Physics 214 UCSD Physics 225a UCSB Experimental Particle Physics

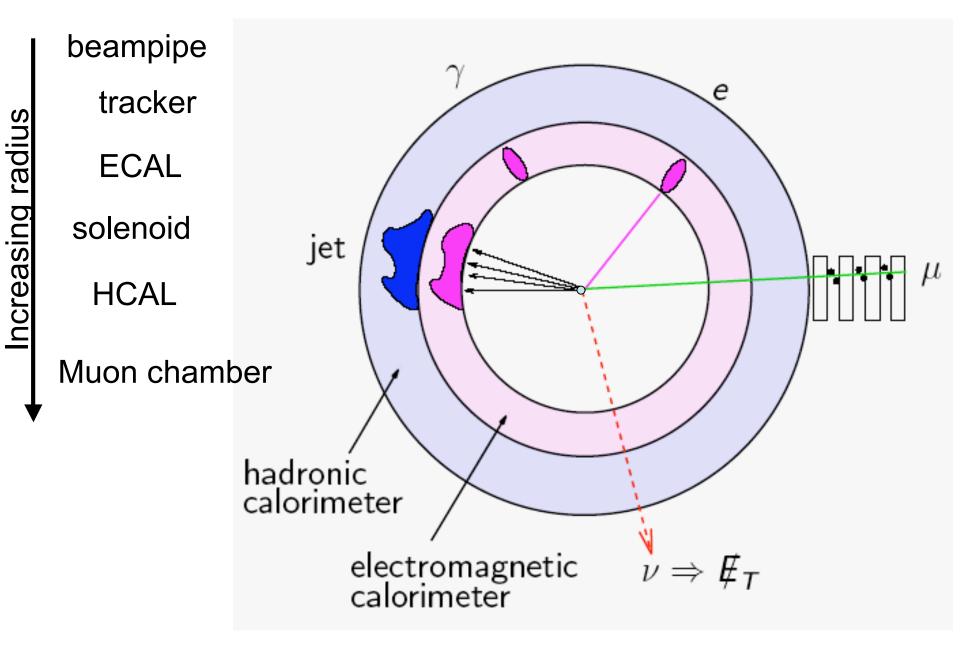
Lecture 17

A few words about the final Attempt at a summary of the quarter.

What do we need to detect?

- Momenta of all stable particles:
 - Charged: Pion, kaon, proton, electron, muon
 - Neutral: photon, K_{s}^{0} , neutron, K_{L}^{0} , neutrino
- Particle identification for all of the above.
- "Unstable" particles:
 - Pizero
 - b-quark, c-quark, tau
 - Gluon and light quarks
 - W,Z,Higgs
 - ... anything new we might discover ...

All modern collider detectors look alike



What you need to know:

- Exprimentalists:
 - Everything we talked about !
- Theorists:
 - The basic model how all collider detectors are built.

Symmetries

- Theorists:
 - Need to know everything, and more.
- Experimentalists:
 - Basic ideas
 - How to apply them to:
 - Know what transitions are allowed
 - Calculate ratios of amplitudes
 - Calculate angular distributions

Summary on Lie groups

- Let L be the N dimensional Lie group of Rank k for the Hamiltonian H.
- Then we have the following set of operators that mutually commute:

 $H, C_1, \ldots, C_k, L_1, \ldots, L_k$

- Any state is thus characterized by 2k quantum numbers.
- The energy E is given as some function of the C_1, \ldots, C_k .

Example:

Group of Rotations in 3-space

- Generators: J_x, J_y, J_z
- Lie algebra: $[J_k, J_l] = i \epsilon_{klm} J_m$
- Rank = 1
- Casimir Operator: J²
- Multiplets are classified by their total angular momentum J
- States are classified by J and J_z, the latter being one of the three generators.

