

# Physics 214 UCSD/225a UCSB

## Lecture 7

### Finish Chapter 2 of H&M

- November revolution, charm and beauty

### CP symmetry and violation

- Simple example
- Unitarity matrix for leptons and quarks

### Beginning of Neutrino Physics

Missed a week due to fire in SD.

Let's skip some stuff!

- Magnetic moment of proton etc.
- November revolution
  - Charm
  - Beauty
  - OZI suppression
- I encourage you to read up on this in chapter 2 of H&M

# CP Symmetry

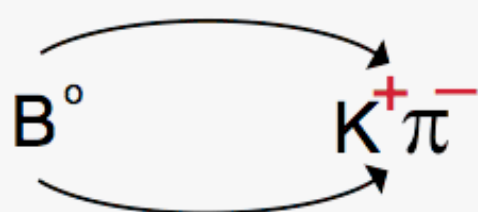


$$\mathcal{L}_{CC} = \frac{g_2}{2\sqrt{2}} J_\mu^+ W^{+\mu} + J_\mu^- W^{-\mu}$$

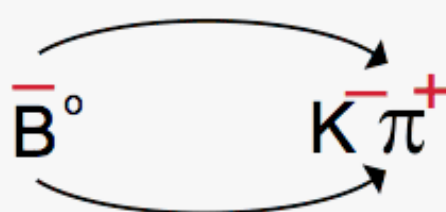
$$J_\mu^+ = (\bar{\nu}_{eL} \bar{\nu}_{\mu L} \bar{\nu}_{\tau L}) \gamma_\mu \begin{pmatrix} e_L^- \\ \mu_L^- \\ \tau_L^- \end{pmatrix} + (\bar{u}_L \bar{c}_L \bar{t}_L) \gamma_\mu \mathbf{V}_{CKM} \begin{pmatrix} d_L \\ s_L \\ b_L \end{pmatrix}$$

Note:

- > This requires CP because weak interactions maximally violate parity.
- > We will ignore subtleties in the difference between lepton and quark sector.
- ⇒ We'll get back to this next quarter.
- ⇒ All we care for now is that there's a *3x3 unitary matrix of couplings* involved.

$$T = |T| e^{-i(\delta - \gamma)}$$


$$P = |P|$$

$$\bar{T} = |T| e^{-i(\delta + \gamma)}$$


$$\bar{P} = |P|$$

$\delta =$  strong phase shift

$\gamma =$  difference in weak phase

$$\text{CP } \gamma = -\gamma \quad \text{CP } \delta = +\delta$$

$$A_{cp} = \frac{\mathcal{B}(B^0 \rightarrow K^+ \pi^-) - \mathcal{B}(\bar{B}^0 \rightarrow K^- \pi^+)}{\mathcal{B}(B^0 \rightarrow K^+ \pi^-) + \mathcal{B}(\bar{B}^0 \rightarrow K^- \pi^+)}$$

$$= \frac{-2|TP| \sin \gamma \sin \delta}{|T|^2 + |P|^2 + 2|TP| \cos \gamma \cos \delta}$$

## Breaking CP is easy

$\Rightarrow$  Add complex coupling  
to Lagrangian.

$\Rightarrow$  Allow 2 or more channels

$\Rightarrow$  Add CP symm. Phase,  
e.g. via dynamics.

# Breaking CP in Standard Model

- Where does the CP violating phase come from?
    - 3x3 unitary matrix  $\Rightarrow$  3 angles + 6 phases
      - $2N^2$  parameters,  $N^2$  constraints from unitarity
    - 6 spinors with arbitrary phase convention
      - Only relative phase matters because only  $|M|^2$  is physical.  
 $\Rightarrow$  Only 5 phases can be used to define a convention.
- $\Rightarrow$  **One phase left in 3x3 matrix that has physical consequences.**

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} c_x c_z & s_x c_z & s_z e^{-i\phi} \\ -s_x c_y - c_x s_y s_z e^{i\phi} & c_x c_y - s_x s_y s_z e^{i\phi} & s_y c_z \\ s_x s_y - c_x c_y s_z e^{i\phi} & -c_x s_y - s_x c_y s_z e^{i\phi} & c_y c_z \end{pmatrix}$$

x,y,z are euler angles. c=cos, s=sin.

Note:  $\sin(z) = 0 \Leftrightarrow$  NO CP violating phase left !!!

# CP violation summary

- CP violation is easy to add in field theory:
  - Complex coupling in Lagrangian
  - Interference of channels with:
    - Different CP violating phase
    - Different CP conserving phase
- Standard Model implements this via:
  - CP violating phase in charged current coupling across 3 families
  - CP conserving phase via:
    - Dynamics, e.g. Breit Wigner resonance lineshape
    - Flavor Mixing & oscillation in neutrino or quark sector

***Let's look at neutrino sector in some detail !***

# Aside:

- If you want to know more about the details, please check out:

[Lecture 9/20/2000 and further reading for it](#)

- It constructs all possible conventions for the CKM matrix in probably more detail than you ever want to know.

# Mixing in Standard Model

- Weak eigenstates not equal mass eigenstates.
    - Mass eigenstates responsible for propagation in time.
    - Weak eigenstates responsible for production and/or decay.
- ⇒ Oscillation between weak eigenstates as a function of time.
- ⇒ Discuss this in detail for Neutrino sector now.



