frac{v}{c^2} x \right) = \frac{1}{\sqrt{1 - 0.8^2}} \left(0 - \frac{0.8c}{c^2} (100 \text{ m}) \right) = 1.667 \left(\frac{80 \text{ m}}{3 \times 10^8 \text{ m/s}} \right) = -4.44 \times 10^{-7} \text{ s}.$$

The time elapsing before A occurs is 444 ns.