Final Exam Question 3

• What's the angular distribution for the decay products in the decay $J^{PC} = 1^{--}$ to two pseudoscalars if the decaying particle has $J_z = +-1$? __Anti-B

e- θ e+ $J_z=0$ in final state, therefore L must be encoded in angular distribution of B anti-B axis !!! $\sigma \propto \left(\left|d_{10}^1\right|^2 + \left|d_{-10}^1\right|^2\right) \propto \sin^2 \theta$

Final Exam Question 8

- Most of you got some fraction of this right.
- Almost everybody missed the fact that there are two reduced matrix elements in play, and that you need to do an isospin decomposition after adding the isospin of the B meson and the isospin of the transition operator $H_{effective}$!!!

Neutrino Physics

- The basic Phenomenology
 - There are 3 families of neutrinos
 - They are produced and observed in their weak eigenstates but propagate in their mass eigenstates => mixing
 - What we know from experiment
 - The research frontier
 - Measuring sin(theta13)
 - Understand hierarchy
 - Majorana vs Dirac

Discussion of QFT and scattering of point particles

- From Feynman Diagrams to Matrix elements.
 - You need to come up with the diagram for a process (e.g. Q6 of final)
 - You need to write down the ME
- Relationship between ME, cross section, and physical observable.
- Understand basic characteristics to know when you've screwed up royally.
- Theorists need to actually master the algebra.

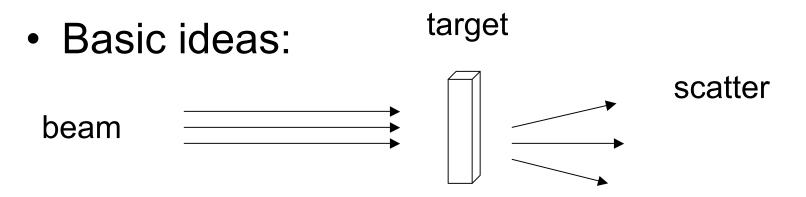
"Rules" for dealing with antiparticles

- Antiparticles get arrow that is backwards in time.
- "Incoming" and "outgoing" is defined by how the arrows point to the vertex.
- Antiparticles get negative energy assigned.

Two-by-two process

- Understand how to relate number of scatters in AB -> CD scattering to "beam & target independent" cross section in terms of W_{fi}.
- Relate W_{fi} to matrix element.
- => Understand relationship between cross section and Matrix Element", and be able to relate it to physical observable.

Cross Section for AB -> CD



of scatters = (flux of beam) x (# of particles in target) x σ

 $\begin{array}{l} \mathsf{W}_{\mathsf{fi}} \\ \mathsf{Cross \ section} = \ \sigma = \underbrace{\mathsf{w}_{\mathsf{fi}}} \\ (\mathsf{initial \ flux}) \end{array} \end{array} (\mathsf{number \ of \ final \ states}) \\ \end{array}$

W_{fi} = rate per unit time and volume

"Cross section" is independent of characteristics of beam and target !!!

Experimental perspective

- We measure event yield within some ragged kinematic corner of phase space.
- We divide by our detector acceptance
- => We get produced yield for well defined corner of phase space.
- We measure integrated luminosity of the colliding beams for our data taking period.

=> σ = produced yield / integrated luminosity

Theoretical perspective

 W_{fi} Cross section = σ = — (number of final states) (initial flux) $\frac{VV_{fi}}{V_{A} (2E_{A}/V) (2E_{B}/V)} \frac{Vdp_{C}^{3}}{(2\pi)^{3} 2E_{C}} \frac{Vdp_{D}^{3}}{(2\pi)^{3} 2E_{D}}$ $\sigma =$ $W_{fi} = (2\pi)^4 \frac{\delta^{(4)}(p_D + p_C - p_A - p_B)}{v^4} |M|^2$

M is obtained from Feynman rules, and the rest is algebra.