# Neutrino mixing

- At the  $W$  vertex an electron-neutrino is created together with a positron.
- That electron-neutrino is a superposition of mass eigenstates:

$$|\nu_e(t)\rangle = \sum_{i=1}^3 U_{ei}^* |\nu_i(t)\rangle$$

- The time evolution of the mass eigenstate can be described either in its rest-frame or in the labframe:

$$|\nu_i(t)\rangle = e^{-im_i t} |\nu_i(0)\rangle = e^{-i(E_i t - p_i L)} |\nu_i(0)\rangle$$

- For interference among the mass eigenstates to be possible, they all have to have the same  $E$  because experimentally we average over time.

Time average demands  $E_i = E$

# Oscillation Amplitude

$$\text{Amp}(\nu_\mu \rightarrow \nu_\tau) = \langle \nu_\tau | e^{-iEt} \sum_{i=1}^3 e^{ip_i L} U_{\mu i}^* | \nu_i \rangle$$

$$\text{Amp}(\nu_\mu \rightarrow \nu_\tau) = e^{-iEt} \sum_{i,j=1}^3 e^{ip_i L} U_{\mu i}^* U_{\tau j} \langle \nu_j | \nu_i \rangle$$

$$\text{Amp}(\nu_\mu \rightarrow \nu_\tau) = e^{-iEt} \sum_{i=1}^3 e^{ip_i L} U_{\mu i}^* U_{\tau i}$$

Next we Taylor expand  $p_i$  using:

$$p_i = \sqrt{E^2 - m_i^2} = E - \frac{m_i^2}{2E} + \dots$$

# Oscillation Probability

$$\text{Amp}(\nu_\mu \rightarrow \nu_\tau) = e^{-iE(t-L)} \sum_{i=1}^3 e^{-i\frac{m_i^2}{2E}L} U_{\mu i}^* U_{\tau i}$$

$$\text{Prob}(\nu_\mu \rightarrow \nu_\tau) = \sum_{i=1}^3 \left| e^{-i\frac{m_i^2}{2E}L} U_{\mu i}^* U_{\tau i} \right|^2$$

In homework, you do this for the general case of N flavors. Here we do it for the simpler case of 2 flavors only.

# Simple math aside

$$\begin{aligned} |1 - e^{ix}|^2 &= (1 - [\cos x + i \sin x])(1 - [\cos x - i \sin x]) \\ &= [1 - \cos x]^2 + \sin^2 x \\ &= 2(1 - \cos x) \end{aligned}$$

We'll need this is a second.

# 2 flavor oscillation probability

$$\begin{aligned} & \left| U_{11}U_{21}e^{-im_1^2 \frac{L}{2E}} + U_{12}U_{22}e^{-im_2^2 \frac{L}{2E}} \right|^2 = \left| U_{11}U_{21} + U_{12}U_{22}e^{i(m_1^2 - m_2^2) \frac{L}{2E}} \right|^2 \\ & = \left| -\cos\theta \sin\theta + \cos\theta \sin\theta e^{i(m_1^2 - m_2^2) \frac{L}{2E}} \right|^2 = \cos^2\theta \sin^2\theta \left| 1 - e^{i(m_1^2 - m_2^2) \frac{L}{2E}} \right|^2 \\ & = \cos^2\theta \sin^2\theta \left[ (1 - \cos\Delta)^2 + \sin^2\Delta \right] = 2\cos^2\theta \sin^2\theta [1 - \cos\Delta] \\ & = \frac{1}{2} \sin^2 2\theta \left[ 2 \sin^2 \frac{\Delta}{2} \right] \\ & \Delta = (m_1^2 - m_2^2) \frac{L}{2E} \end{aligned}$$

This is a bit simplistic, as it ignores matter effects.  
We'll discuss those on Wednesday.

# Discussion of Oscillation Equation

$$\text{Prob}(\nu_e \rightarrow \nu_\mu) = \sin^2 2\theta \left[ \sin^2 \frac{(m_1^2 - m_2^2)L}{4E} \right]$$

- Depends on difference in mass squared.
  - No mixing if masses are identical
  - Insensitive to mass scale
  - Insensitive to mass hierarchy
- Depends on  $\sin^2(2\theta)$ 
  - Need large angle to see large effect
- Depends on  $L/4E$ 
  - Exp. with unfortunate  $L/E$  won't see any effect.
  - Exp. with variable  $L/E$  can measure both angle and mass squared difference.
  - Exp. with  $\Delta m^2 L/4E \gg 1$  and some energy spread average over  $\sin^2 \rightarrow 1/2$

# Experimental situation

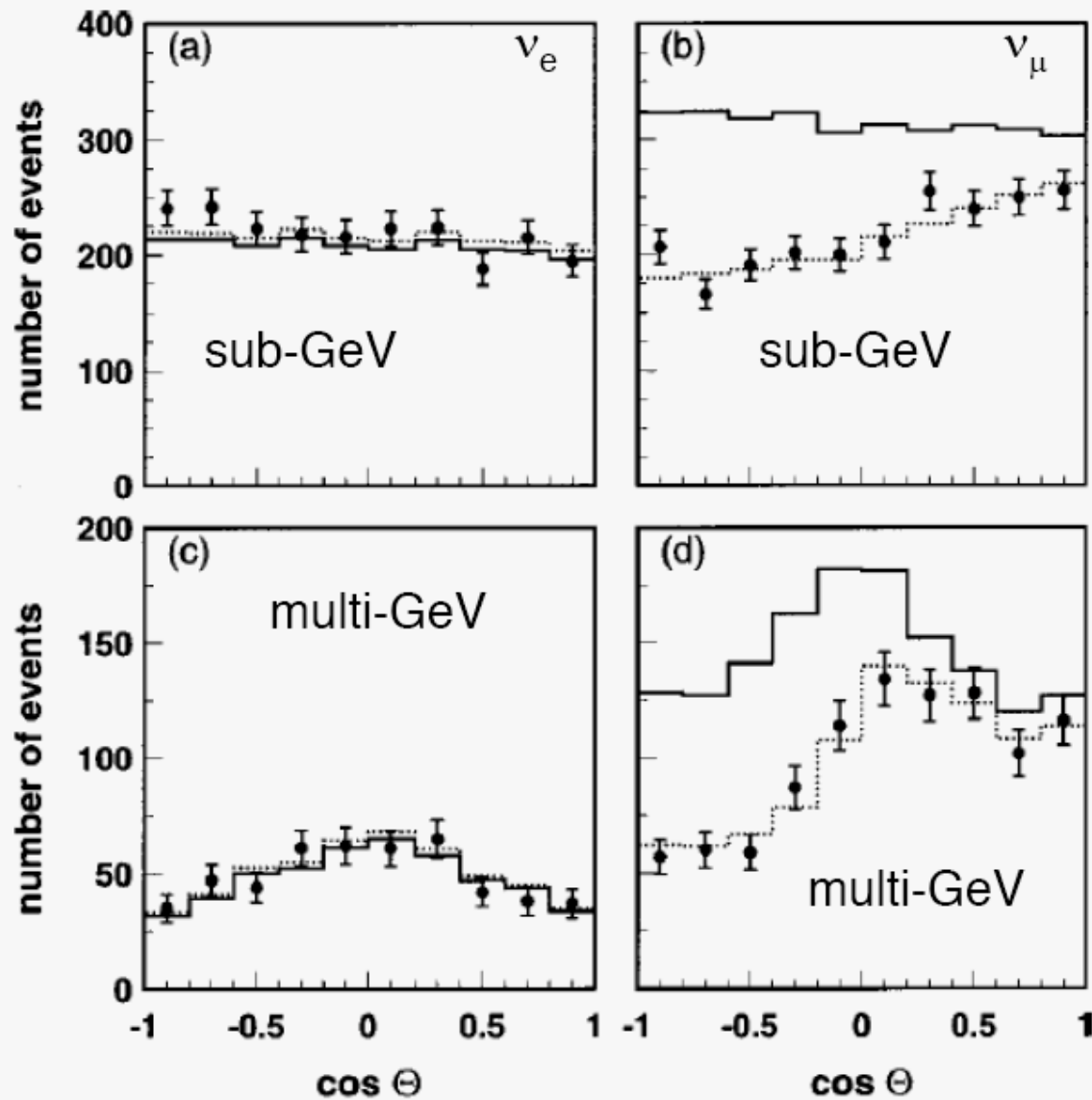
- Sources of electron neutrinos
  - Sun
  - Reactors
  
- Sources of muon neutrinos
  - From charged pion beams
  - From charged pion decay in atmosphere



# Atmospheric neutrinos

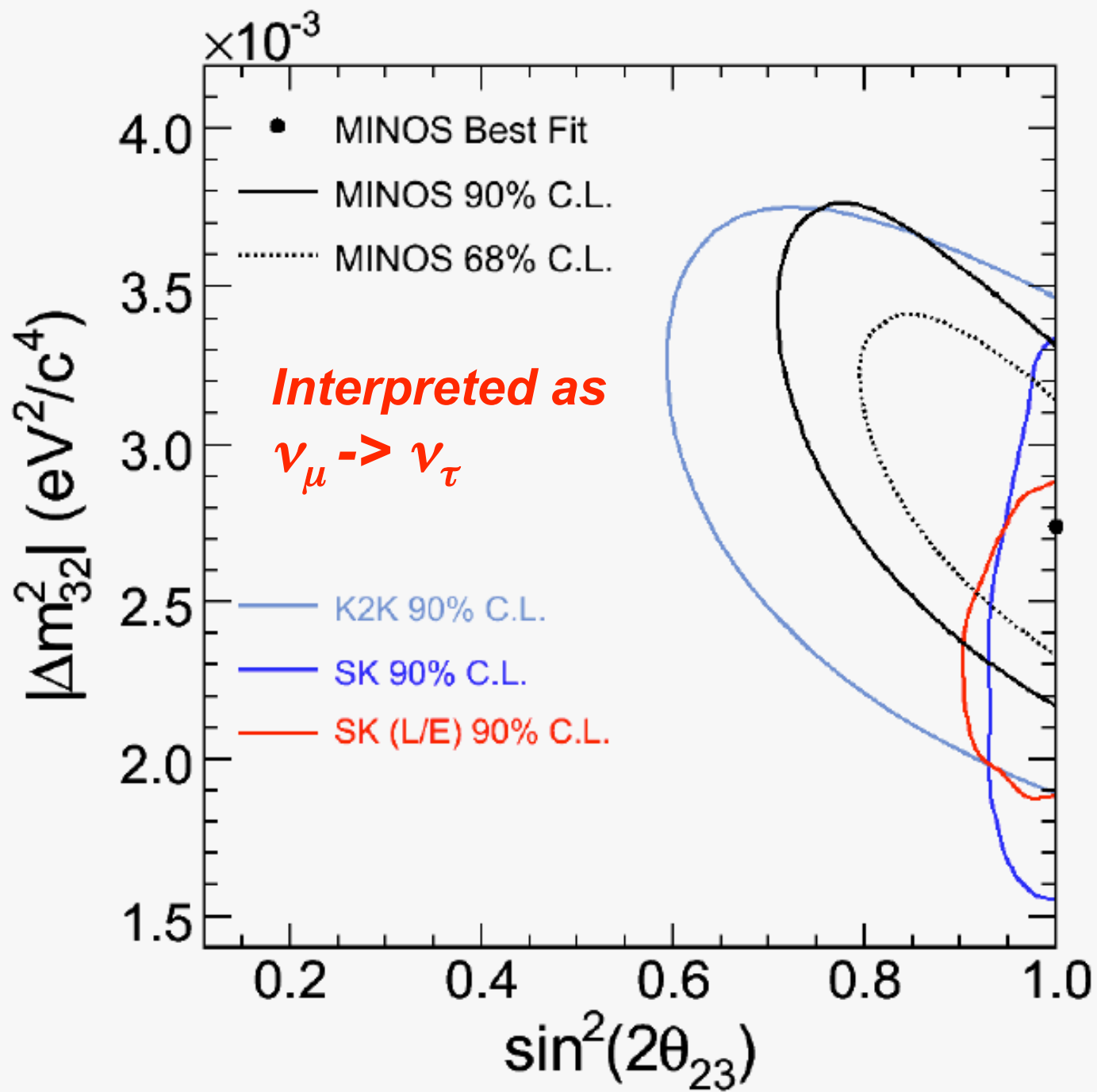
- Expect  $\nu_\mu$  anti- $\nu_\mu$  in equal numbers
- Expect  $\nu_e$  half as many as  $\nu_\mu + \text{anti-}\nu_\mu$
- Can change L as a function of Zenith angle. (L  $\sim$  15km to L  $\sim$  13,000km)
- $\nu_e$  Oscillation to  $\nu_\mu$   
=> See excess of  $\nu_\mu$  vs zenith angle
- $\nu_\mu$  Oscillation to  $\nu_e$   
=> See excess of  $\nu_e$  vs zenith angle
- $\nu_e$  Oscillation to  $\nu_\tau$   
=> Deficit of  $\nu_e$  vs zenith angle
- $\nu_\mu$  Oscillation to  $\nu_\tau$   
=> Deficit of  $\nu_\mu$  vs zenith angle