It is customery to re-express

$$d\sigma = \frac{\left|M\right|^2}{F} dQ$$

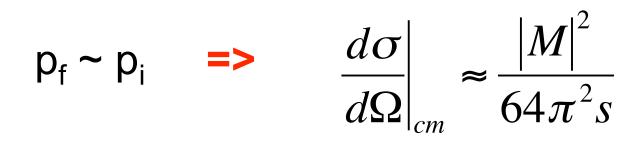
F = flux factor:
$$F = 4\sqrt{(p_A^{\mu} p_{\mu}^{B})^2 - m_A^2 m_B^2}$$

dQ = Lorentz invariant phase space:

$$dQ = \frac{1}{16\pi^2} \delta^{(4)} (p_D + p_C - p_A - p_B) \frac{dp_c^3 dp_D^3}{E_c E_D}$$

e+e- -> f+ f-

- Question 1 on final is probably the most important example process to understand.
- I promise to revisit this next quarter, and discuss it in some detail then.



$$t = -2 k^{2} (1 - \cos\theta)$$

$$u = -2 k^{2} (1 + \cos\theta) \implies |M|^{2} = 2e^{4} \frac{8k^{4} (1 + \cos^{2}\theta)}{16k^{4}}$$

$$s \sim 4k^{2}$$

 $\alpha = e^2 / 4\pi$

Relativistic limit

$$\frac{d\sigma}{d\Omega}\Big|_{cm} \approx \frac{\alpha^2}{4s} \left(1 + \cos^2\theta\right)$$
$$\sigma(e^+e^- \to \mu^+\mu^-) \approx \frac{4\pi\alpha^2}{3s}$$

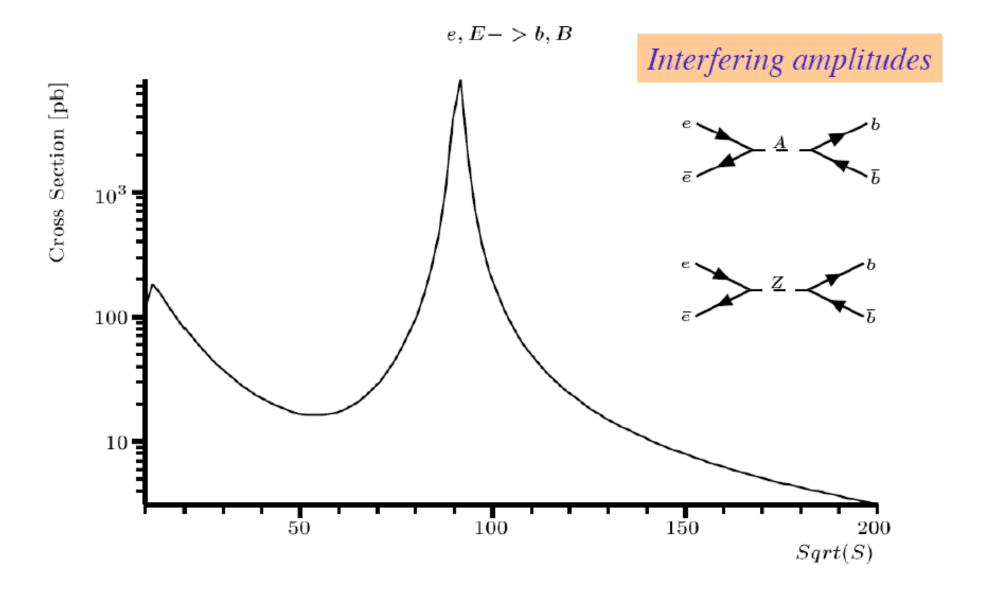
- e_L⁻ e_R⁺ -> mu_L⁻ mu_R⁺
- $e_{L}^{-} e_{R}^{+} -> mu_{R}^{-} mu_{L}^{+}$
- $e_R^- e_L^+$ -> $mu_L^- mu_R^+$
- e_R⁻ e_L⁺ -> mu_R⁻ mu_L⁺
- $J_z + 1 -> +1$ $J_z + 1 -> -1$ $J_z -1 -> +1$ $J_z -1 -> -1$
- Next look at the rotation matrices:

$$d_{11}^{1}(\theta) = \frac{1}{2} (1 + \cos \theta) \approx \frac{-u}{s}$$
$$d_{-1-1}^{1}(\theta) = \frac{1}{2} (1 + \cos \theta) \approx \frac{-u}{s}$$
$$d_{-1-1}^{1}(\theta) = \frac{1}{2} (1 - \cos \theta) \approx \frac{-t}{s}$$
$$d_{-1-1}^{1}(\theta) = \frac{1}{2} (1 - \cos \theta) \approx \frac{-t}{s}$$