# Super Kamiokande Results



cos $\theta$ =-1 up-going,  
 $L \approx 13,000$  km

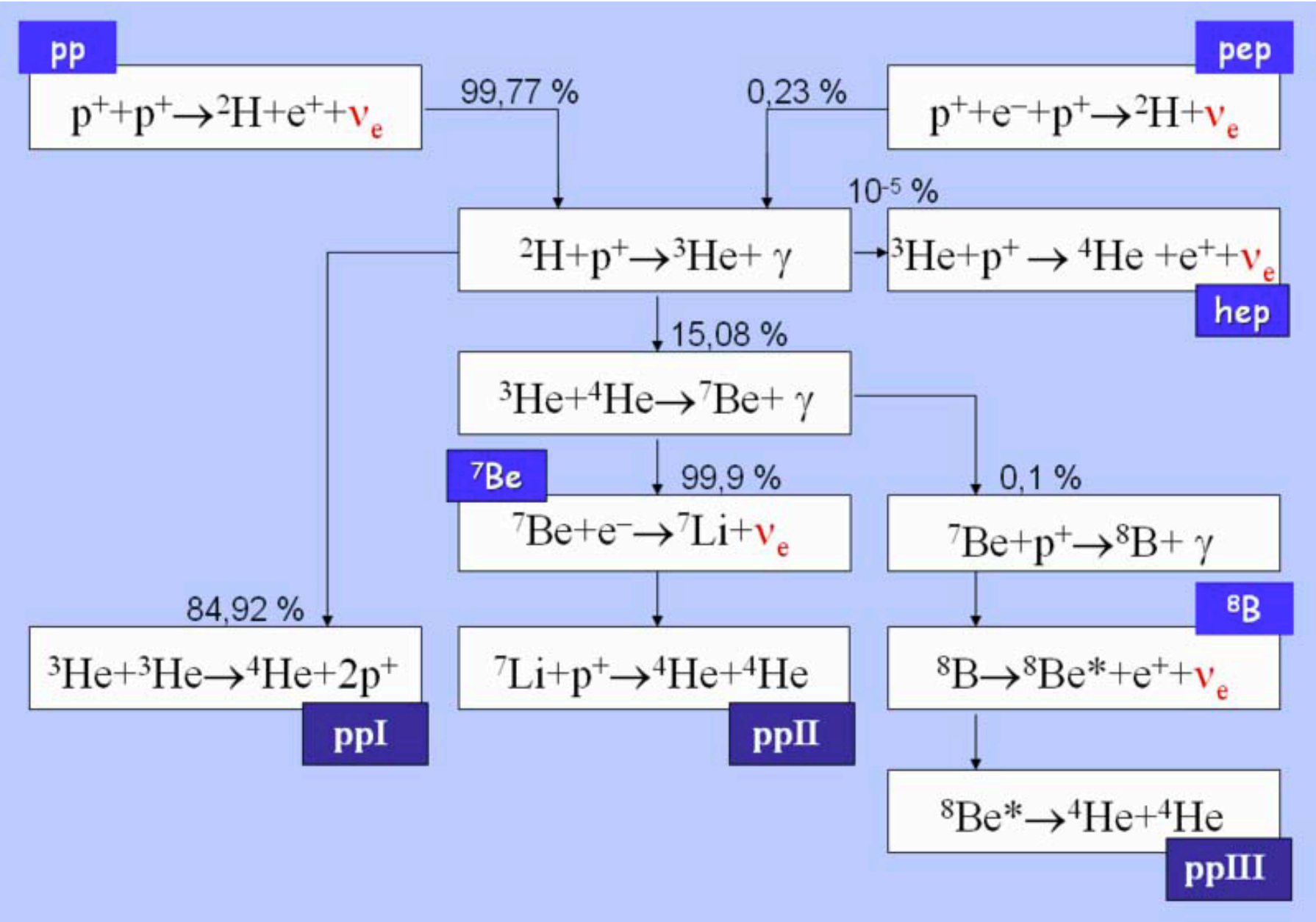
cos $\theta$ =1 down-going,  
 $L \approx 15$  km



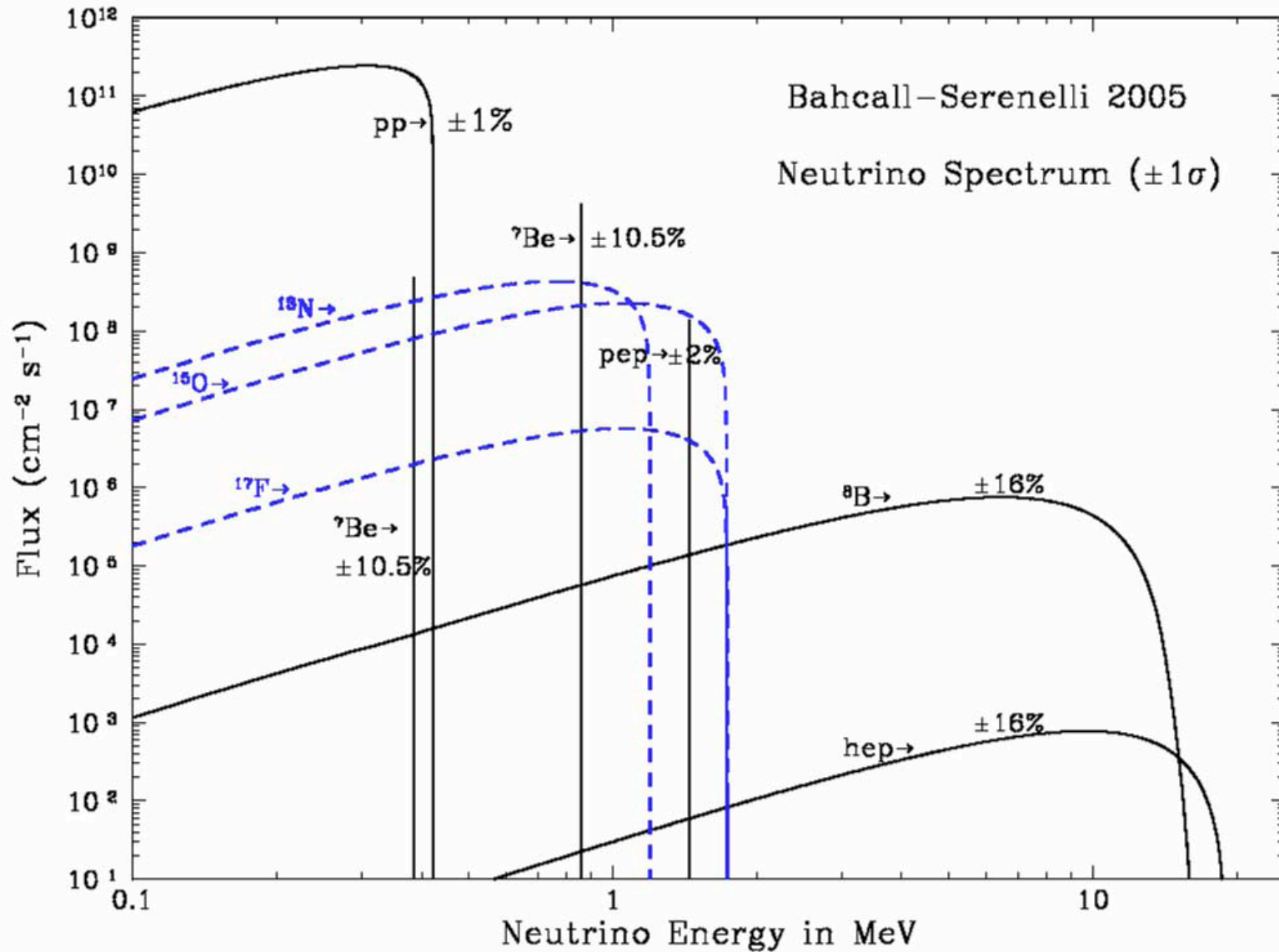
# Neutrinos from the Sun

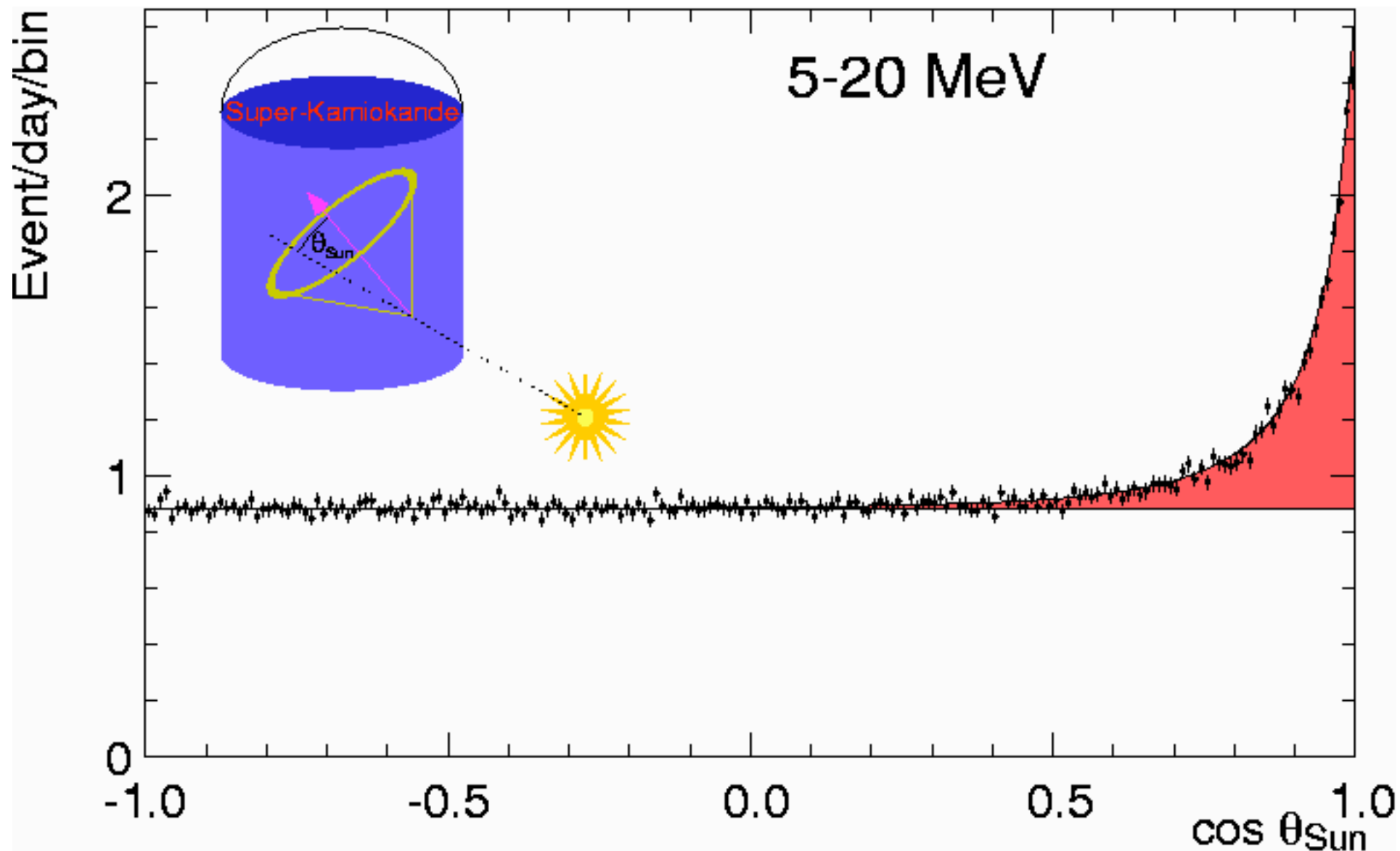
- Many mechanisms, all leading to electron neutrinos with varying energies.
  - Expect:  $0.5 \sin^2(2\theta)$  of solar model flux convolved with energy dependent efficiency.
- Neutrino energy too low to produce either muons or taus.
  - Electron disappearance experiments only in all but one experiment (SNO).