Cross products cancel in Spin average:

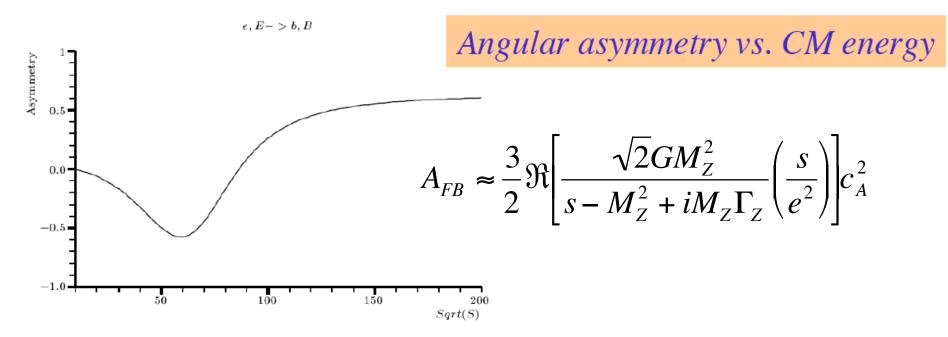
$$\overline{\left|M\right|^2} \propto \left(1 + \cos^2\theta\right)$$

QED piece goes like 1/E² => -2logE on log scale. Weak piece must have a Z-pole.



Forward-backward asymmetry

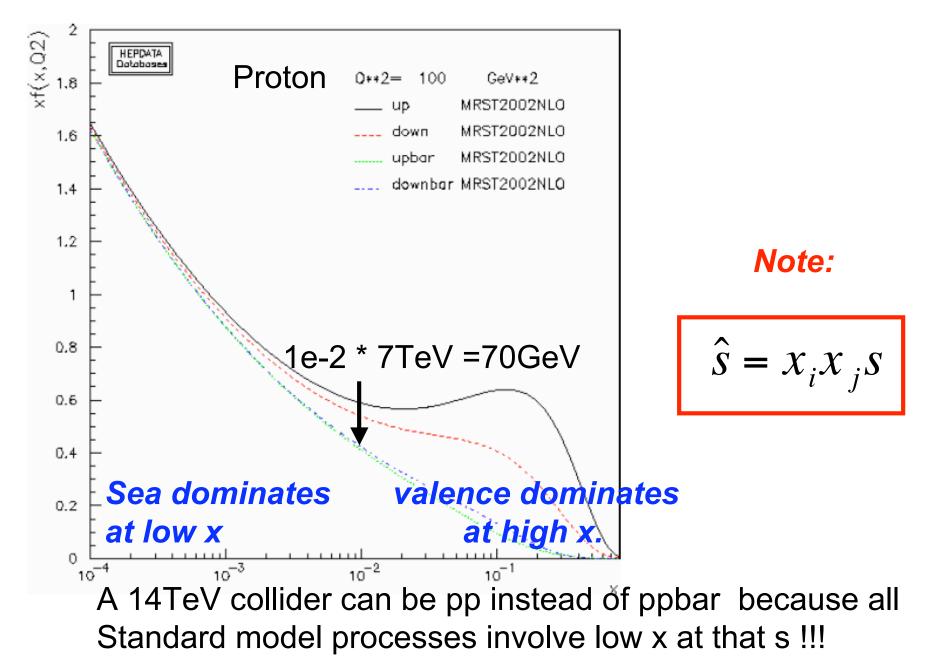
Must start at 0 because of 1+cos² dependence. For dependence look at Eq. 13.66 and 13.61 in H&M.