# Solar Model is Quite Complex



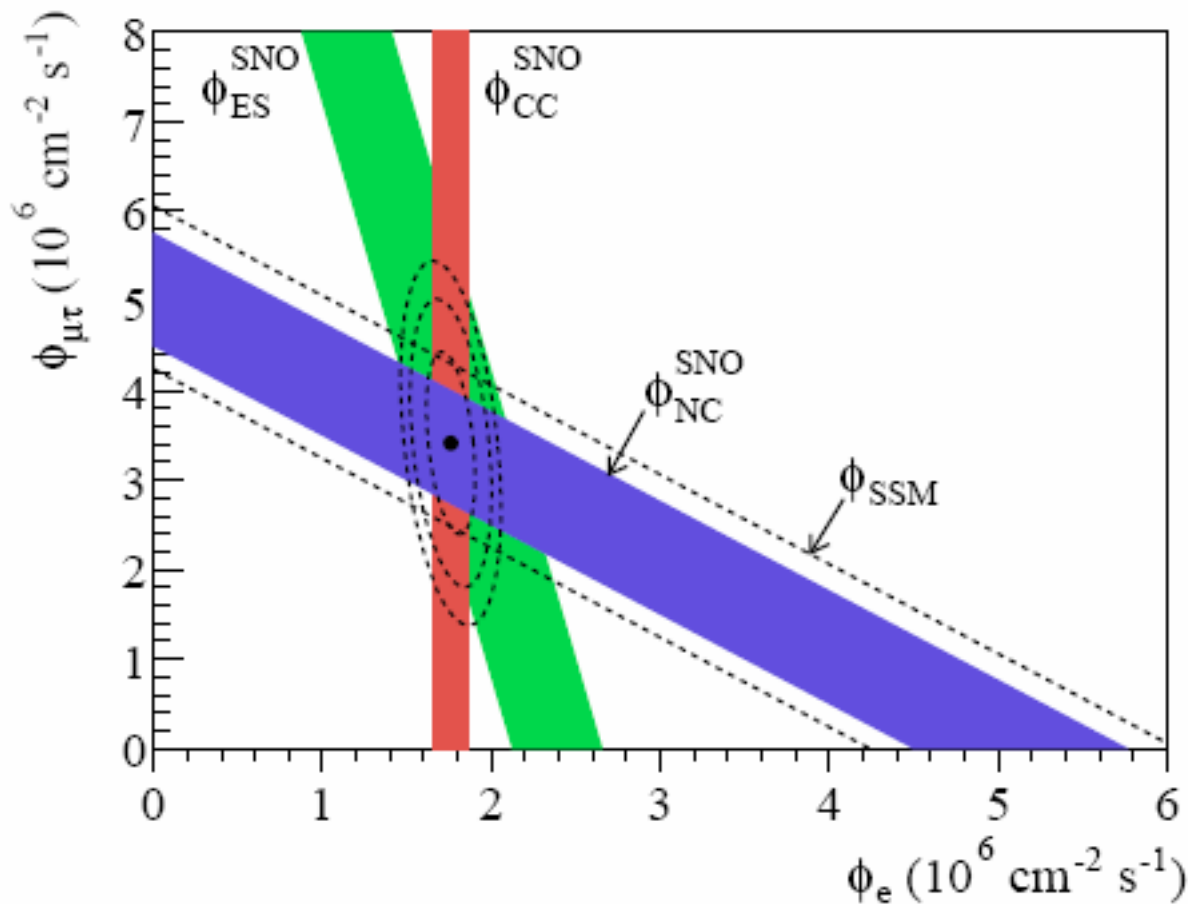
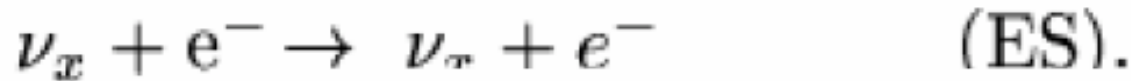
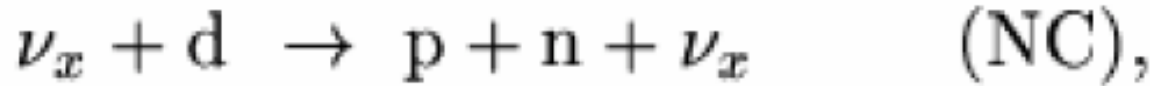
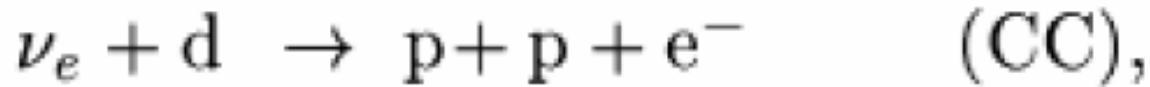
# Neutrino Energies are quite small Very Challenging Experimentally for many decades





$0.465 \pm 0.005^{+0.016}_{-0.015}$  of expectation

SNO allowed CC and NC, and was thus sensitive to all neutrino flavors => measures solar flux and electron neutrino flux.