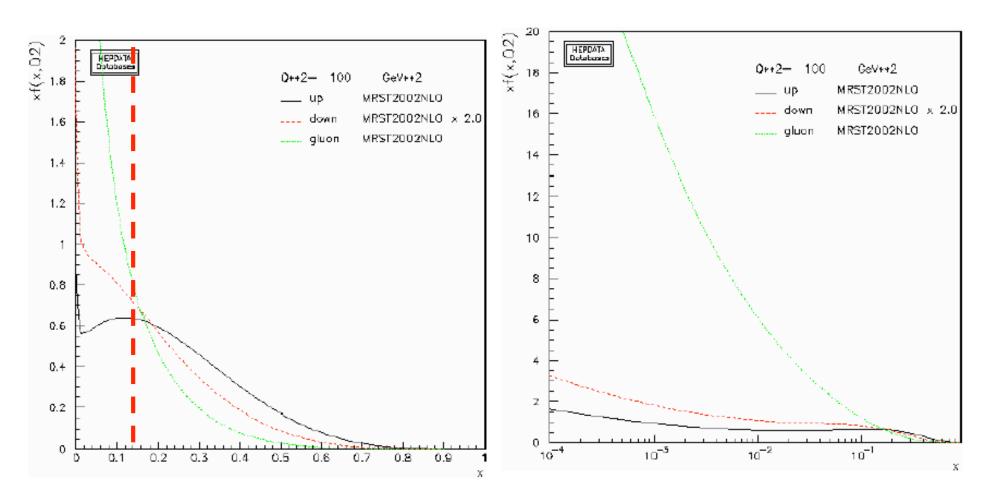
It goes negative like -s as s increases from 0. It reaches a minimum t a place that's not immediatey obvious. It goes asymptotically towards a positive value.

Deep Inelastic Scattering

- Bjorken Scaling is a sign of point particles inside the proton.
- The parton picture, and pdf's that describe the structure of the proton
 - At what x do valence quarks dominate
 - At what x do sea quarks dominate
 - At what x do gluons dominate
 - How do I use this information to gain some intuition about proton proton and proton antiproton collisions as a function of sqrt(s).



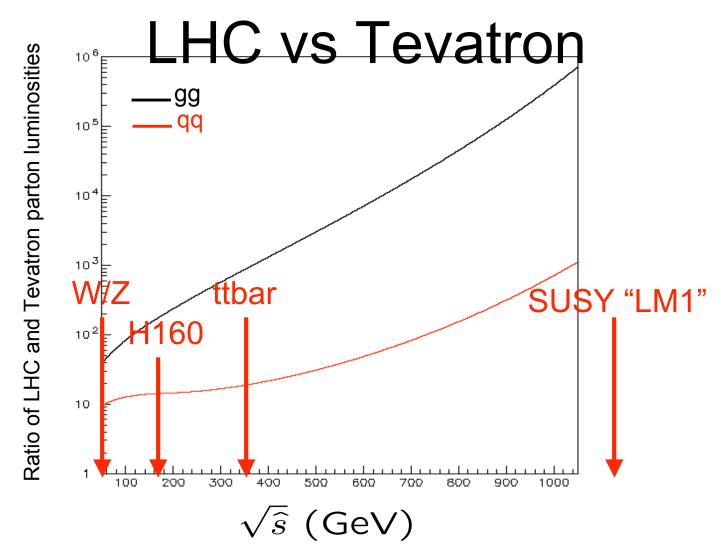
PDFs from http://durpdg.dur.ac.uk/hepdata/pdf3.html



Gluons dominate at low x.

To set the scale, x = 0.14 at LHC is 0.14 * 7TeV = 1TeV

=> The LHC is a gluon collider !!!



simplistic rule of thumb:

- For 1 TeV gg processes, 1 fb⁻¹ at FNAL is like 1 nb⁻¹ at LHC
- For 1 TeV qq processes, 1 fb⁻¹ at FNAL is like 1 pb⁻¹ at LHC

Cross sections at 1.96TeV versus 14TeV Tevatron vs LHC

	Cross section		Ratio
Ζ→μμ	260pb	1750pb	6.7
WW	10pb	100pb	10
H _{160GeV}	0.2pb	25pb	125
mSugra _{LM1}	0.0006pb	50pb	80,000

At 10³²cm⁻²s⁻¹ CMS might accumulate 10pb⁻¹ in one day!

... and SUSY might not exist in nature.

In case you want to prepare for next quarter during the break.

Make yourself comfortable with question 1 on final. Play around a bit with comphep, madgraph, etc.

Have a great holiday!

And see you all back in the new year in the continuation of this lecture.