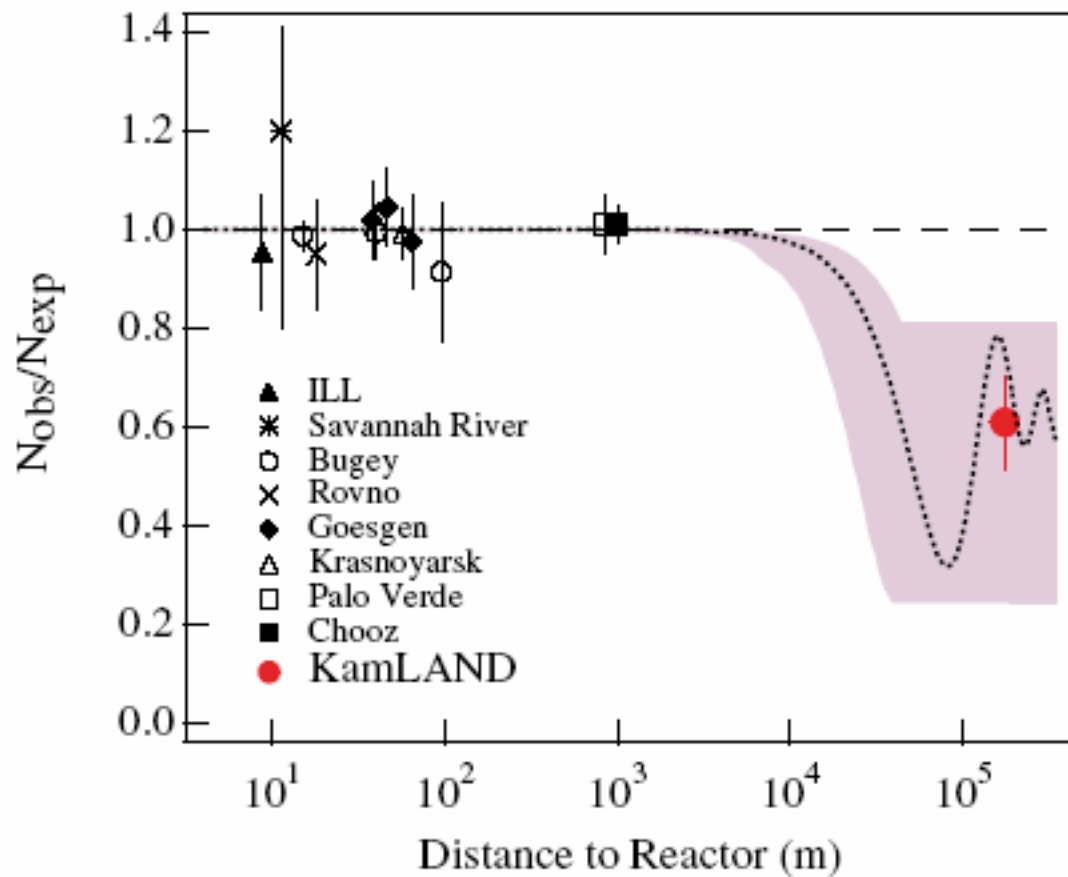


*Interpreted as*  
 $\nu_e \rightarrow \nu_\mu$



# Reactor Experiments

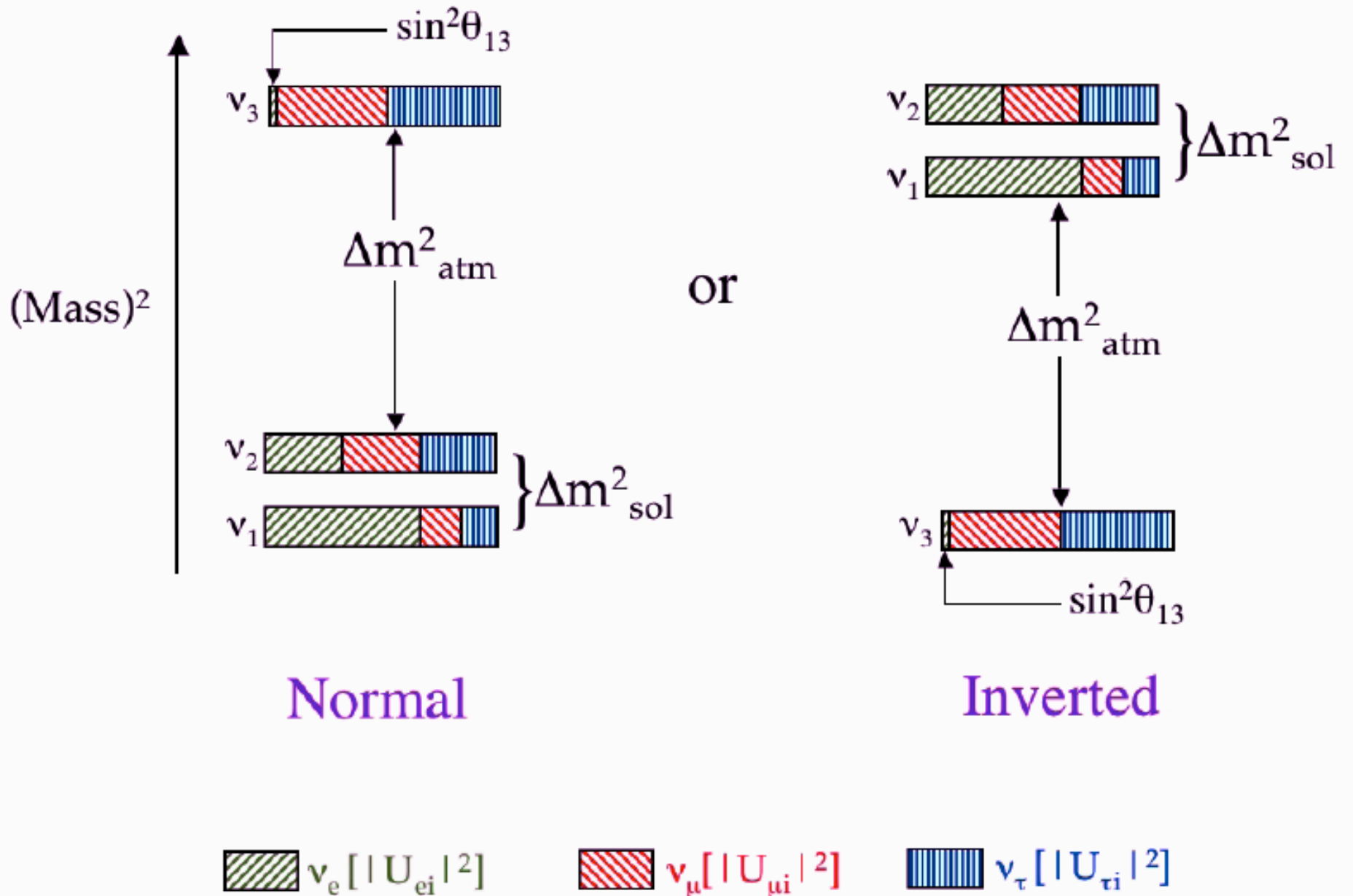
All except KamLAND had L that is too small!  
=> Only KamLAND saw oscillations !!!



# Interpretation

- Atmospheric must be  $\nu_{\mu} \rightarrow \nu_{\tau}$ 
  - Though tau appearance has never been seen.
  - However, electron appearance is ruled out.
  - The state that is far in mass from the other two must have very little electron neutrino content!

## Two Possible Mass Hierarchies



# Things we have not discussed yet.

- Majorana Neutrinos -> see **homework**
- “Size of CP violation” -> see **homework**
- Getting well collimated E via off-axis -> see **homework**
- Reactor neutrinos and  $\sin\theta_{13}$  -> see **homework**
- Resolving the mass hierarchy -> Wednesday.



